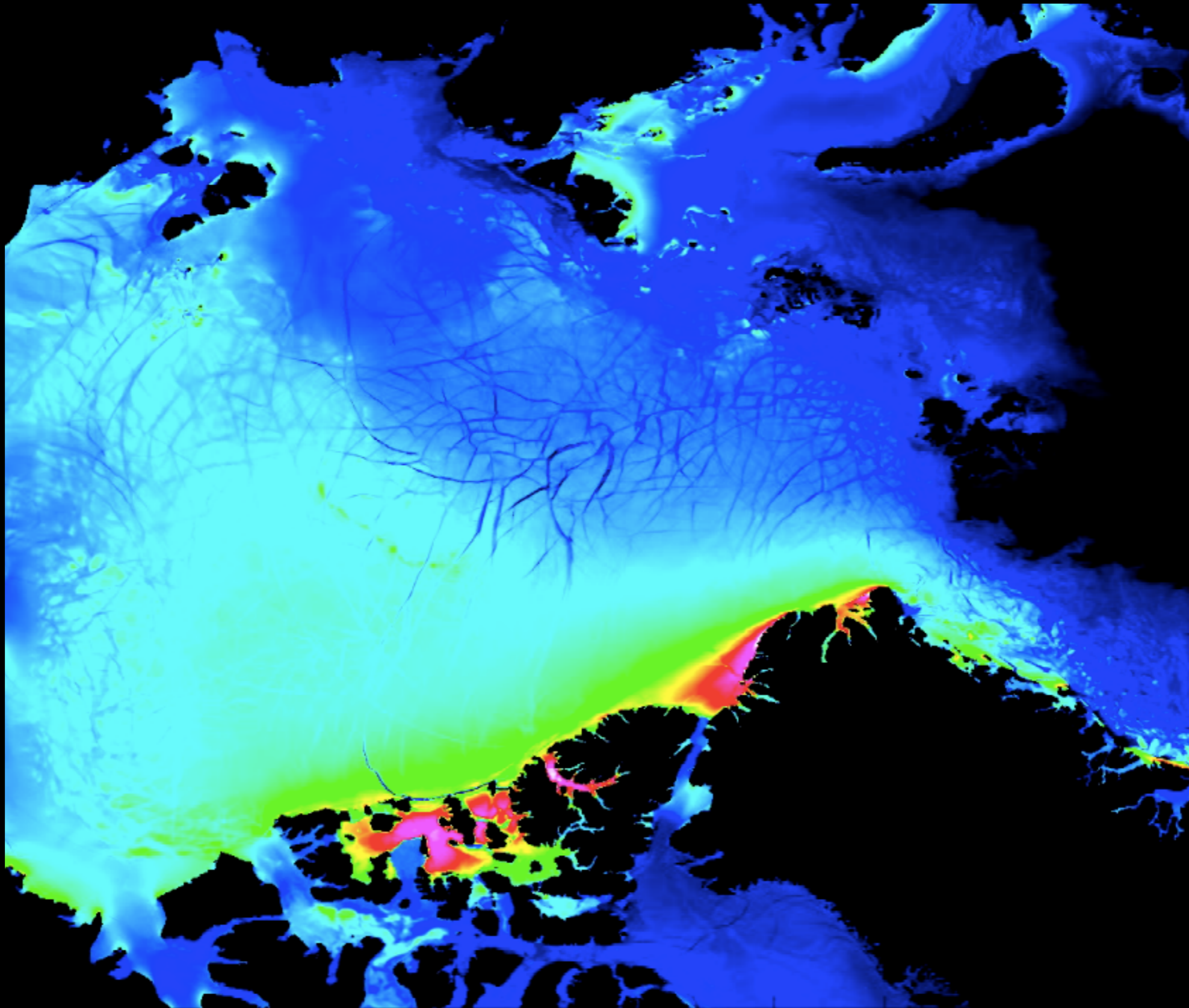


Convergence and accuracy of sea ice dynamics solvers

Martin Losch (AWI)

Annika Fuchs (AWI), Jean-François Lemieux (EC)



Sea ice thickness with MITgcm (JPL)

Review of 2D momentum equations

$$m \frac{D\mathbf{u}}{Dt} = -m f \mathbf{k} \times \mathbf{u} + \boldsymbol{\tau}_{air} + \boldsymbol{\tau}_{ocean} - m \nabla \phi(0) + \mathbf{F}, \quad (1)$$

$F_j = \partial_i \sigma_{ij}$ = divergence of symmetric stress tensor of rank 2

Viscous-Plastic (VP) constitutive law (rheology):

$$\sigma_{ij} = 2\eta \dot{\epsilon}_{ij} + [\zeta - \eta] \dot{\epsilon}_{kk} \delta_{ij} - \frac{P}{2} \delta_{ij}. \quad (2)$$

EVP equations (NOT intended to be a different rheology):

$$\frac{1}{E} \frac{\partial \sigma_{ij}}{\partial t} + \frac{1}{2\eta} \sigma_{ij} + \frac{\eta - \zeta}{4\zeta\eta} \sigma_{kk} \delta_{ij} + \frac{P}{4\zeta} \delta_{ij} = \dot{\epsilon}_{ij}. \quad (3)$$

$$\dot{\epsilon}_{ij} = \frac{1}{2} (\partial_i u_j + \partial_j u_i) = \text{strain rates}$$

$$P = P^* hc \cdot e^{-C^*(1-c)} \quad \zeta = \min \left(\frac{P}{2\Delta}, \zeta_{\max} \right) \quad \eta = \frac{\zeta}{e^2}$$

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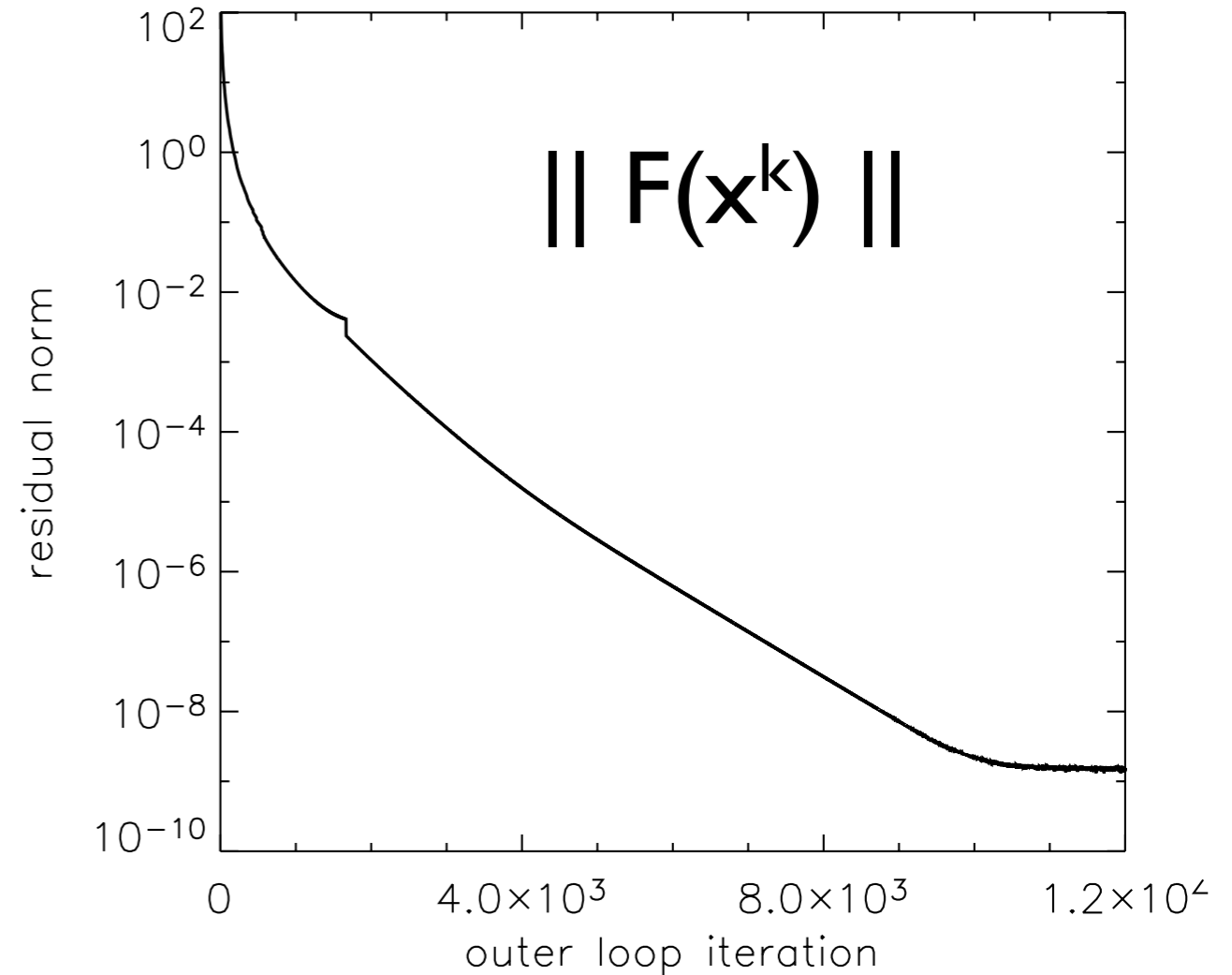
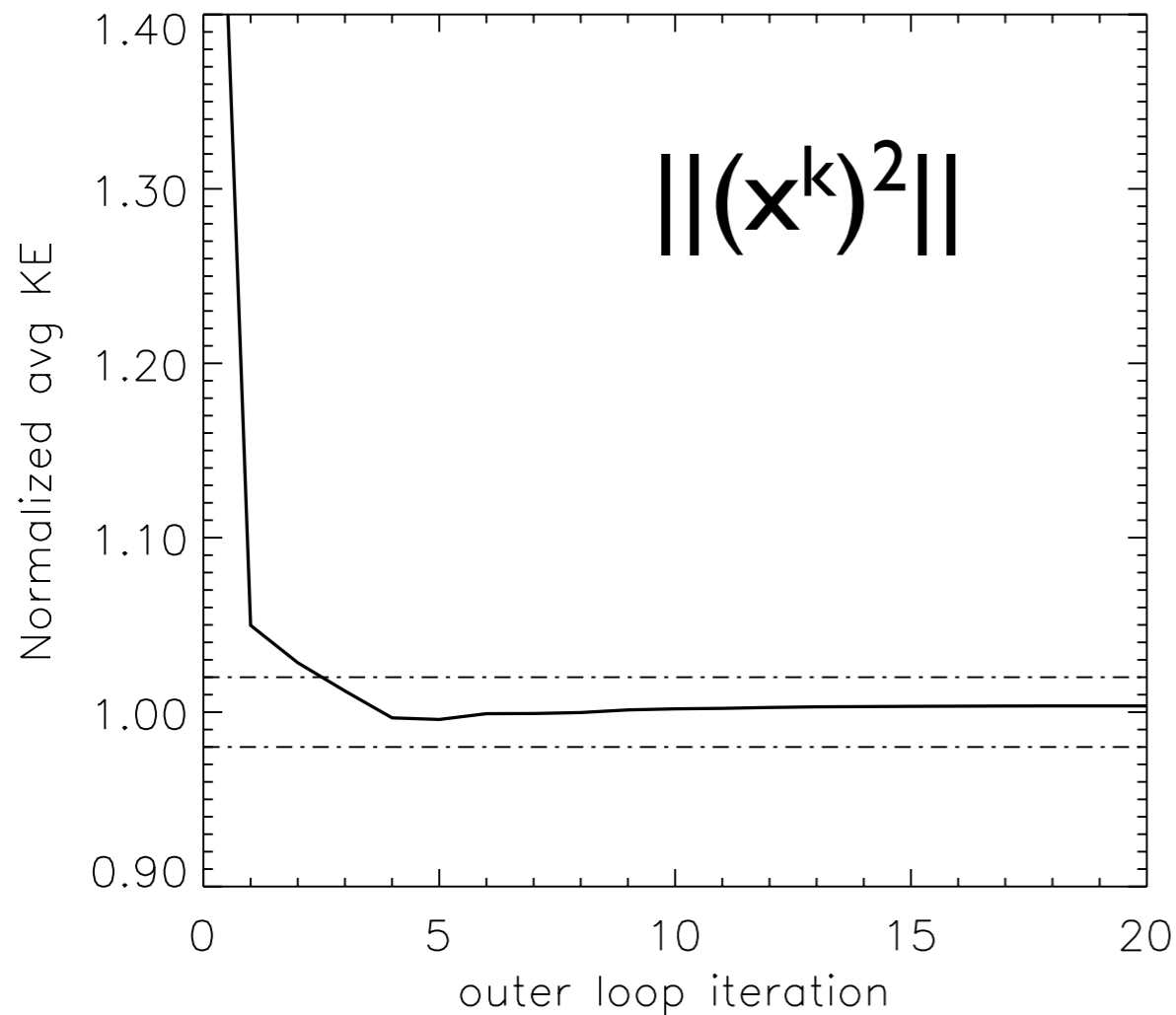
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classical implicit VP-solver: the Picard solver
for $A(x) x = b(x)$

$$A(x^{k-1}) x^{k-1} - b(x^{k-1}) = F(x^{k-1})$$

outer (non-linear) loop:

until (some criterion), solve $A(x^{k-1}) x^k = b(x^{k-1})$

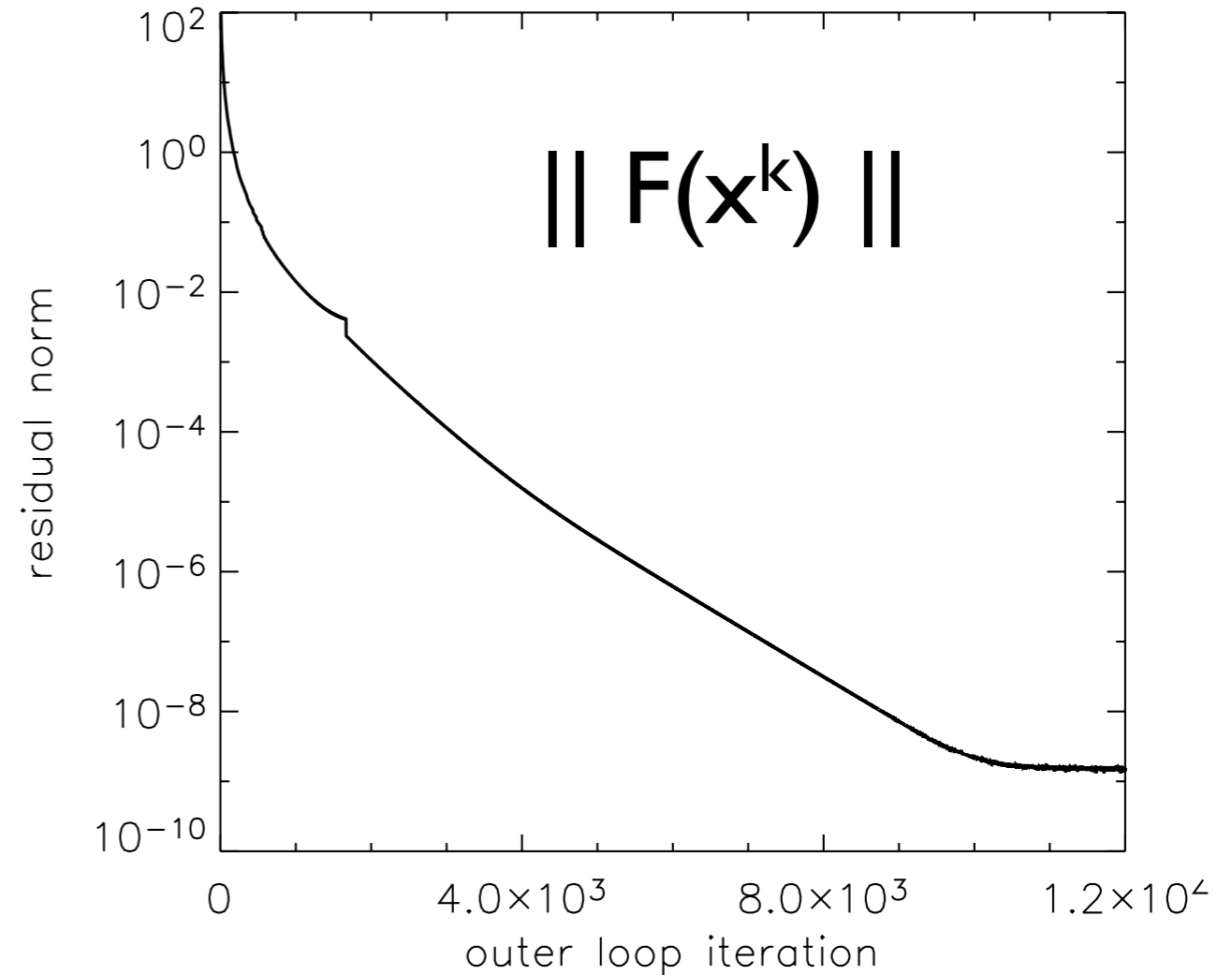
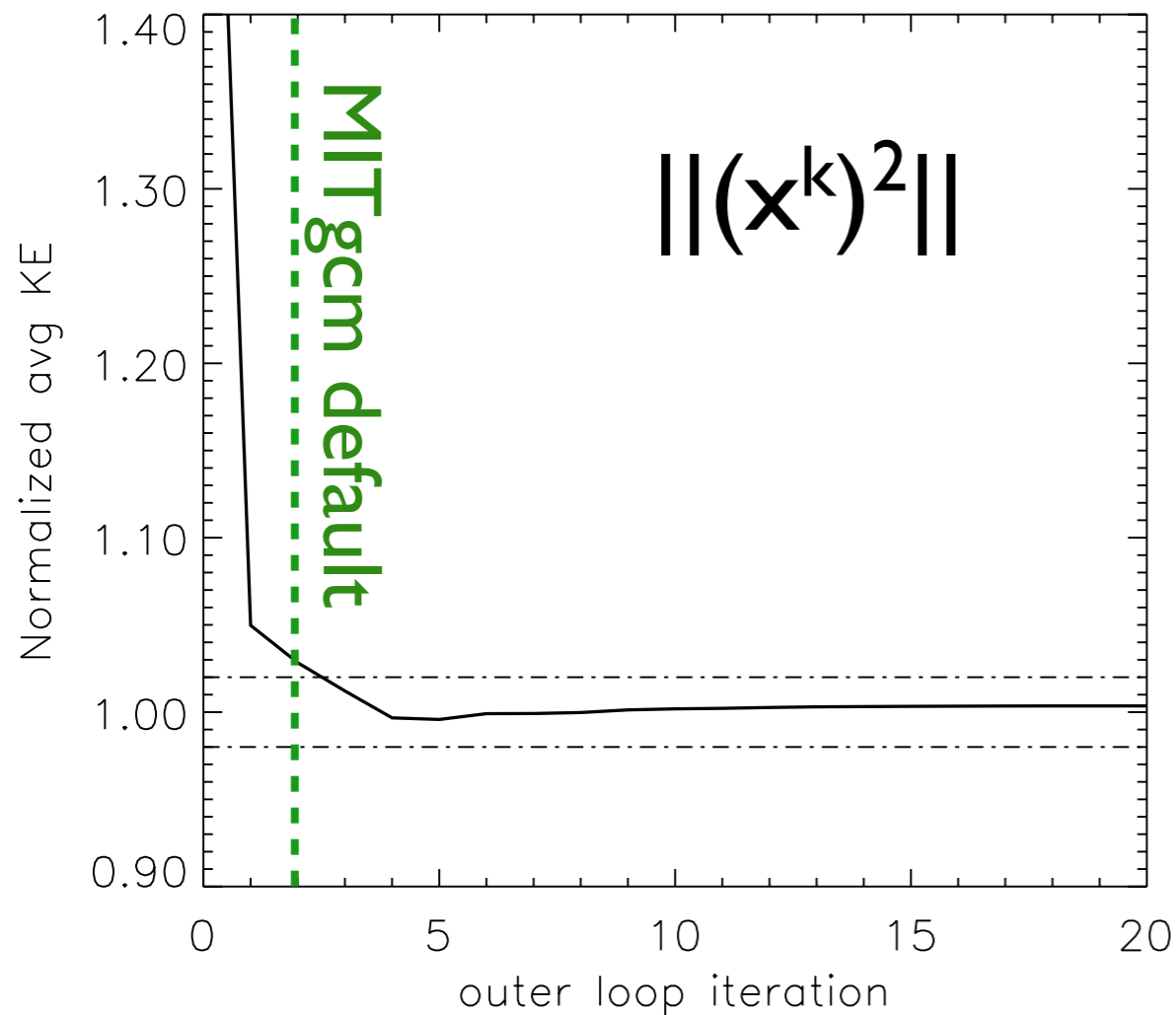


classical implicit VP-solver: the Picard solver
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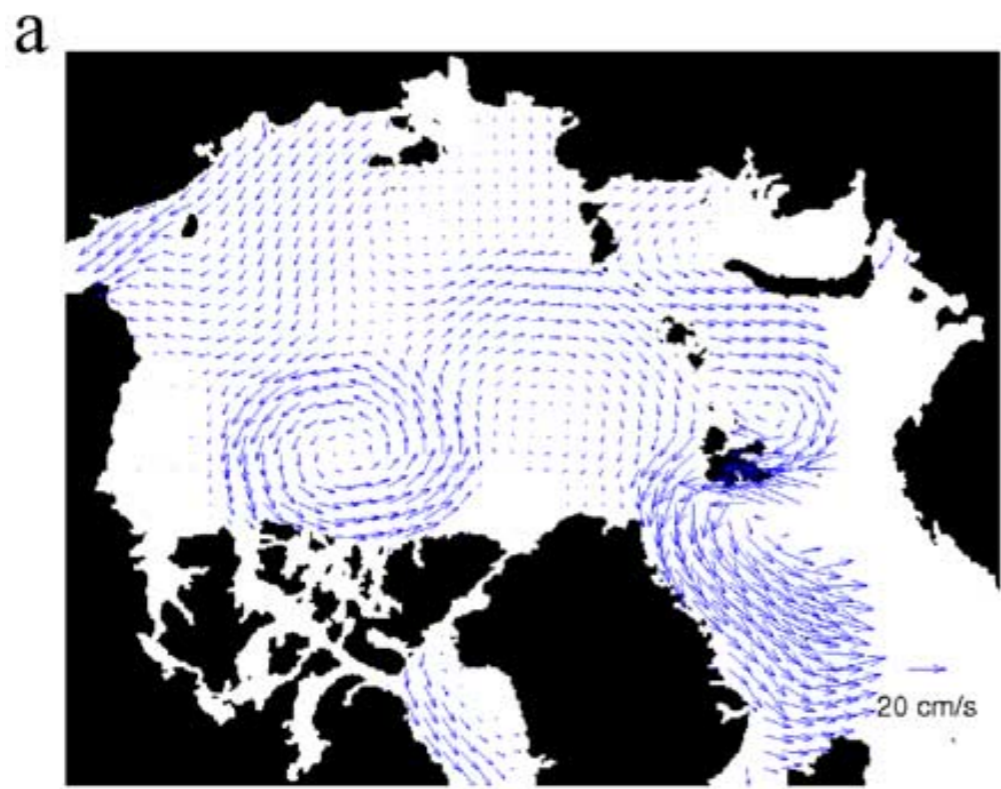
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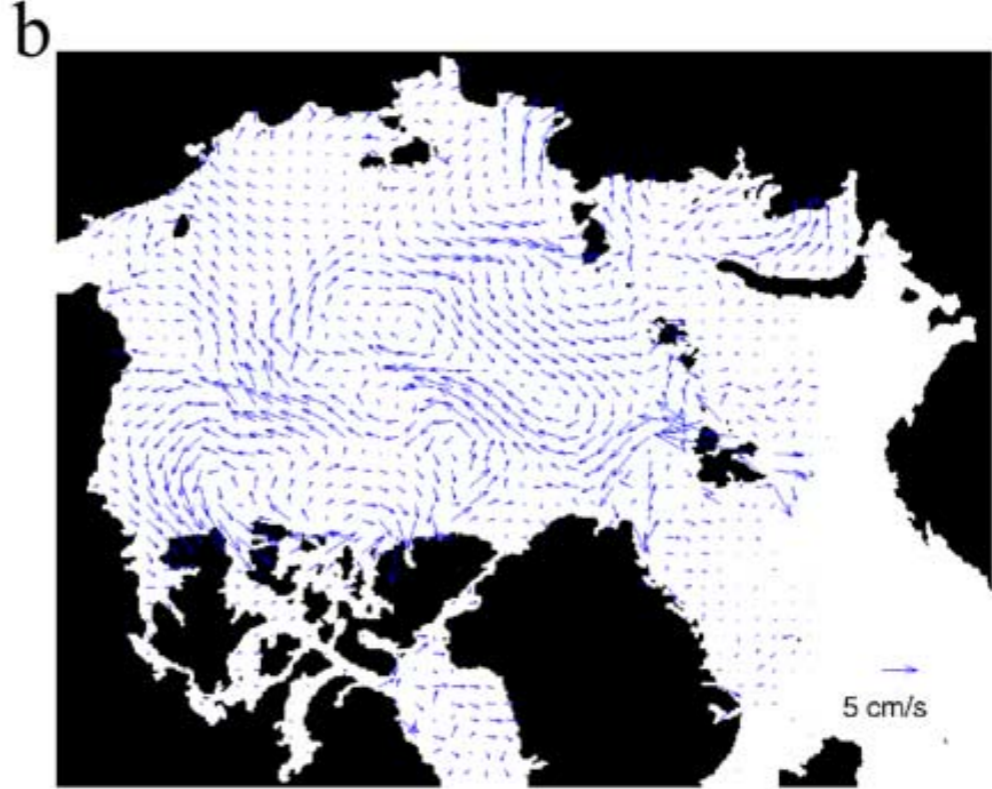
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Lemieux and Tremblay (2009)

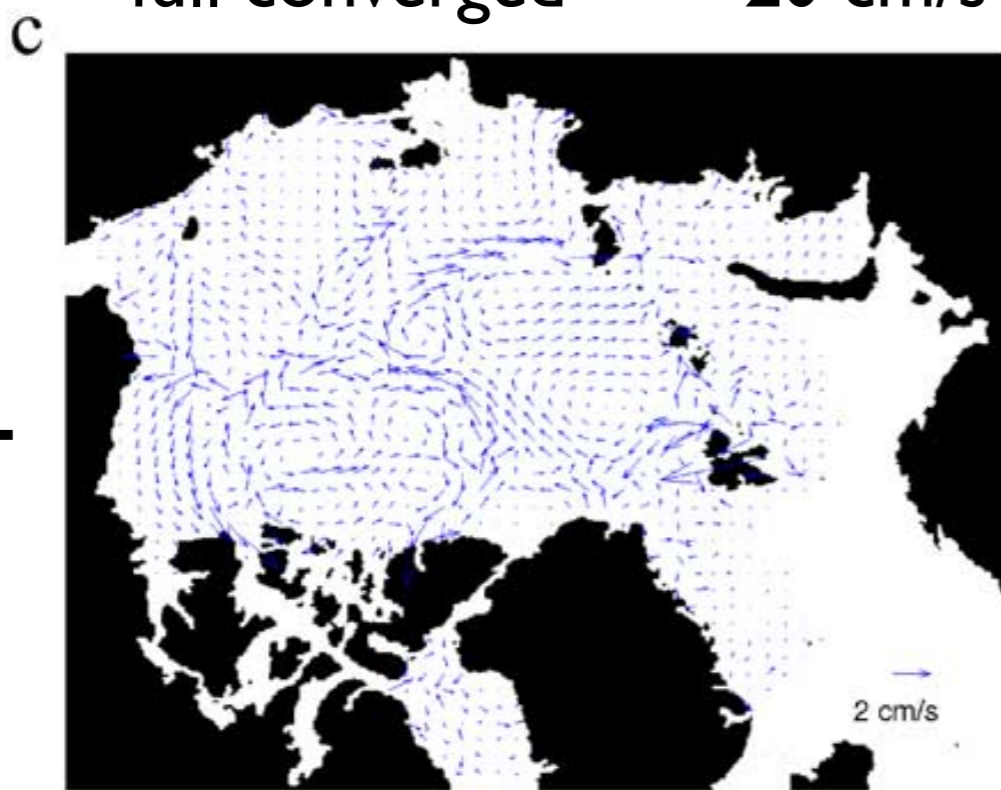


full converged 20 cm/s



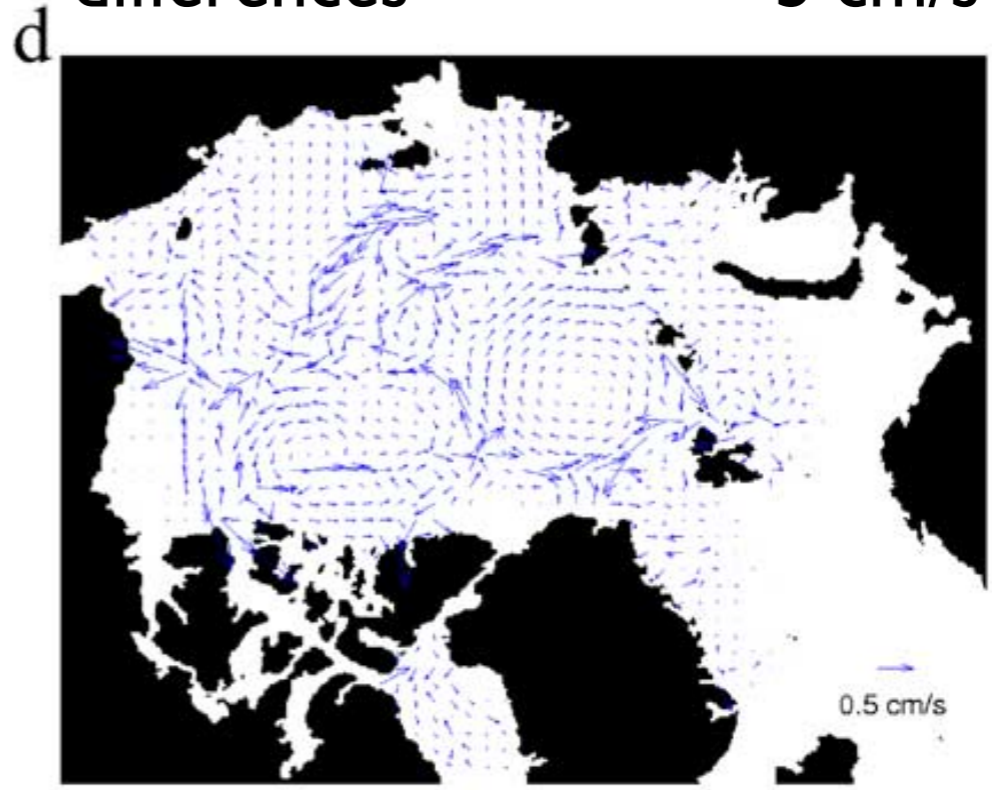
differences 5 cm/s

2 OL



10 OL

2 cm/s



40 OL

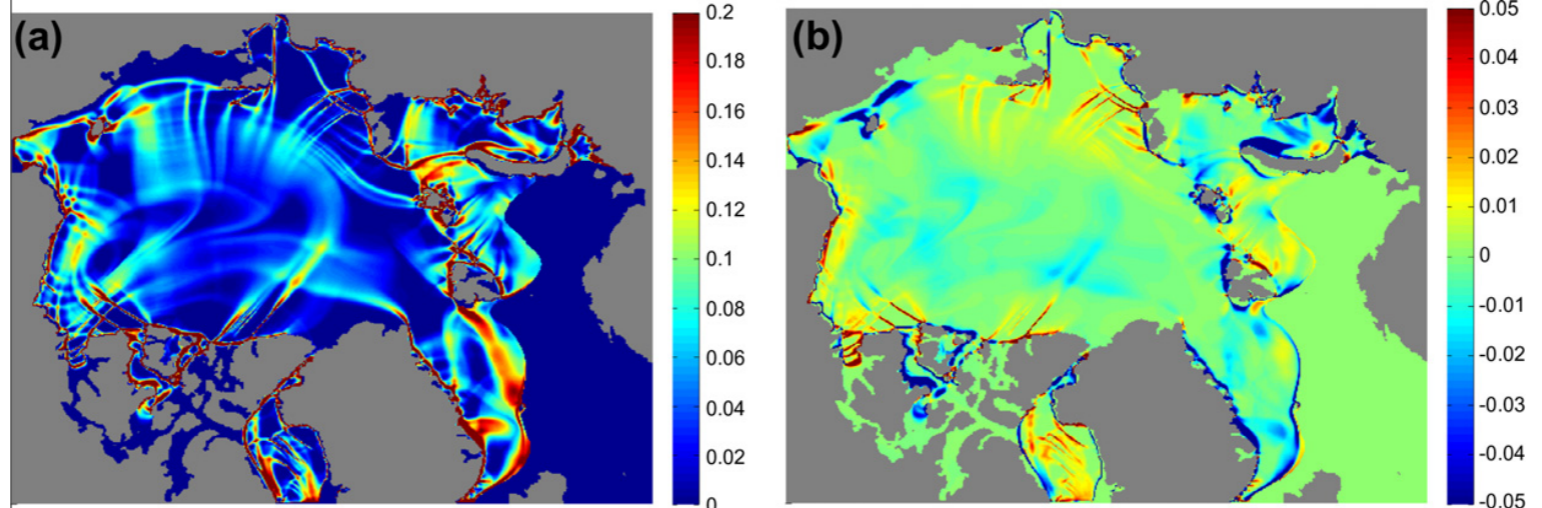
0.5 cm/s

Lemieux and Tremblay (2009)

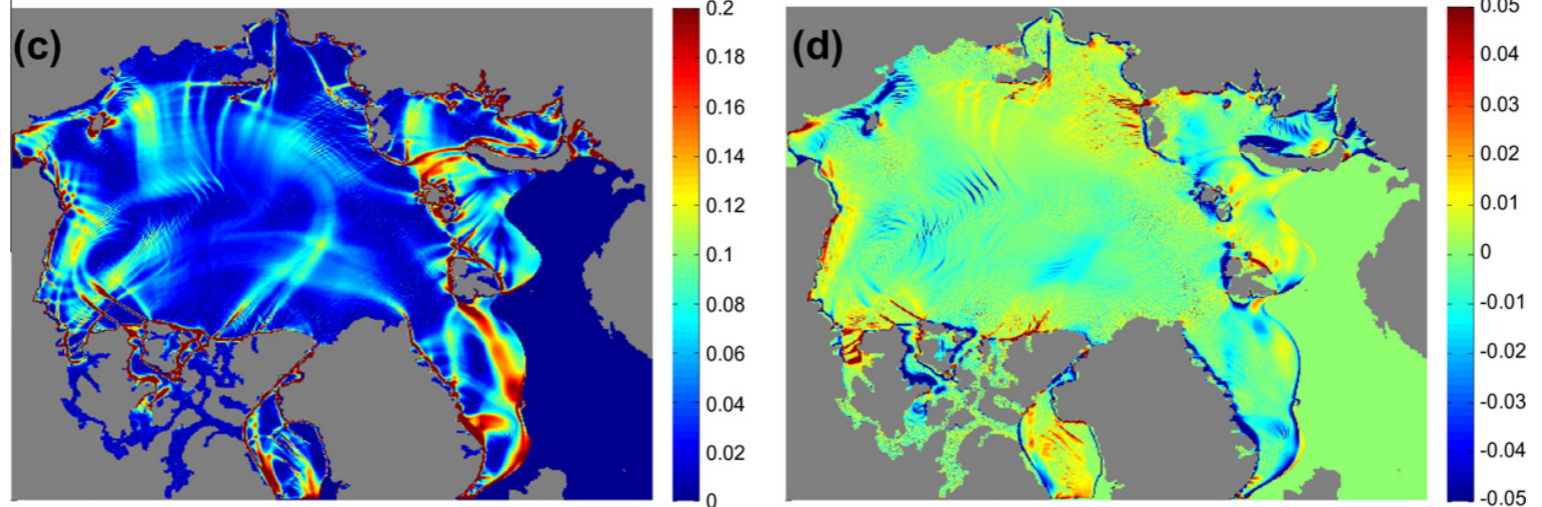
(for long time steps)

EVP isn't any better (but faster)

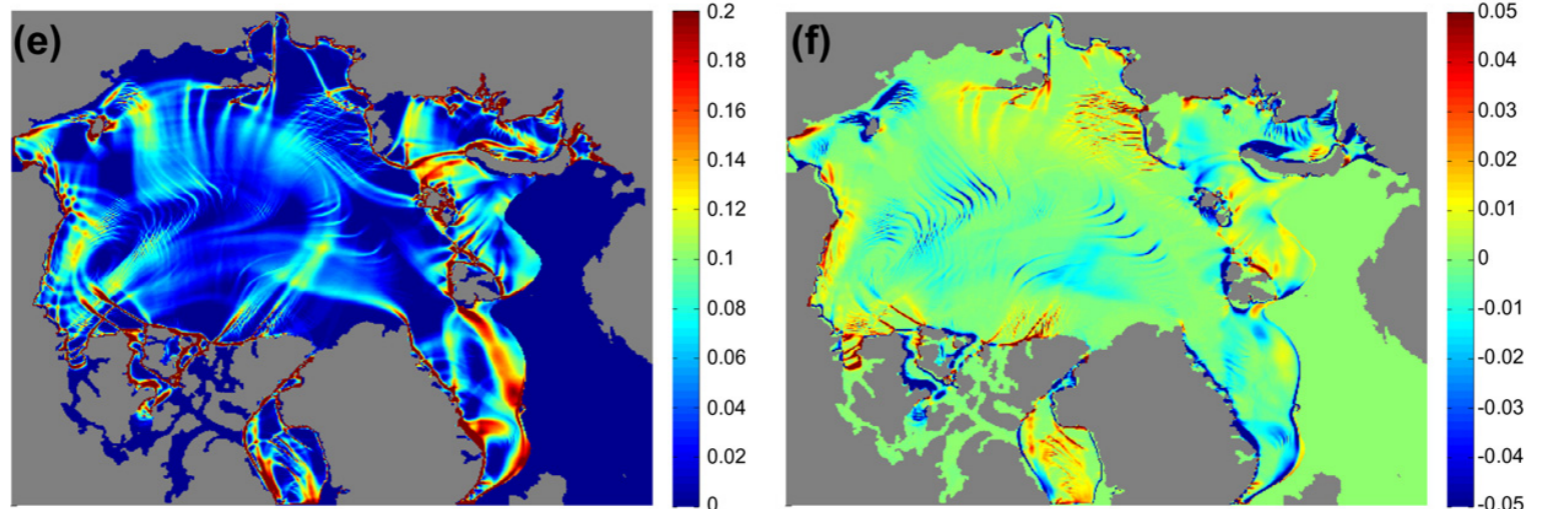
reference



EVP, 120
sub-cycles



EVP, 1980
sub-cycles

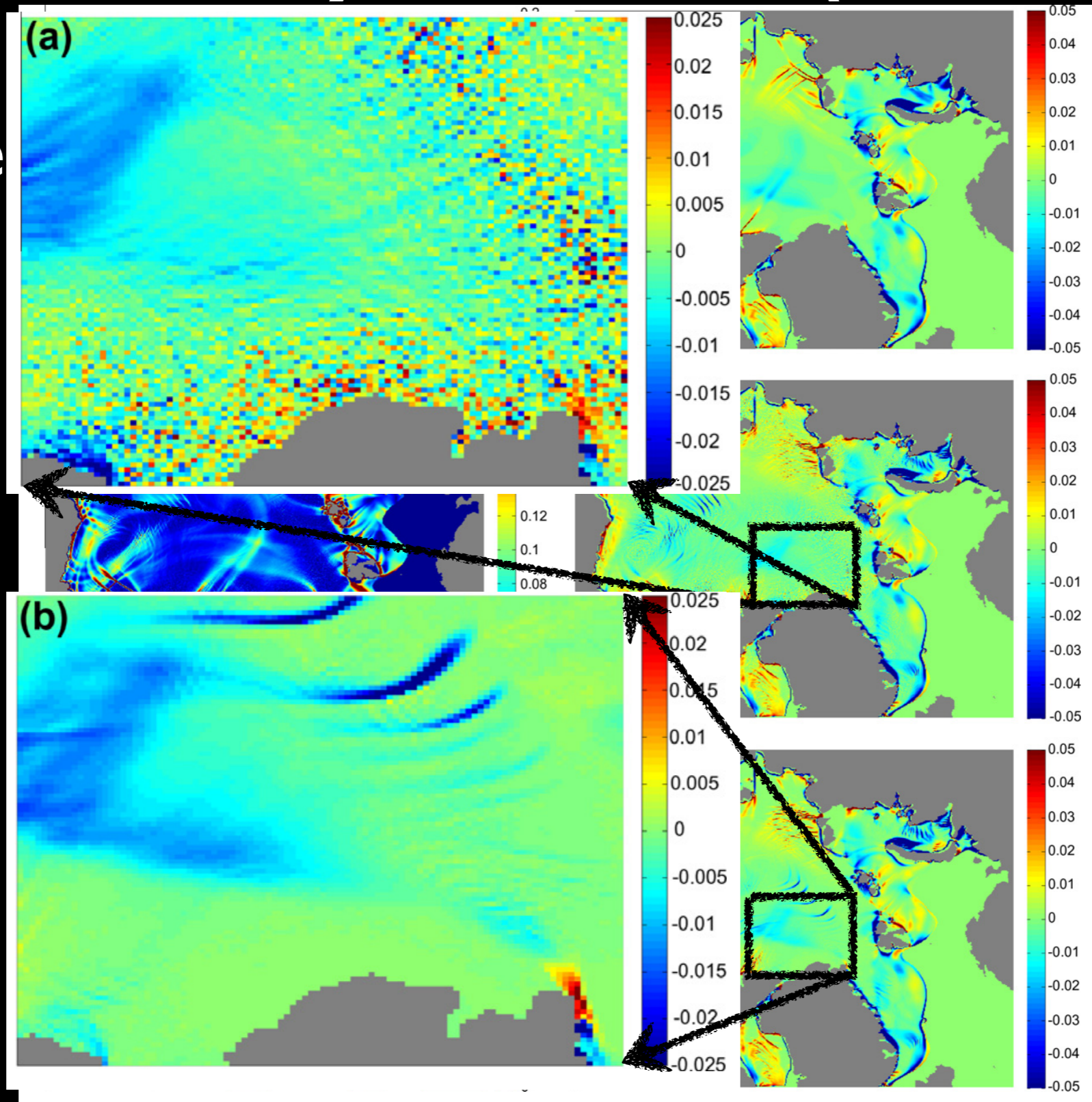


Lemieux et al. (2012), shear and divergence (per day)

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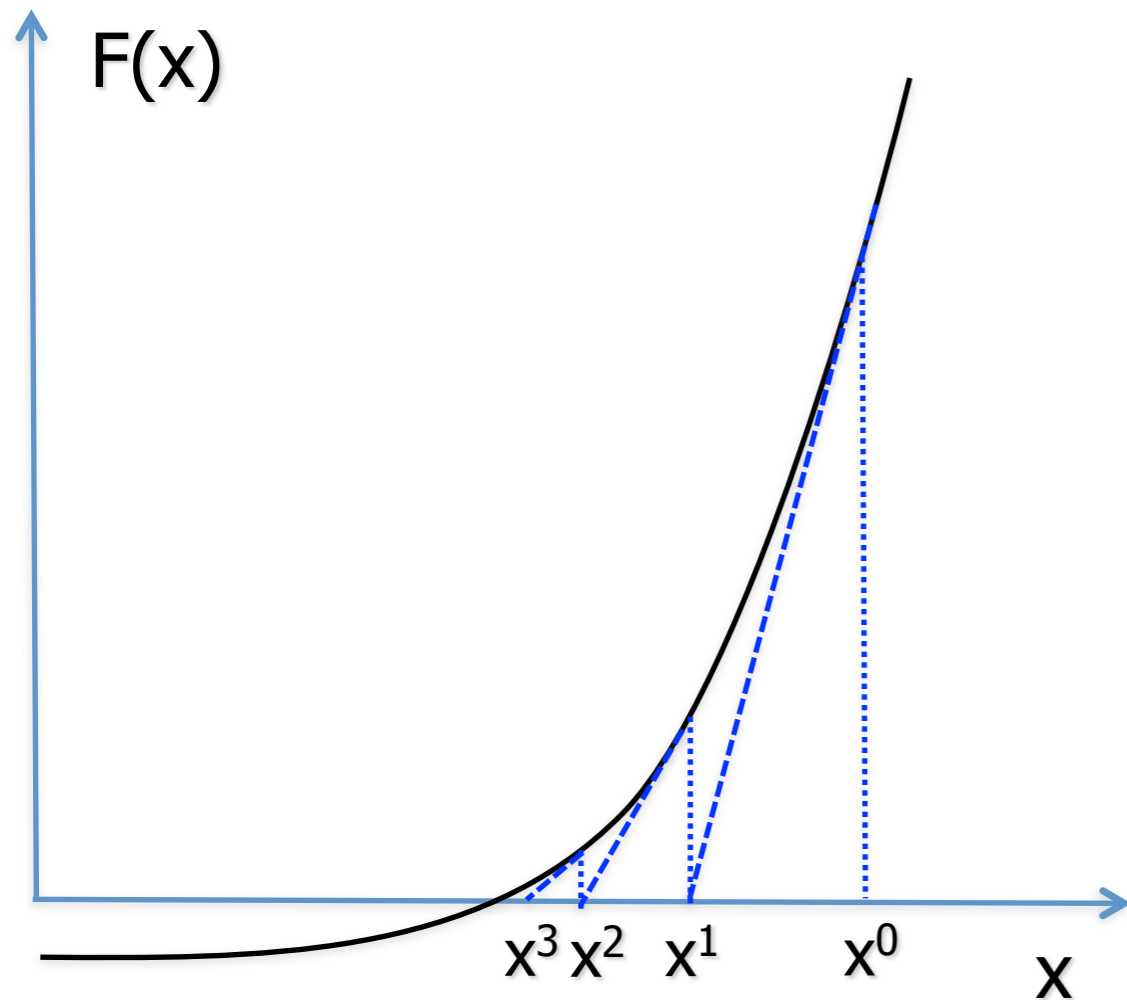
new to sea ice dynamics: the JFNK solver

$$[A(\mathbf{x}^{k-1}) \mathbf{x}^k - \mathbf{b}(\mathbf{x}^{k-1}) = \mathbf{F}(\mathbf{x}^k)]$$

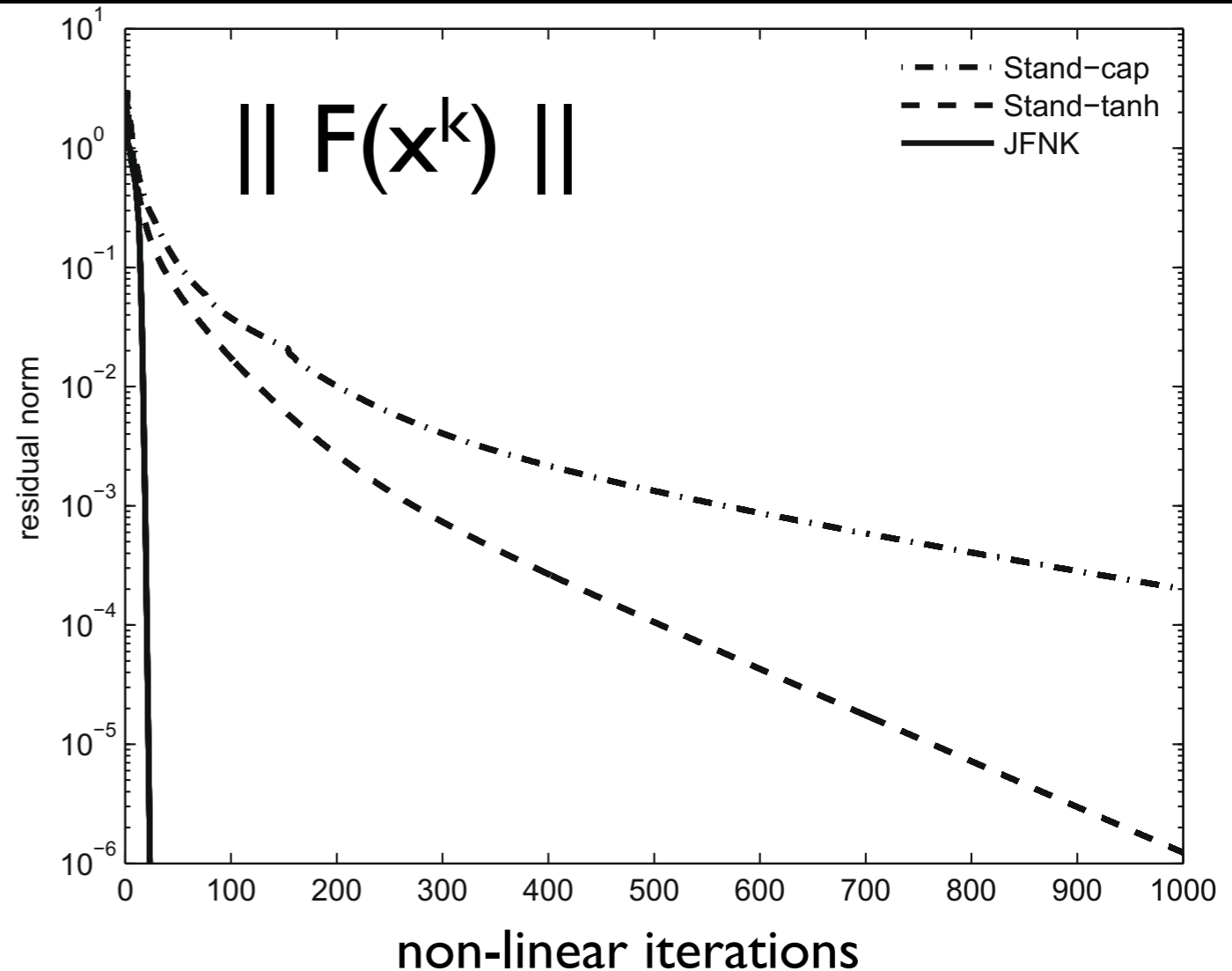
$$(0 \approx) \mathbf{F}(\mathbf{x}^{k-1} + \Delta \mathbf{x}^k) = \mathbf{F}(\mathbf{x}^{k-1}) + \mathbf{F}'(\mathbf{x}^{k-1}) \Delta \mathbf{x}^k, \quad \mathbf{F}' = \mathbf{J}(\text{acobian})$$

until $\|\mathbf{F}(\mathbf{x}^k)\| < \text{tol}$, solve $\mathbf{J}(\mathbf{x}^{k-1}) \Delta \mathbf{x}^k = -\mathbf{F}(\mathbf{x}^{k-1})$

and update $\mathbf{x}^k = \mathbf{x}^{k-1} + \Delta \mathbf{x}^k$



Newton method

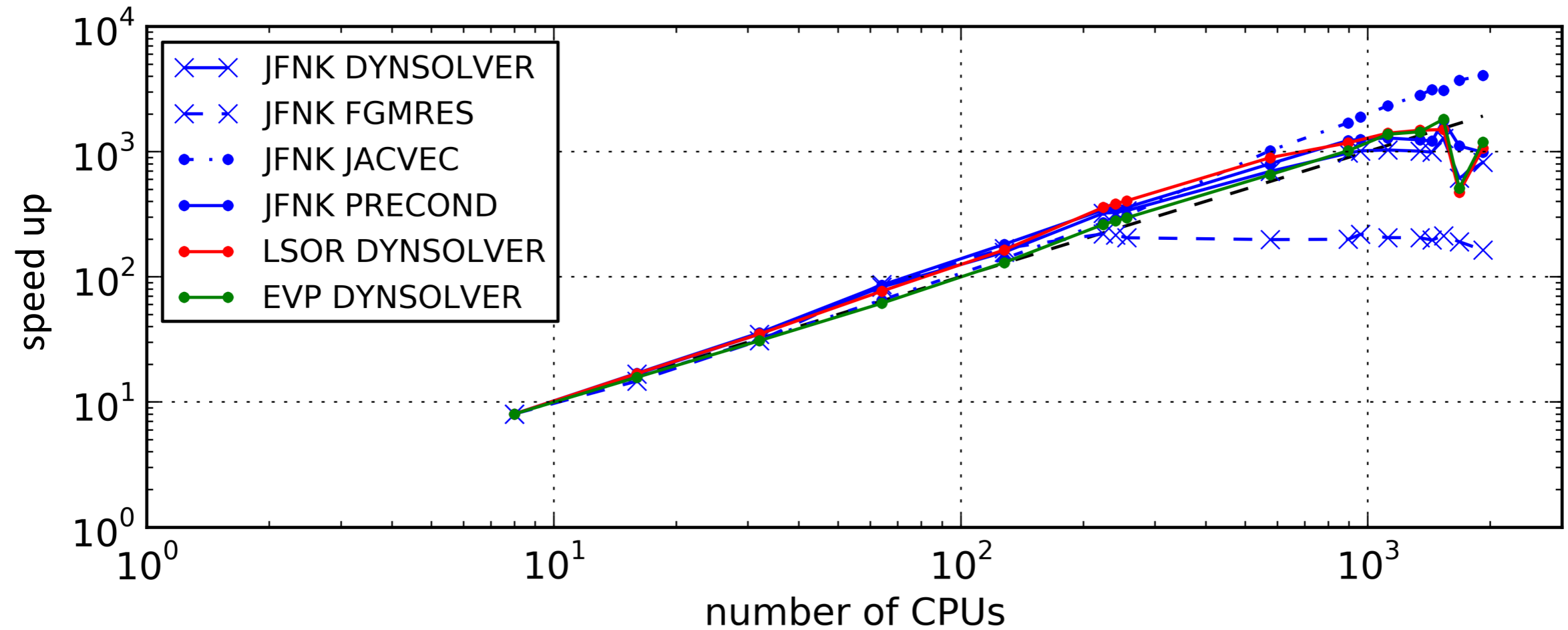
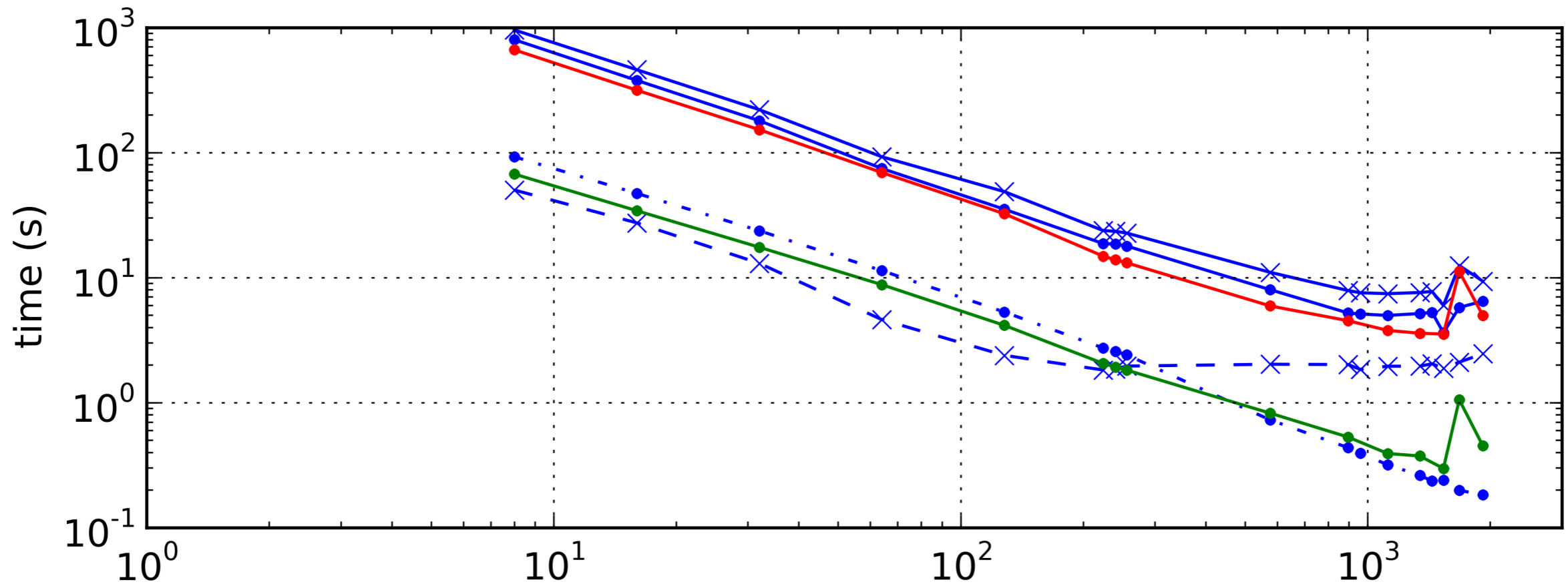


Lemieux et al. (2010)

JFNK in the MITgcm

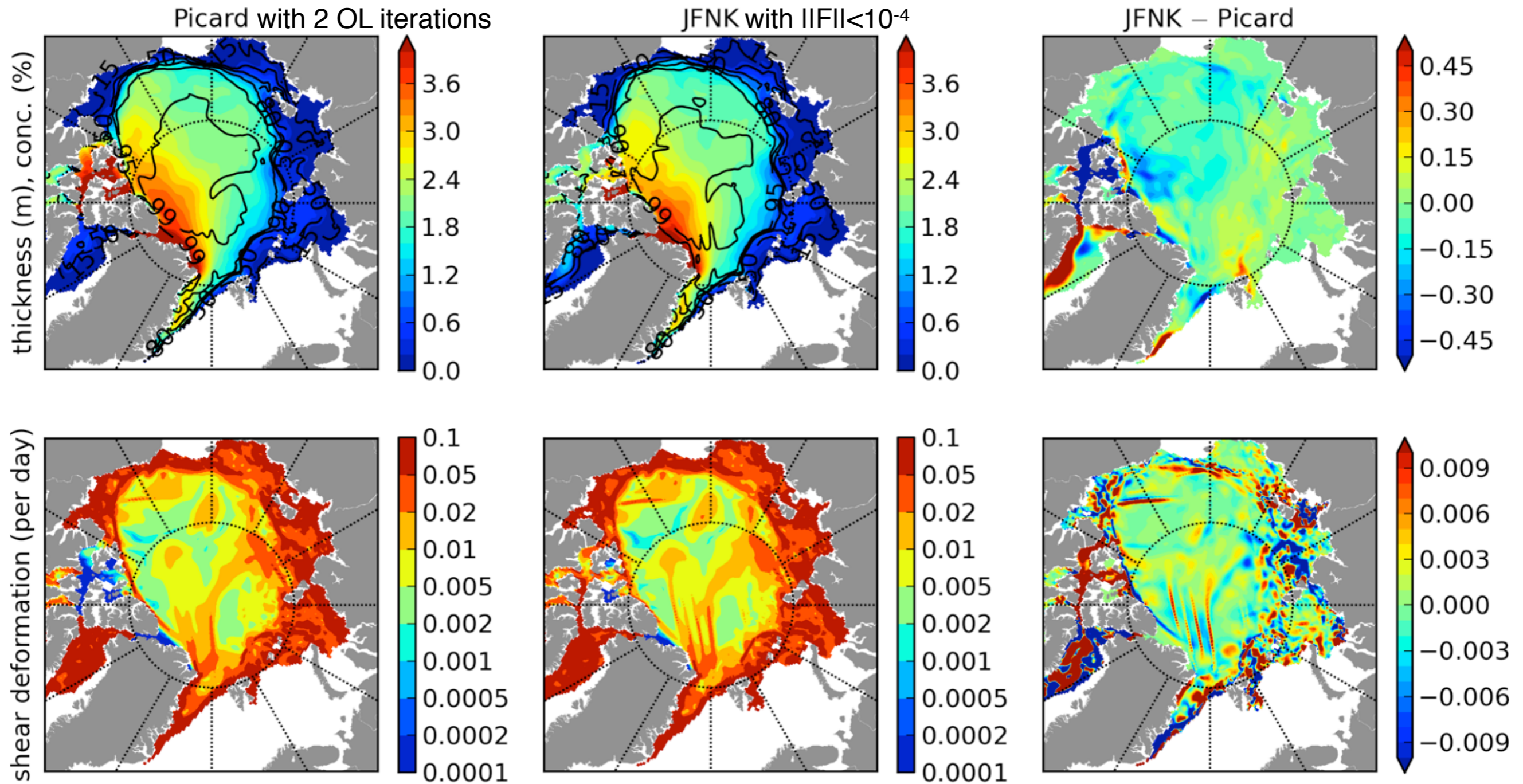
- FGMRES with preconditioner (LSR, Zhang and Hibler, 1997)
- exact Jacobian times vector possible by AD
- parallel code:
 - scalar products in FGMRES
 - restricted additive Schwarz (RAS) method in LSR
- vector code:
 - iterative preconditioner (LSR)
 - coloring (zebra) method

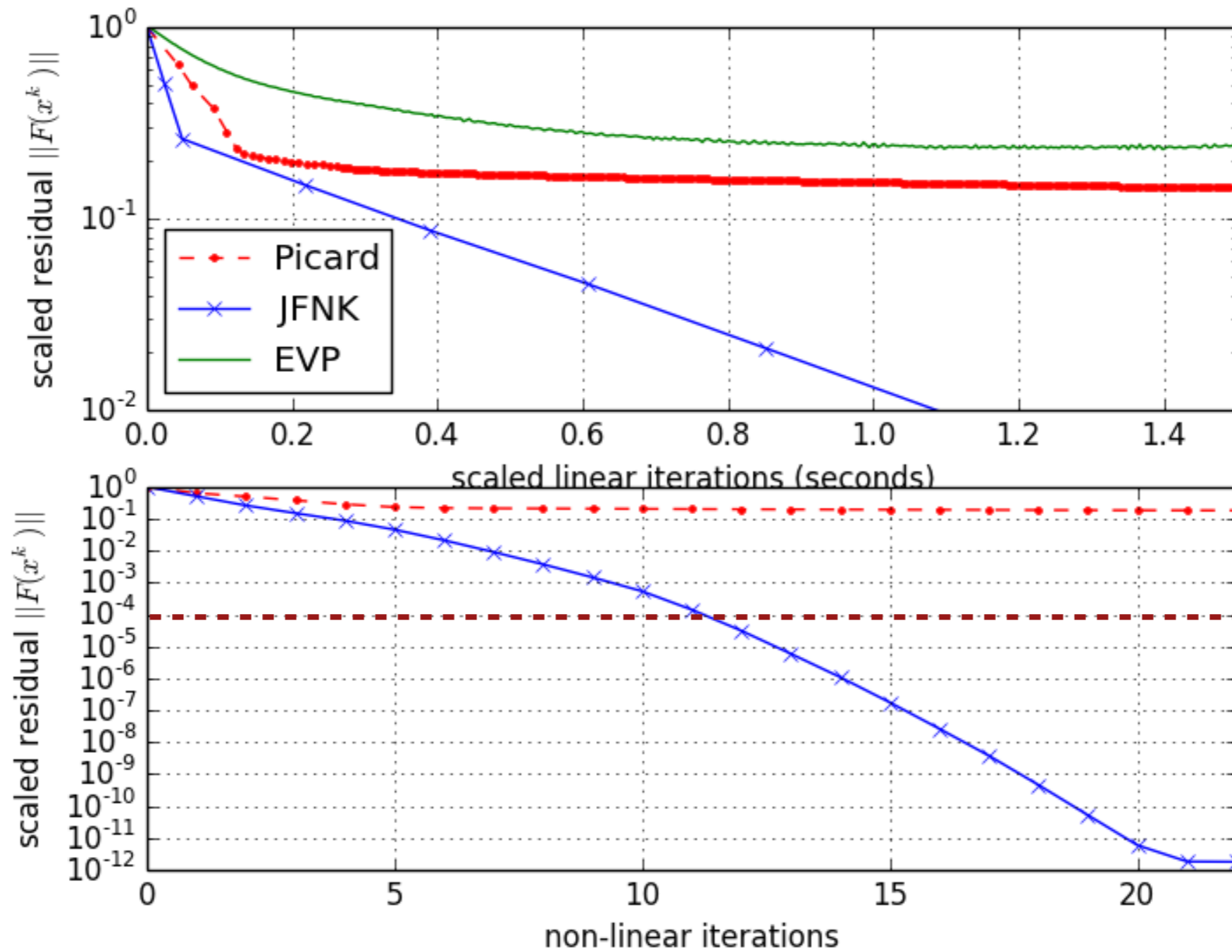
Parallel performance of solvers (4km Arctic configuration)



Does it matter?

after nearly 40 years of simulation: average of Oct, 1995





“Timing” of solvers
 (Is JFNK really faster?)

Accurate solvers affect

- ice thickness distribution
- linear kinematic features
- computer time!!!

Summary

- traditional Picard solver converges slowly
- EVP solver does not converge at all (to VP)
- JFNK-solver is efficient but expensive

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- traditional Picard solver converges slowly
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JFNK-solver is available for
large-scale problems