Size distribution and ionic composition of marine summer aerosol at the continental Antarctic site Kohnen

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Abstract. We measured aerosol size distributions and conducted bulk and size-segregated aerosol sampling during two summer campaigns in January 2015 and January 2016 at the continental Antarctic station Kohnen (Dronning Maud Land). Physical and chemical aerosol properties differ conspicuously during the episodic impact of a distinctive low-pressure system in 2015 (LPS15) compared to the prevailing clear sky conditions. The approximately 3-day LPS15 located in the eastern Weddell Sea was associated with the following: marine boundary layer air mass intrusion; enhanced condensation particle concentrations (1400 ± 700 cm⁻³ compared to 250 ± 120 cm⁻³ under clear sky conditions; mean ± SD); the occurrence of a new particle formation event exhibiting a continuous growth of particle diameters (Dₚ) from 12 to 43 nm over 44 h (growth rate 0.6 nm h⁻¹); peaking methane sulfonate (MS⁻), non-sea-salt sulfate (nss–SO₂⁻⁴), and Na⁺ concentrations (190 ng m⁻³ MS⁻, 137 ng m⁻³ nss–SO₂⁻⁴, and 53 ng m⁻³ Na⁺ compared to 24 ± 15, 107 ± 20, and 4.1 ± 2.2 ng m⁻³, respectively, during clear sky conditions); and finally an increased MS⁻/nss–SO₂⁻⁴ mass ratio βMS of 0.4 up to 2.3 (0.21 ± 0.1 under clear sky conditions) comparable to typical values found at coastal Antarctic sites. Throughout the observation period a larger part of MS⁻ could be found in super-micron aerosol compared to nss–SO₂⁻⁴, i.e., (10 ± 2) % by mass compared to (3.2 ± 2) %, respectively. On the whole, under clear sky conditions aged aerosol characterized by usually mono-modal size distributions around Dₚ = 60 nm was observed. Although our observations indicate that the sporadic impacts of coastal cyclones were associated with enhanced marine aerosol entry, aerosol deposition on-site during austral summer should be largely dominated by typical steady clear sky conditions.

1 Introduction

The impact of aerosols on global climate, which is in particular mediated by governing cloud droplet concentrations and hence cloud optical properties (Rosenfeld et al., 2014; Seinfeld et al., 2016), is of crucial importance but likewise notoriously charged with the largest uncertainties (Boucher et al., 2013; Seinfeld et al., 2016). In a seminal review, Carslaw et al. (2013) concluded that uncertainties in cloud radiative forcing are inter alia caused by uncertainties in natural emissions of dimethylsulfide (DMS) producing biogenic sulfur aerosol (i.e., MS⁻ and nss–SO₂⁻⁴) and sea spray. The radiation budget over the Southern Ocean is a particular challenge for current climate models, as they disturbingly underpredict the aerosol optical depth, pointing to a missing source of aerosols influencing cloud properties in this region (Bodas-Salcedo et al., 2014; Humphries et al., 2016). The competing role of biogenic sulfur (linked with phytoplankton presence) and otherwise sea salt aerosol (linked with stormy sea) in controlling cloudiness above the Southern Ocean is still up for debate (Meskhidze and Nenes, 2006; Korhonen et al., 2008; Quinn and Bates, 2011; Gras and Keywood, 2017).

Biogenic sulfur aerosol, i.e., secondary aerosol produced by the photooxidation of DMS and primary sea salt aerosol, by far dominate the aerosol mass over the Southern Ocean.
around Antarctica (Raes et al., 2000; Quinn and Bates, 2011). This original marine aerosol is partly transported to continental Antarctica and eventually deposited on the ice shield. Ideally, deposited aerosol constituents are archived in chronological order in firm (densified snow) and ice (Legrand and Mayewski, 1997). Therefore ice core records of biogenic sulfur and sea salt tracers potentially provide invaluable information on their (strictly speaking local) atmospheric budget, which is intrinsically tied to the Southern Ocean climate in the past (Wolff et al., 2006; Kaufmann et al., 2009; Mayewski et al., 2009; Abram et al., 2013). More specifically, investigations of sea salt and biogenic sulfuric acid from ice core archives is consequently desirable. Certainly, the knowledge of the source region and major transport processes, in view of the poor knowledge of aerosol source strength and atmospheric concentrations regarding the Southern Ocean realm, retrieving representative historic data for the South Pole (Arimoto et al., 2004) and the two deep ice core drilling sites of the EPICA project (Kohnen: Weller and Wagenbach, 2007, and Concordia: Preunkert et al., 2008, and Becagli et al., 2012). Based on an extended (2006–2015) record of aerosol, Legrand et al. (2017a) found that low values of $\beta_{\text{MS}}$ in midsummer at Concordia (0.05 ± 0.02) coincided with periods of high photochemical activity as indicated by the presence of locally photochemically produced O$_3$. This outcome strongly suggests the occurrence of an efficient chemical destruction of MS$^-$ over the Antarctic Plateau in midsummer. In addition, dedicated aerosol investigations within EPICA at Kohnen (Piel et al., 2006) have demonstrated the conspicuous impact of a cyclone on aerosol transport: in the aftermath of an intense low-pressure system over the eastern Weddell Sea region in combination with a blocking high-pressure ridge to the east in early January 2002 (Birnbaum et al., 2006), MS$^-$ and nss–SO$_4^{2-}$ aerosol entry showed an exceptional maximum (Piel et al., 2006). So far, a distinct impact of cyclones on aerosol concentrations has been frequently observed at coastal sites (e.g., Ito and Iwai, 1981; for Syowa and Pant et al., 2010, for Maitri), but only rarely on the Antarctic Plateau region (Hogan and Barnard, 1978, for the South Pole). In the most elaborate study about this topic, Pant et al. (2010) reported an almost constant background aerosol concentration under calm conditions, while during the passage of cyclones coarse-mode sea salt aerosol increased by an order of magnitude and occasionally new particle formation could be detected in the aftermath of the storms.

Up to now, all investigations at Kohnen were purely based on bulk aerosol sampling. The aim of our present study is focused on the variability of aerosol number concentration, aerosol size distribution, and size-segregated aerosol sampling to investigate the impact of different general weather situations on the physical and chemical properties of the aerosol for a site located on the Antarctic Plateau (Weller et al., 2017). This extended synoptic documentation of the physics and chemistry of Antarctic aerosol primarily concentrates on biogenic sulfur aerosol due to its distinct seasonal summer concentration peak caused by the seasonality of marine biogenic activity in the surrounding Southern Ocean (Weller and Wagenbach, 2007). Hence deposition in Dronning Maud Land should be virtually entirely governed by the atmospheric concentration maximum on-site, which was sufficiently covered by our dedicated observations.

2 Experimental techniques and data evaluation methods

2.1 Site description

During two summer seasons between 16 January and 3 February 2015 and 12 and 29 January 2016, we conducted aerosol size distribution measurements and size-segregated aerosol samplings at the continental Kohnen station (75°00′ S, 00°04′ E, altitude: 2892 m a.g.l.; http://www.awi.de/en/expedition/stations/kohnen-station.html; last access: 7 December 2017). This summer camp is the former deep ice core drilling location within the EPICA project in Dronning Maud Land about 550 km away from the ice shelf edge. Apart from the given internet link, another detailed description of the site comprising technical issues and a retrospect of the scientific activities can be found in Oerter et al. (2009). In order to minimize the impact of contamination from the permanently running diesel generator at the station, all experiments were installed inside (in situ measurements) or around (aerosol sampling) a bivouac hut located in the clean air sector about 250 m to the northeast of Kohnen (Fig. 1). The power supply (7 kW) was provided by cable from the main station. No fuel-driven generator was operated in the designated clean air sector (Fig. 1) and motor vehicle traffic was strictly prohibited. Contamination-free sampling was controlled by the permanently recorded wind velocity and direction. Contamination was indicated for each of the...
following criteria: wind direction within a 260–340° sector and/or wind velocity < 1.5 m s\(^{-1}\). In the case of contamination as given by these criteria, we interrupted the aerosol sampling experiments (low volume and impactor).

For convenience we will use throughout this work the notion “day of the year” (doy) instead of the calendar date and time is indicated in UTC, which is virtually identical to the solar time. All trace compound concentrations refer to standard volume at 273 K and 1013 hPa (STP).

2.2 Experimental setup

2.2.1 Particle concentration and size distribution

An overview of the experimental setup during both campaigns is given in Table 1, comprising the respective measuring periods of the different configurations and the relevant specifications of the deployed instruments. The size distribution of the sub-micrometer aerosol at Kohnen was determined by a scanning mobility particle sizer (SMPS; TSI classifier model 3080; Wang and Flagan, 1990). Setup and respective data evaluation methods were virtually identical to our installation run at Neumayer Station, as already described in detail elsewhere (Weller et al., 2015). Hence we simply highlight here the most relevant points. The SMPS was generally run in combination with the same condensation particle counter (TSI, model CPC 3775; 50 % cutoff diameter \(D_{p(50\%)}\) of 2.5 nm) controlled by the TSI software AIM (Aerosol Instrument Manager\(^{®}\) version 9.0). The classifier was operated alternately with a so-called nano-DMA (nano-differential mobility analyzer; TSI model 3085) and a DMA model 3081. Equipped with the nano-DMA 3085, we adjusted the aerosol and sheath flow of the classifier to achieve nominal aerosol size distribution measurements between 2 and 64 nm, while in use with the DMA 3081, flows were adjusted to cover the size range between 10 and 420 nm (in both cases with 64 channel resolution). Both DMA types were operated with a scanning time of 120 s (retrace time 15 s) and the average size distribution of four consecutive (multiple charge and diffusion corrected) spectra was finally evaluated, resulting in a temporal resolution of 600 s.

Particle size distributions were complemented with continuous condensation particle (CP) concentration measurements (TSI, model CPC 3775; \(D_{p(50\%)}\) of 4 nm in 2015 and CPC 3025A, \(D_{p(50\%)}\) of 3 nm in 2016). Below, we will use the terms nucleation mode for particles with \(D_{p}<25\) nm and Aitken mode for the size range between 25 and 100 nm (Dal Maso et al., 2005).

2.2.2 Low-volume and impactor aerosol sampling

The low-volume sampler (ISAP\(^{®}\) 1050; Schulze Automatisierungstechnik, Germany) and the five-stage low-pressure Berner-type impactor (GIV, model 80/0.05/2.88) were installed outdoors ca. 15 m northeast of the bivouac hut with ambient air inlets about 1.70 m above the ground. Low-volume sampling (1.66 m\(^3\) h\(^{-1}\) (STP), face velocity 0.33 m s\(^{-1}\)) was regularly conducted in 24 h time periods, while the impactor was run with a flow rate of 4.74 m\(^3\) h\(^{-1}\) (STP) in 2-day or 3-day intervals (Table 1; sample exchange usually in the early afternoon around 14:00). For bulk aerosol sampling with the low-volume sampler we used a Teflon and a nylon (Nylasorb) filter in series (Pall Corp.; 47 mm diameter and 1 μm pore size). The low-volume sampling procedure and filter handling was the same as in our Kohnen campaigns in 2001 and 2002, as described in Piel et al. (2006). The impactor was equipped with Tedlar foils (Hauke GmbH, Austria) and the aerodynamic cutoffs of the stages were corrected for ambient conditions according to the Stokes law (Table 1).

According to Piel et al. (2006) sample extracts were analyzed by using ion chromatography (IC) for methane sulfitonate (\(\text{CH}_3\text{SO}_3^-\), \(\text{MS}^-\)), \(\text{Cl}^-\), \(\text{Br}^-\), \(\text{NO}_3^-\), \(\text{SO}_4^{2-}\), \(\text{Na}^+\), \(\text{NH}_4^+\), \(\text{K}^+\), \(\text{Mg}^{2+}\), and \(\text{Ca}^{2+}\). The nss–\(\text{SO}_4^{2-}\) part was calculated using the sea salt tracer \(\text{Na}^+\) according to Piel et al. (2006). Unfortunately, due to unbearable nitrate blanks (caused by an unperceived and unauthorized use of nitric acid in the IC laboratory), all measured \(\text{NO}_3^-\) concentrations were treated as meaningless and discarded. In addition, \(\text{Br}^-\), \(\text{K}^+\), \(\text{Mg}^{2+}\),
Table 1. Survey of the experimental setup during both summer campaigns at Kohnen station.

<table>
<thead>
<tr>
<th>Measured parameter</th>
<th>Summer campaign 2015 (between 16 Jan and 3 Feb 2015)</th>
<th>Summer campaign 2016 (between 12 Jan and 29 Jan 2016)</th>
</tr>
</thead>
<tbody>
<tr>
<td>Particle number concentration</td>
<td>CNC 3775 (TSI), cutoff $D_{50%} = 4$ nm; 60 s resolution</td>
<td>CNC 3025A (TSI), cutoff $D_{50%} = 3$ nm; 60 s resolution</td>
</tr>
<tr>
<td>Particle size distribution</td>
<td>SMPS 3936/WCPC 3788 (TSI); 10 min res. 1. nano-DMA 3085 (TSI), range (nominal): 2 to 64 nm, measuring period: 16 Jan to 27 Jan 2. long DMA 3081 (TSI), range (nominal): 10 to 420 nm, measuring period: 27 Jan to 2 Feb</td>
<td>SMPS 3936/WCPC 3788 (TSI); 10 min res. 1. nano-DMA 3085 (TSI) Range (nominal): 2 to 64 nm Measuring period: 12 to 23 Jan 2. long DMA 3081 (TSI), range (nominal): 10 to 420 nm, measuring period: 23 to 28 Jan</td>
</tr>
<tr>
<td>Bulk aerosol sampling</td>
<td>Teflon–nylon filter combination, about 24 or 48 h sampling intervals</td>
<td>Teflon–nylon filter combination, about 24 h sampling interval</td>
</tr>
<tr>
<td>Size-segregated aerosol sampling</td>
<td>Five-stage Berner-type impactor: cutoff (µm): 3.5, 1.2, 0.4, 0.12, 0.042, about 48 or 72 h sampling intervals</td>
<td>Five-stage Berner-type impactor: cutoff (µm): 3.5, 1.2, 0.4, 0.12, 0.042, about 48 or 72 h sampling interval</td>
</tr>
<tr>
<td>Meteorology</td>
<td>Weather transmitter WXT520 (Vaisala), $U, D, T, P, RH$; 60 s resolution</td>
<td>Weather transmitter WXT520 (Vaisala), $U, D, T, P, RH$; 60 s resolution</td>
</tr>
</tbody>
</table>

and Ca$^{2+}$ concentrations were frequently close to the analytical detection limit (DL) of around 0.2 ng m$^{-3}$. Figure 2 summarizes the overall experimental uncertainty and ascertained DL, which are valid for the main ionic compounds to be discussed below. The circles are representative percentage error at different concentrations above the given DL ($\xi_i$) considering both the analytical (filter handling and IC) accuracy and the uncertainty of the sampling volume (<± 5 %). By using this procedure the error close to the DL was typically ±25 % for all ions. The DL was derived from dedicated procedure blanks (in total 10 blanks for low-volume sampling and 4 for each impactor stage) to be 3 times the standard deviation of these blanks. The blue line is an exponential fit through the values. Above a concentration of around 10 ng m$^{-3}$, the uncertainty was mainly governed by the concentration-independent volume ascertainment of around ±5 %.

The original concentration data from the different impactor stages were inverted according a procedure described in Winklmayr et al. (1990). In order to assess the validity of the inversion algorithm, we first compared the mass balance of the original with the inverted data. The discrepancy was typically within 1 to 2 % by mass, but never exceeds 5 %. Next we had to consider the influence of the analytical uncertainty of the ion concentrations on the inversion procedure. To this end we ran a Monte Carlo simulation, i.e., 1000 realizations of the inversion for a given compound were calculated, independently varying the concentrations of each impactor stage within 2 SD of the experimental error (we used individual concentration-dependent errors derived from the exponential fit in Fig. 2). From these 1000 realizations, the mean size distribution and the confidence intervals were determined. The result for the two most critical cases is shown in the Supplement (Fig. S1). For clarity, we refrain below from displaying confidence intervals in the presented impactor data.

Comparing ionic concentrations derived from low-volume and impactor sampling (including results from both campaigns covering the same sampling period, i.e., 11 out of 12 impactor samples), we found that concentration ratios from the impactor versus low-volume sampling were

![Figure 2. Relative uncertainty $\epsilon_{rel}(i)$ of the ion measurements derived from bulk (low volume) and size-segregated (Berner impactor) sampling as a function of $\xi_i$. The circles represent the percentage error at different concentrations $\xi_i$, i.e., the measured concentration of the ion $[x_i]$ minus the corresponding DL$_i$. The blue line is an exponential fit through these values. Above a concentration of around 10 ng m$^{-3}$, the uncertainty was mainly governed by the concentration-independent volume ascertainment of around 5 % (red horizontal line).](image)
0.64 ± 0.15 for MS$^-$, 0.79 ± 17 for nss–SO$_4^{2-}$, 1.1 ± 0.2 for Na$^+$, 0.96 ± 0.4 for Cl$^-$, and 0.72 ± 0.5 for NH$_4^+$.

The somewhat low sampling efficiency of the impactor for MS$^-$ may be due to partial reevaporation from the Tedlar foils (Teinilä et al., 2014). Finally, denuder experiments conducted by Piel et al. (2006) demonstrated that spurious post-sampling reactions like mobilization of volatile acids could generally be regarded as negligible for Teflon–nylon filter combinations.

### 2.2.3 Meteorological measurements and backward trajectories

The basic meteorological parameters wind speed ($U$) and wind direction ($D$), temperature ($T$), barometric pressure ($P$), and relative humidity (RH) were measured in 60 s resolution by a compact automatic weather station (Vaisala Weather Transmitter WXT520). The origin and pathway of the air masses advected to the measuring site was assessed by backward trajectories (HYSPLIT 4.0, Hybrid Single-Particle Lagrangian Integrated trajectory; http://www.arl.noaa.gov/documents/reports/hysplit_user_guide.pdf; last access: 7 December 2017) relying on GDAS meteorological data with a spatial resolution of $1^\circ \times 1^\circ$ (longitude × latitude grid). The accuracy of back trajectories is generally dependent on the availability of reliable meteorological data and their spatial coverage, which is of particular concern for remote regions like the Southern Ocean and Antarctica (see, e.g., the review by Stohl, 1998). As for coastal site Neumayer we already discussed this point (Weller et al., 2014, 2015), but for continental Antarctica, the persisting surface inversions and katabatic winds pose additional challenges to air mass trajectory analyses. Though a detailed accuracy assessment is beyond the scope this work, we tried to validate the consistency of the back trajectories used by (i) using the 3-D wind fields of the GDAS data and the isentropic approximation, which could be more reliable for regions with sparse meteorological input data (Harris et al., 2005); (ii) varying the starting height above the ground between 10 and 100 m; and (iii) calculating trajectory ensembles whose starting points were varied by ±1$^\circ$ longitude and ±1$^\circ$ latitude each, while for the height above the ground 0, 250, and 500 m were chosen.

The impact of these different initial conditions on our conclusions will be appraised in the Discussion section. Finally, weather charts generated by the Antarctic Mesoscale Prediction System (AMPS) were used to assess the general weather situation (Powers et al., 2003; http://www2.mmm.ucar.edu/rt/amps/information/amps_esg_data_info.html; last access: 7 December 2017).

### 3 Results

#### 3.1 Meteorological conditions

Figures 3 and 4 give a detailed synopsis of the measured meteorological parameters. Regarding local meteorology, bright weather conditions, i.e., clear sky conditions with a cloud amount typically well below 2 oktas, prevailed, except the period between 17 and 20 January 2015 (doy 17 at 12:00 to doy 20 at 18:00) when the impact of a cyclone (LPS15) with a center in the southeastern Weddell Sea reached the site (a corresponding weather chart is shown in Fig. S2). Note that the cloud amounts given in Figs. 3 and 4 were assessed by...
everyday casual visual inspection between 04:00 and 22:00 and should be rated as subjective but reasonable estimates. In the following we will especially focus on this episodic and prominent weather situation, which is preluded by a sharp pressure drop of 15 hPa and associated with increasing wind velocities to around 10 m s$^{-1}$, drifting snow, overcast sky, and some snow precipitation (which was notoriously difficult to distinguish from drifting snow). Corresponding HYSPLIT calculations predicted only marginal local precipitation around 0.2 mm water equivalent in the evening on 19 January 2015. A second prominent high wind event occurred on 15 January 2016 (LPS16; Figs. 4 and S3), again induced by an intense low-pressure system in the Weddell Sea region. In both cases the nearby coastal site Neumayer (70$^\circ$ 39′S, 8$^\circ$ 15′W) experienced severe blizzards with wind velocities up to 25 m s$^{-1}$. But in contrast to LPS15, the sky was then only temporarily overcast in the case of LPS16 at Kohnen and no precipitation has been realized.

Due to our limited meteorological database, a thorough characterization of the atmospheric boundary layer was not feasible. Though to this end it seems reasonable to access a detailed investigation of the summertime atmospheric boundary layer at Kohnen by Van As et al. (2006). Concerning clear sky conditions, our meteorological conditions ($U$, $D$, $T$, and $P$) were virtually comparable to a 4-day period described therein. Hence, according to Van As et al. (2006), we may assume the presence of a stable boundary layer overnight with strong temperature inversion (typically around 5 K m$^{-1}$ and a thickness up to 50 m), developing to a slightly convective layer (up to 50 m thick) around noon with near-neutral stability and capped by an inversion layer (Van As et al., 2006).

3.2 Particle concentrations and related size distributions

In Figs. 5 and 6, CP concentrations measured with both CPCs (CP$_{4\text{nm}}$ and CP$_{3\text{nm}}$ during 2015 and 2016, respectively) along with the ultrafine particle concentrations between 3 and 25 nm (UCP$_{3-25}$) and between 10 and 25 nm (UCP$_{10-25}$) are displayed. CP concentrations were strikingly enhanced during LPS15 ($1400 \pm 700 \text{ cm}^{-3}$ compared to $250 \pm 120 \text{ cm}^{-3}$ during clear sky conditions), though on the whole clearly lower compared to coastal sites like Neumayer (Weller et al., 2011). Both UCP time series were derived from the nanoDMA 3085 and DMA 3081 data, which are conflated as contour plots in Figs. 7 and S4–S6. The striking feature was the strong rise in CP, most notably UCP$_{3-25}$ concentrations, during the onset of LPS15. This period was characterized by elevated UCP$_{1-25}$ / CP$_{4\text{nm}}$ ratios of 0.75 ± 0.5. After a recurring weather situation, i.e., the remaining observation period...
in 2015 and throughout the campaign in 2016, including the stormy period LPS16, UCP / CP ratios were typically below 0.3 (more precisely: UCP$_{3–25}$ / CP$_{4 \text{ nm}}$ and UCP$_{3–25}$ / CP$_{3 \text{ nm}}$ stayed around 0.28 ± 0.2, while UCP$_{10–25}$ / CP$_{4 \text{ nm}}$ and UCP$_{10–25}$ / CP$_{3 \text{ nm}}$ were 0.18 ± 0.07 and 0.11 ± 0.1, respectively, Figs. 5 and 6). Table 2 summarizes particle concentrations and the ionic composition of the aerosol during the impact of LPS15, LPS16, and clear sky conditions. While during 2015 CPC and UCP concentrations were nearly an order of magnitude higher compared to clear sky conditions, the effect of LPS16 was not obvious.

Expectedly, the time series of the particle number size distributions (PNSDs) exhibited a corresponding distinct feature (Fig. 7): during LPS15, we observed according to Dal Maso et al. (2005) a pronounced class 1 or so-called “banana-type” NPF event. Particle growth started at doy 19 (02:50) at a modal maximum of 12 nm, reaching 43 nm at doy 20 (24:00; Fig. 7). While CP and UCP concentrations exhibited distinct breaks around noon at doy 19 (Fig. 5), steady particle growth was observed throughout. The initial nucleation particle formation rate (doy 19 between 02:50 and 17:20) in the size range 3 to 25 nm ($J_{3–25} = \Delta \text{UCP}_{3–25} / \Delta t$) was around 0.1 s$^{-1}$. The shape of the size distributions during this NPF event were “closed” (i.e., decreasing from the modal maxima towards the lowest size bin at 3 nm). From their temporal evolution we derived a continuous growth rate of 0.6 ± 0.08 nm h$^{-1}$ for the whole event between doy 19 (02:50) and doy 20 (24:00). In addition, a separate calculation for the first and second part of the event (boundary: noon on doy 19) resulted in virtually identical growth rates of 0.6 ± 0.1 nm h$^{-1}$ for both sections. A condensation sink of typically around $(2.0 \pm 0.2) \times 10^{-4}$ s$^{-1}$ was calculated according to Kulmala et al. (2001) from SMPS data when operated with DMA 3081, covering the extended size range between 10 and 420 nm.

In addition, we observed enhanced UCP concentrations between 10 and 25 nm during doy 18 in the morning hours, just before the actual NPF, and in the evening on doy 19 (numbered 2, 4, and 5 in Fig. 7). Apart from that, discernible natural nucleation bursts occurred around noon on doy 17 and 18 (numbered 1 and 3 in Fig. 7). All these transient UCP maxima did not show any detectable particle growth. The nucleation bursts were characterized by increasing particle concentrations from slightly above 5 nm downward towards the lower instrumental size limit (“open” distribution), indicating local nucleation. Table 3 provides a summary of the
Table 2. Particle number concentrations and ionic composition (mean ± SD) classified into different general weather conditions (impact of low-pressure system and clear sky conditions during both campaigns); note that the ionic composition during LPS16 was based on one sample, and hence no SD can be given. NA – not available.

<table>
<thead>
<tr>
<th>Measured parameter</th>
<th>Summer campaign 2015 LPS15 clear sky cond.</th>
<th>Summer campaign 2016 LPS16 clear sky cond.</th>
</tr>
</thead>
<tbody>
<tr>
<td>CPC_{4\text{nm}} (cm\textsuperscript{-3})</td>
<td>1400 ± 700</td>
<td>250 ± 120</td>
</tr>
<tr>
<td>CPC_{3\text{nm}} (cm\textsuperscript{-3})</td>
<td>NA</td>
<td>NA</td>
</tr>
<tr>
<td>UCP_{3–25\text{nm}} (cm\textsuperscript{-3})</td>
<td>880 ± 730</td>
<td>90 ± 100</td>
</tr>
<tr>
<td>UCP_{10–25\text{nm}} (cm\textsuperscript{-3})</td>
<td>NA</td>
<td>46 ± 24</td>
</tr>
<tr>
<td>MS\textsuperscript{−} (ng m\textsuperscript{-3})</td>
<td>111 ± 70</td>
<td>14.8 ± 3.7</td>
</tr>
<tr>
<td>nss–SO\textsubscript{4}\textsuperscript{2−} (ng m\textsuperscript{-3})</td>
<td>96 ± 35</td>
<td>81 ± 9</td>
</tr>
<tr>
<td>β_{MS} (MS\textsuperscript{−} / nss–SO\textsubscript{4}\textsuperscript{2−})</td>
<td>1.14 ± 0.8</td>
<td>0.18 ± 0.03</td>
</tr>
<tr>
<td>Cl\textsuperscript{−} (ng m\textsuperscript{-3})</td>
<td>15 ± 11</td>
<td>2.6 ± 2</td>
</tr>
<tr>
<td>Na\textsuperscript{+} (ng m\textsuperscript{-3})</td>
<td>33 ± 20</td>
<td>2.8 ± 1.4</td>
</tr>
<tr>
<td>Cl\textsuperscript{−} / Na\textsuperscript{+}</td>
<td>0.45 ± 0.3</td>
<td>1.24 ± 1.2</td>
</tr>
<tr>
<td>NH\textsubscript{4}\textsuperscript{+} (ng m\textsuperscript{-3})</td>
<td>4.0 ± 3.8</td>
<td>5.7 ± 2.4</td>
</tr>
</tbody>
</table>

3.3 Ionic composition of bulk and size-segregated aerosol

Though super-micrometer particles were not captured by the SMPS data, PNSD appeared clearly governed by submicrometer aerosol (Fig. 8), indicating a crucial role of nss–SO\textsubscript{4}\textsuperscript{2−} in the chemical composition of the aerosol. Indeed this was supported by the results from bulk and size-segregated aerosol samplings. Throughout both campaigns, about (75 ± 6)% of the aerosol mass consisted of biogenic sulfur aerosol (nss–SO\textsubscript{4}\textsuperscript{2−} and MS\textsuperscript{−}), leaving only about (9 ± 5)% for sea salt aerosol, while the highly variable NH\textsubscript{4}\textsuperscript{+} portion contributed to about (6 ± 4)% of the sea salt aerosol. The amount of sea salt aerosol was calculated from the measured Na\textsuperscript{+} concentrations referring to standard mean ocean water composition with an Na\textsuperscript{+} / sea salt mass ratio of 0.306 (Holland, 1993).

Note that due to analytical problems we assumed an NO\textsubscript{3}\textsuperscript{−} portion of 7.4% in the above mass balance estimate, according to results from three former summer campaigns at Kohnen (Piel et al., 2006). Ion balance considerations revealed that the sampled aerosol was constantly acidic, i.e., (0.7 ± 0.3) neq m\textsuperscript{-3} H\textsuperscript{+}, corresponding to (40 ± 15)% H\textsuperscript{+} in 2015 and (0.8 ± 0.6) neq m\textsuperscript{-3} H\textsuperscript{+} or (25 ± 16)% H\textsuperscript{+} equivalent contingent in 2016. Note that these figures are not a lower limit because we ignore here any HNO\textsubscript{3} part due to given analytical problems. Although Cl\textsuperscript{−} / Na\textsuperscript{+} mass ratios were highly variable (1.0 ± 0.9 in 2015 and 0.7 ± 0.3 in 2016), a significant Cl\textsuperscript{−} depletion relative to the seawater composition (Cl\textsuperscript{−} / Na\textsuperscript{+} = 1.8) was evident, indicating HCl mobilization by acids like HNO\textsubscript{3} or H\textsubscript{2}SO\textsubscript{4}.
Table 3. List of distinct enhanced UCP concentrations and natural nucleation burst events apart from the main NPF (event numbers refer to Fig. 7). The initial nucleation rate (mean ± SD) is specified together with the period of the UCP3–25 rise and the range of the particle diameter $D_p$.

<table>
<thead>
<tr>
<th>Event</th>
<th>Event start</th>
<th>Event end</th>
<th>$J_{3–25}$ (s$^{-1}$)</th>
<th>Period of UCP3–25 rise</th>
<th>$D_p$ range (nm)</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>17 Jan, 12:00</td>
<td>17 Jan, 15:00</td>
<td>0.05 ± 0.02</td>
<td>12:00 to 15:00</td>
<td>&lt;3.0 to 6.0</td>
</tr>
<tr>
<td>2</td>
<td>18 Jan, 00:00</td>
<td>18 Jan, 12:00</td>
<td>0.06 ± 0.02</td>
<td>00:00 to 02:00</td>
<td>8.0 to 25</td>
</tr>
<tr>
<td>3</td>
<td>18 Jan, 13:00</td>
<td>18 Jan, 14:00</td>
<td>0.15 ± 0.03</td>
<td>13:00 to 14:00</td>
<td>&lt;3.0 to 8.0</td>
</tr>
<tr>
<td>4</td>
<td>18 Jan, 18:00</td>
<td>18 Jan, 24:00</td>
<td>0.24 ± 0.04</td>
<td>18:00 to 21:00</td>
<td>6.0 to 40</td>
</tr>
<tr>
<td>5</td>
<td>19 Jan, 17:00</td>
<td>19 Jan, 24:00</td>
<td>0.11 ± 0.03</td>
<td>17:00 to 19:00</td>
<td>8.0 to 12</td>
</tr>
</tbody>
</table>

**Figure 9.** Time series of the measured MS$^-$ and nss–SO$_4^{2-}$ concentrations as well as the MS$^-$ / nss–SO$_4^{2-}$ mass ratio $\beta_{MS}$ from bulk aerosol (low volume) sampling during the campaign 2015. The period of LPS15 is shaded in yellow.

Figures 11 and 12 show the size-segregated composition of the aerosol derived from impactor samples taken 2015 and 2016. Note that in case of 2016 we could not assign particular impactor results to LPS16 due to the short duration of this event compared to the sampling period. In Fig. 11, percentage entries denote the portion of the corresponding ion mass in the size range >1 µm (super-micrometer or coarse mode) calculated from the inverted profiles. Generally consistent with comparable measurements conducted at Concordia (Legrand et al., 2017a), a significantly larger part of MS$^-$ compared to nss–SO$_4^{2-}$ resided in the coarse mode, resulting in higher $\beta_{MS}$ ratios there (range: 0.4 up to 2.3 during LPS15 when MSA reached 190 ng m$^{-3}$; Fig. 13). Concerning sea salt aerosol, nearly half of the Na$^+$ mass was present as sub-micrometer aerosol (Fig. 11). On the whole, ion balance considerations (again inevitably neglecting HNO$_3$) indicated acidic sub-micrometer aerosol, while in super-micrometer samples nss–SO$_4^{2-}$ and MS$^-$ were roughly counterbalanced by Na$^+$ and NH$_4^+$, which is in agreement with previous results from Concordia during summer (Becagli et al., 2012; Legrand et al., 2017a). However, ion balance results appeared not as clear-cut as found in the bulk low-volume Teflon–nylon filter combination samples: for sub-micrometer aerosol the mean excess anionic portion was merely 0.074 neq m$^{-3}$ in 2015 and 0.03 neq m$^{-3}$ in 2016, while for super-micrometer aerosol we found a mean excess cation part of 0.015 neq m$^{-3}$ (2015) and 0.084 neq m$^{-3}$ (2016). There are two main plausible reasons for this discrepancy: (i) a more pronounced loss of acidic gases from the impactor foils and (ii) an efficient sampling of gaseous HCl and HNO$_3$ on nylon filters (Piel, 2004; Piel et al., 2006). Again during LPS15, impactor results (covered mainly by the first and partly by the second impactor sampling period) exhibited some conspicuous features: particle mass size distributions (PMSDs) for MS$^-$ appeared somewhat broader and a considerably greater part of nss–SO$_4^{2-}$ and NH$_4^+$ resided in the super-micrometer mode (Fig. 11). Regarding the 2016 campaign, impactor results (Fig. 12) were apparently comparable to the 2015 campaign for prevailing clear sky conditions (covered by impactor samples 3–5; Fig. 11).
Figure 10. Time series of the measured \( MS^- \) and nss–\( SO_{4}^{2-} \) concentrations as well as the \( MS^- / nss-SO_{4}^{2-} \) mass ratio \( \beta_{MS} \) from bulk aerosol (low volume) sampling during the campaign 2016. The period of LPS16 is shaded in yellow.

4 Discussion

4.1 Case study LPS15: cyclone-induced marine air advection and NPF

4.1.1 Bulk and size-segregated chemical composition

In Dronning Maud Land cyclone-driven marine air intrusions are infrequent, sporadic events that are often associated with high precipitation rates (Birnbaum et al., 2006; Schlosser et al., 2010; Welker et al., 2014; Kurita et al., 2016). Such a cyclone-induced advection of marine boundary layer air masses towards the Antarctic Plateau had essentially coined the physical and chemical properties of the aerosol on-site, most noticeably by a maximum of biogenic sulfur concentrations and the occurrence of an NPF event. During a previous similar general weather situation at Kohnen (10 and 11 January 2002; LPS02), Piel et al. (2006) reported even strikingly higher nss–\( SO_{4}^{2-} \) (1100 ng m\(^{-3}\)) and \( MS^- \) (350 ng m\(^{-3}\)) maxima, though about 48 h after the passing of the frontal system and the transition from marine to continental air mass origin. In contrast to the LPS15, LPS02 came along with heavy snowfall (Birnbaum et al., 2006). Another peculiarity of our recent observations was a preceding and well-defined \( MS^- \) peak pursued by a distinct nss–\( SO_{4}^{2-} \) maximum (Fig. 9) compared to their simultaneous emergence in the aftermath of LPS02 described in Piel et al. (2006). Given that \( MS^- \) should have been primarily formed by heterogeneous liquid-phase chemistry prevalent in the marine boundary layer (Legrand et al., 2001; Bardouki et al., 2002; Hoffmann et al., 2016), the segregated \( MS^- \) peak indicated a temporary and efficient advection of such air masses. In addition, a striking peak of the sea salt tracer Na\(^+\) appeared along with the \( MS^- \) maximum (Supplement, Fig. S7), emphasizing the impact of marine boundary layer air masses during this part of LPS15. Note that Na\(^+\) remained at typical mean concentrations around 7 ng m\(^{-3}\) during the former biogenic sulfur maximum between 14 and 15 January 2002 (Piel, 2004).

Figure 11. Results from size-segregated (Berner impactor) sampling during the campaign 2015. Bold reddish lines are the mass size distributions during LPS15. Percentage entries in the legend denote the portion of the corresponding ion mass in the size range > 1 \( \mu \)m (super-micrometer or coarse mode) calculated from the inverted profiles. The bold gray line is the median mass size distribution of the corresponding ion.

The 5-day back trajectories confirmed these conclusions (Fig. 14; starting point 100 m above Kohnen): air masses during LPS15 were generally marine 2 to 3 days before arrival at Kohnen. Trajectories representing doy 19 (NPF event; reddish and yellow traces in Fig. 14a) spent several hours within the marine boundary layer close to the East Antarctic coast before arrival at Kohnen, largely following the contour lines of the local topography. All trajectories started under cyclonic curvature, finally approaching Kohnen in an anticyclonic bow from northerly directions (Fig. 14). In order to estimate the reliability of this trajectory-based finding, we repeated the calculations with an initial height of 10 m above Kohnen, calculated trajectory ensembles, and made an extra attempt using the isentropic approach instead of the 3-D wind field from GDAS data (Figs. S8–S10). Though appreciable differences regarding the geographic location of the corresponding source regions and trajectory course were obvious, in the end the basic aforementioned implications appeared consistent. In contrast to this finding, air masses originated and stayed within continental Antarctica during the biogenic sulfur peaks observed during 14 and 15 January 2002 (Piel et al., 2006), which could be confirmed by a reanalysis with HYSPLIT trajectories based on NCEP meteorological data (Fig. S11). Similar air mass trajectories (not shown) were ob-
Results from size-segregated (Berner impactor) sampling during the campaign 2016 with the following mass portion of the respective ions in the super-micrometer range (>1 µm): MS\(^-\) 14 % (a), nss–SO\(_4\)\(^{2-}\) 3 % (b), Na\(^+\) 46 % (c), and NH\(_3\)\(^+\) 15 % (d). The bold gray line is the median mass size distribution of the corresponding ion.

Figure 12. Results from size-segregated (Berner impactor) sampling during the campaign 2016 with the following mass portion of the respective ions in the super-micrometer range (>1 µm): MS\(^-\) 14 % (a), nss–SO\(_4\)\(^{2-}\) 3 % (b), Na\(^+\) 46 % (c), and NH\(_3\)\(^+\) 15 % (d). The bold gray line is the median mass size distribution of the corresponding ion.

Figure 13. Mean size-segregated results for the MS\(^-\) / nss–SO\(_4\)\(^{2-}\) mass ratio \(\beta_{\text{MS}}\) determined for each of the five impactor stages for both seasons and for the event LPS15.

...served during and in the aftermath of a short stormy period LPS16, which was again characterized by enhanced Na\(^+\) and biogenic sulfur loadings (Figs. 10 and S7b).

Regarding the chemical composition of the aerosol during the final stage of LPS15, subsequent increasing nss–SO\(_4\)\(^{2-}\) along with declining MS\(^-\) concentrations indicated a minor importance of liquid-phase chemistry (Hoffmann et al., 2016). We may speculate that intrusions of marine boundary layer into the so-called buffer layer were responsible for the efficient advection of gaseous DMS photooxidation products like SO\(_2\) and DMSO (Davis et al., 1998; Russel et al., 1998). Russel et al. (1998) assumed that the buffer layer typically extends from the turbulent marine boundary layer (400 to 700 m) up to a capping inversion (1400 to 1900 m). While transported to continental Antarctica, gas-phase photooxidation processes should have dominated, leading finally to a preferred formation of H\(_2\)SO\(_4\) at the expense of RSA (Preunkert et al., 2008). This plausible but subtle transport route could not be unequivocally deduced from respective backward trajectory analyses because in the case at hand the presence and the extent of a buffer layer could not be ascertained from available meteorological data.

Another conspicuous point was the comparatively large particle concentration and ionic composition maxima at Neumayer. This coastal site was governed by the same low-pressure system (LPS15) provoking a blizzard there around 19 January 2015. From there, daily low-volume sampling (Teflon–nylon filter combination), CP concentration (measured with a CPC 3022A, TSI; \(D_p(50\%) = 7\) nm), and meteorological data were available. Again, biogenic sulfur (particularly MS\(^-\)), Na\(^+\), and CP concentrations showed distinctive maxima though about 1 day after the LPS15 (Fig. S12), similar to the previous situation described by Piel et al. (2006). Consequently we can assume that this characteristic weather situation transported marine aerosol throughout DML. But in contrast to Kohnen the impact of LPS15 on the delayed particle concentration and ionic composition maxima at Neu-
mayer was less pronounced. Pant et al. (2010) concluded from previous particle concentration and size distribution data measured at coastal Maitri that during the impact of cyclones, coarse-mode sea salt aerosol increased by an order of magnitude compared to calm weather conditions, which is similar to our results during LPS15 (Table 2).

4.1.2 NPF, particle growth, and size-segregated chemical composition

Obviously the most striking feature during LPS15 was a distinctive NPF event. The closed shape of PNSD, starting with a modal maximum not less than 12 nm, implied that the actual particle nucleation event should have occurred upwind of Kohnen. Assuming a constant growth rate of 0.6 nm h\(^{-1}\) throughout advection to the measuring site, particle nucleation thus happened about 20 h before, i.e., according to backward trajectories roughly 700 km away from Kohnen at 73° S, 19° E and 160 m above the ground (all virtually independent from the choice of the starting height). Note that the observed particle growth was confined to the nucleation mode and consequently did not reach a size range potentially relevant for acting as CCN. As described in Weller et al. (2015), NPF events at Neumayer generally showed almost comparable growth rates. Therein we used a simple estimate based on Nieminen et al. (2010) and Yli-Juuti et al. (2011) that demonstrated the need for condensable vapors other than sulfuric acid to sustain the observed particle growth (Weller et al., 2015). We employed the same approach in connection with the NPF event at Kohnen (i.e., Eq. 3 in Weller et al., 2015, using \( T = 250 \) K), resulting in \( 1.4 \times 10^{7} \) molec cm\(^{-3}\) gaseous H\(_2\)SO\(_4\) necessary for the determined growth rate of 0.6 nm h\(^{-1}\). Again, this appeared at least an order of magnitude too high compared to the respective values published (e.g., Mauldin III et al., 2004), emphasizing the importance of other yet unknown condensable vapors, most probably low-volatility organic compounds (Metzger et al., 2010; Tröstl et al., 2016). Interestingly, Kyrö et al. (2013) identified biogenic emissions by nearby melting ponds as a potential source for condensable vapor. The surroundings of Kohnen, however, are completely ice covered throughout as typical for the Antarctic Plateau region. The nearest rocky outcrops are more than 200 km away.

In this context it is interesting to compare our results with previous PNSD measurements conducted at the South Pole during two summer campaigns (Park et al., 2004) and year-round observations at Concordia (Järvinen et al., 2013). Results from the South Pole showed similar PNSDs with mean a \( D_{p} \) between 60 and 98 nm (Park et al., 2004, Table 4 therein), while PNSDs at Concordia were appreciably lower with \( D_{p} = 39 \) nm (Järvinen et al., 2013). Järvinen et al. (2013) could detect NPF throughout the year under condensation sink values of around \( 1.8 \times 10^{-4} \) s\(^{-1}\), comparable to the condensation sink range found at Kohnen. At Concordia, NPF events with a determinable growth rate were not connected with low-pressure systems but essentially observed under air mass advection from the upper atmosphere (Järvinen et al., 2013). This fact indicates that also at Kohnen, NPF may potentially occur during clear sky conditions, but could not be detected within our admittedly limited observation period. Furthermore, the growth rates determined so far at coastal sites (Virkkula et al., 2007; Asmi et al., 2010; Kyrö et al., 2013; Weller et al., 2015) appeared comparable to GR reported from continental Antarctica (Järvinen et al., 2013; Chen et al., 2017) and were within a similar broad range between 0.2 and up to 8.8 nm h\(^{-1}\).

Pant et al. (2010) presented a detailed analysis about the impact of passing cyclones on particle size distributions at coastal Maitri. The authors observed bimodal PNSD with a coarse-mode maxima around 2 µm and a broad Aitken mode between 0.04 and 0.1 µm when a storm approached the site. Occasionally NPF occurred just after the passage of a cyclone associated with particle growth rates between 0.2 and 0.6 nm h\(^{-1}\). From meteorological data the authors conclude that the observed NPF events were linked with the mixing of marine and continental air during the subsidence of free tropospheric air after the storm (Pant et al., 2010).

Figure 14. The 5-day backward trajectories during the NPF event calculated with a trajectory starting height of 100 m above Kohnen at the points in time given in the legend (a). Below, the travel height above the ground (local topography) is illustrated in a color-coded scale (point interval 1 h) (b).
4.2 The standard case: clear sky conditions and aged aerosol

On the one hand, LPS15 was attended with an expeditious advection of marine boundary layer air, leading to an enhanced entry of marine aerosol and secondary aerosol precursors, highly variable PNSD, and an NPF event. On the other hand, clear sky conditions largely prevailed during our recent campaign, but also during previous summer campaigns (Piel et al., 2006). The present observations at Kohnen showed that throughout this characteristic synoptic situation, appreciably lower particle number concentrations restricted within the accumulation mode were typical. Concerning the ionic composition, MS$^-$ and Na$^+$ concentrations were considerably lower compared to their maxima accompanied with the impact of LPS15 (Figs. 9 and S7a) in contrast to virtually constant, though highly variable mean NH$_4^+$ concentrations ($5.2 \pm 2.6$ ng m$^{-3}$ in 2015 and $17.9 \pm 6$ ng m$^{-3}$ in 2016). We have no explanation at hand for the approximately threefold higher mean NH$_4^+$ concentrations in 2016, but the results presented in Piel et al. (2006) also showed large annual fluctuations. We finally calculated a mass ratio of NH$_4^+$ to nss–SO$_4^{2-}$ of 0.06 for 2015 and 0.15 for 2016. At least the former value was in the range of what was concurrently observed at Neumayer in January 2015. Most probably, atmospheric NH$_3$ / NH$_4^+$ ratios were correspondingly variable depending on the availability of acidic trace gases (note that NH$_3$ could not be captured by the sampling methods used).

Turning towards biogenic sulfur aerosol, $\beta_{MS}$ ratios decreased to values around 0.2 under clear sky conditions, as is typical for continental Antarctica (Piel et al., 2006; Weller and Wagenbach, 2007; Preunkert et al., 2008; Becagli et al., 2012), while during LPS15, $\beta_{MS}$ was more comparable to coastal sites like Neumayer and Dumont D’Urville (Legrand and Pasteur, 1998).

The air mass history during clear sky conditions was assessed by composite backward trajectory calculations and is summarized in Fig. 15. Even 10 days before arrival at Kohnen with a characteristic anticyclonic curvature, trajectory origins remained principally inside the Antarctic continent and thus remote from marine source regions. Varying initial start height did not essentially change this general feature, but employing the isentropic instead of the 3-D approach showed an increased relevance of marine source regions (see Fig. S13). Upon extending the trajectory travel time to 20 days, the origin of the air masses eventually became marine (Fig. 16), covering a large part of the Southern Ocean, except the western part of the Weddell Sea and the Bellingshausen Sea. The relevance of air mass transport via the free troposphere was difficult to assess, mainly due to the generally highly variable and poorly characterized depth of the marine boundary layer and especially the vertical extent of the atmospheric boundary layer over Antarctica. According to Russell et al. (1998), the border between the marine boundary layer and the free troposphere typically varies between 1400 and 1900 m, while on the Antarctic Plateau only a shallow atmospheric boundary layer of no more than a few hundred meters is typical (Van As et al., 2006). Inspection of the calculated trajectories revealed that they typically originated within the marine boundary layer and essentially below 1500 m and 500 m above the ground across continental Antarctica.

Surprisingly, while the amount of sea salt aerosol was highly variable, the super-micrometer-mode fraction of sea salt aerosol remained constant at around 50% during rapid and efficient marine boundary layer air mass advection under LPS15 and during long-range transport under clear sky conditions (Figs. 11 and 12). Since our sea salt mass size distribution also appeared similar to that typically observed at Concordia (Jourdain et al., 2008; Legrand et al., 2017b),
we infer that the transport of coarse-mode sea salt particles to continental Antarctica was generally inefficient, regardless of the general weather situation and transport time. Finally, the fact that observed $\beta_{MS}$ ratios were in general higher in the coarse mode (supported by measurements in coastal Antarctica by Teinilä et al., 2000, and Rankin and Wolff, 2003) implied a preferential loss of $MS^-$ during transport towards the Antarctic Plateau region. In contrast, recent results from Concordia give strong evidence for a preferential photochemical depletion of MSA in the atmosphere above the plateau during austral summer (Legrand et al., 2017a) but less indication for fractional loss en route. Nevertheless, both processes of fractionation and photochemistry could be potential explanations for typically lower $\beta_{MS}$ ratios commonly found in continental Antarctica compared to the coastal regions (Piel et al., 2006; Weller and Wagenbach, 2007; Preunkert et al., 2008).

5 Conclusions

We measured aerosol size distributions and conducted bulk and size-segregated aerosol sampling during two summer campaigns at the continental Antarctic station Kohnen. This extended approach allowed for a detailed synopsis of the physical and chemical properties of summer aerosol in this region. For the first time, the impact of passing cyclones on aerosol advection into the Antarctic Plateau region was examined. Based on these admittedly still limited investigations, we may conclude that during austral summer, the transport of marine aerosol to Kohnen in particular and to continental DML in general was mediated by two different synoptic situations: (i) the impact of low-pressure systems in the western part of the South Atlantic associated with temporarily exceptional marine aerosol concentrations and (ii) persistent long-range transport providing a background aerosol level during clear sky conditions over DML. In the present study, a distinct low-pressure event (LPS15) was additionally associated with NPF. Under prevailing clear sky conditions, on the other hand, aged aerosol and less aerosol (by mass and number concentration) entered DML in air masses that were typically continental for about 10 days before. We tentatively infer that our recent observation, i.e., NPF and peaking marine aerosol concentrations during LPS15, could be of fortuitous occurrence since unexpectedly just the trace compounds seemed to hardly be depleted by precipitation. In contrast, during the mentioned blizzard in 2002 (LPS02) reported by Piel et al. (2006), biogenic sulfur concentrations stayed first quite low but peaked at about 48 h in the aftermath of the storm. Though we refer to two recent and three previous studies, just three LPS and one pronounced NPF event occurred that could be analyzed in detail, emphasizing their sporadic nature. Hence, a worthwhile confirmation of our conclusions would clearly require similar investigations at this site. Such effort is important to better understand the role of biogenic aerosol in general and in particular the impact of NPF events on regional climate forcing.

Though the efficient transport of biogenic sulfur (and also sea salt) aerosol to continental DML may be associated with cyclonic activity in the South Atlantic, in the long run the crucial transport pathway of marine aerosol during austral summer should be long-range transport under typical clear sky conditions. In particular for biogenic sulfur, showing a pronounced summer maximum (Weller and Wagenbach, 2007), we suppose that transport to DML, deposition, and final storage in firn and glacial ice will be dominated by prevailing clear sky conditions. Thus dry deposition, but to an only minor extent wet deposition (partly associated with clear sky precipitation), would be decisive. Consequently future research activities should also envisage assessing dry deposition velocities at this site, e.g., by gradient and/or eddy correlation studies (Grönlund et al., 2002; Contini et al., 2010). On the other hand, retrieving meaningful historic aerosol concentrations from ice core archives also needs a thorough consideration of snow accumulation since snow accumulation co-determines trace compound concentrations in firn and ice (Fischer et al., 1998), which is evidently governed by the infrequent impact of low-pressure systems (Birnbaum et al., 2006; Schlosser et al., 2010; Welker et al., 2014; Kurita et al., 2016). Finally, trajectory analyses indicated that a large part of the Southern Ocean should be considered as a potential source region representative of aerosol deposition in continental DML in contrast to coastal Neumayer where the dominance of the South Atlantic was evident (Minikin et al., 1998).

Data availability. Data from both campaigns reported here are available at https://doi.org/10.1594/PANGAEA.882375 (Weller et al., 2017) for scientific purposes. In this case, we expect collaboration with Rolf Weller (contact: rolf.weller@awi.de).

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Competing interests. The authors declare that they have no conflict of interest.

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