

1 Landfast ice affects the stability of the Arctic
2 halocline: evidence from a numerical model

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Abstract.

Landfast ice covers large surface areas of the winter Siberian Seas. The im-
mobile landfast ice cover inhibits divergent and convergent motion, hence
dynamical sea ice growth and re-distribution, decouples winter river plumes
in coastal seas from the atmosphere and positions polynyas at the landfast
ice edge offshore. In spite of the potentially large effects, state-of-the-art nu-
merical models usually do not represent landfast ice in its correct extent. A
simple parametrization of landfast ice based on bathymetry and internal sea
ice strength is introduced its effects on the Arctic Ocean are demonstrated.
The simulations suggest that the Siberian landfast ice impacts the Arctic halo-
cline stability through enhanced brine production in polynyas located closer
to the shelf break and by re-directing river water to the Canadian Basin. These
processes strengthen the halocline in the Canadian Basin, but erode its sta-
bility in the Makarov and Eurasian Basin.

1. Introduction

17 One of the dominant characteristics of the winter Arctic shelf seas is landfast ice (also
18 land-fast, fast or shore-fast ice), sea ice that is immobile and mechanically fastened to
19 the coast or to the sea floor. As there is no compressive deformation, landfast ice grows
20 only thermodynamically and rarely exceeds thicknesses of 1.5 m [Romanov, 2004]. It
21 can extend a few (Beaufort Sea, Chukchi Sea, Western Laptev Sea) to several hundred
22 kilometers from the coast into the ocean (Kara Sea, Eastern Laptev Sea, East Siberian
23 Sea, see Fig. 1). The mechanisms that determine the landfast ice formation, extent and
24 decay are not fully understood. To complicate things further, these mechanisms differ
25 regionally. In the Chukchi and Beaufort Sea the landfast ice edge is found at relatively
26 shallow depth of 18 m [Mahoney *et al.*, 2007], while on the Eurasian shelf this depth is
27 between 25 and 30 m [Dmitrenko *et al.*, 2005a; Proshutinsky *et al.*, 2007]. In the Chukchi
28 and Beaufort Sea the landfast ice is immobilized behind a row of bottom reaching pressure
29 ridges [Mahoney *et al.*, 2007]. In the Kara Sea the landfast ice is formed behind a row
30 of small islands parallel to the coast [Divine *et al.*, 2004]. At the landfast ice edge in
31 the Laptev and East Siberian Seas grounded pressure ridges are rare [Reimnitz *et al.*,
32 1994; Eicken *et al.*, 2005]. Proshutinsky *et al.* [2007] proposed that the landfast ice edge
33 occurs where the warm intermediate Atlantic water reaches the surface after upwelling
34 at the shelfbreak. König Beatty and Holland [2010] attributed the landfast ice extent to
35 the mechanical properties of the sea ice. Strong freshwater and brackish sea ice [Dethleff
36 *et al.*, 1993; Eicken *et al.*, 2005] formed in low salinity shelf seas with high river water
37 content might ground at wide and shallow sand banks [Dethleff *et al.*, 1993; Reimnitz

38 *et al.*, 1994] and from there extend with long tongues toward the landfast ice edge. Tides
39 in the region are very weak and the tidal amplitudes of up to 10 cm in the Southeastern
40 Laptev Sea *Fofonova et al.* [2014] might not be large enough to deform the extensive
41 landfast ice cover.

42 Especially in the Siberian Seas, where landfast ice has the largest extent, it is thought
43 to have three roles:

44 1. Landfast ice limits the sea ice thickness by preventing sea ice compression (e.g.
45 pressure ridges) in convergent motion of sea ice. By the same token it lowers the sea ice
46 production by preventing sea ice divergence and lead formation. If these processes are not
47 adequately represented in a numerical model, simulated thickness fields are not realistic
48 [*Johnson et al.*, 2012].

49 2. An immobile lid of the landfast ice effectively decouples the inner shelf from the
50 atmosphere and affects the river water distribution. The seasonal runoff of the large
51 Siberian rivers (Ob, Lena, Yenisei) determines the water masses of the shallow Siberian
52 shelf. In a summer with strong onshore winds large parts of the summer maximal discharge
53 can be held back on the inner shelf until winter [*Dmitrenko et al.*, 2005b]. The large
54 amount of the fresh river water in the Arctic surface layer enhances the ocean stratification.

55 3. The extent of landfast ice determines the areas of low sea ice concentration—so called
56 flaw polynyas (also flaw lead polynyas or flaw leads)—along the landfast ice edge during
57 offshore wind conditions. The water in these polynyas is more saline than directly at the
58 coast, where ocean salinities remain low due to fresh river water inflow and consequently
59 more brine is rejected during sea ice production. The brine formed in the winter shelf
60 seas maintain, along with the cool water formed by winter convection north of the Barents

61 Sea [*Rudels et al.*, 1996], the Arctic halocline [*Aagaard et al.*, 1981; *Martin and Cavalieri*,
62 1989; *Cavalieri and Martin*, 1994; *Winsor and Björk*, 2000]. The cold and saline water
63 flows off the shelves and sinks along the shelf break where it feeds into the halocline layer
64 (Fig. 2), which decouples the cold Arctic surface layer from the warm intermediate-depth
65 Atlantic layer. The surface layer stays cold with temperatures close to the freezing point
66 due to a seasonal sea ice melting - freezing cycle and the layer is fresh due to a strong
67 river runoff. In contrast, the ocean temperature at mid-depth is up to several degrees
68 above zero. This warm and saline Atlantic water enters the ocean through the Nordic
69 Seas (Fig. 1).

70 Landfast ice has been found important for the accurate simulation of the sea surface
71 height [*Proshutinsky et al.*, 2007] and sea ice thickness *Johnson et al.* [2012]. Although
72 landfast ice is also important for the processes maintaining the Arctic halocline, it is
73 usually not properly represented in state-of-the-art sea ice-ocean models. In this study
74 we demonstrate the effects of the landfast ice for the Arctic halocline.

75 The outline of this paper is as follows. In section 2 we describe the numerical model for
76 the sensitivity study. In Section 2.1 we describe the details of the landfast ice parametriza-
77 tion. In sections 3, 4 and 5 we present and discuss our results. A summary of our findings
78 and final remarks are given in section 6.

2. Model setup

79 Our model is a regional coupled sea ice - ocean model based on the Massachusetts
80 Institute of Technology General Circulation Model code - MITgcm [*Marshall et al.*, 1997;
81 *MITgcm Group*, 2014] with a model domain covering the Arctic Ocean, Nordic Seas and
82 northern North Atlantic. The horizontal resolution of is $1/4^\circ$ (~ 28 km) on a rotated grid

83 with the grid equator passing through the geographical North Pole. The model has 36
84 vertical levels unevenly distributed in a way that the surface layer is well resolved at the
85 cost of the poor resolution in the deep layer. The shelf bottom topography has realistic
86 details that allow dense brine to flow downslope and off the shelf. Vertical mixing in
87 the ocean is parameterized by a K-Profile Parameterization (KPP) scheme [*Large et al.*,
88 1994] and tracers (temperature and salinity) are advected with an unconditionally stable
89 seventh-order monotonicity preserving scheme [*Daru and Tenaud*, 2004] that requires no
90 explicit diffusivity. The sea ice model is a dynamic-thermodynamic sea-ice model with a
91 viscous-plastic rheology [*Losch et al.*, 2010]. The model has a non-linear free surface and
92 sea ice that depresses the surface ocean layer according to its mass. We avoid numerical
93 issues associated with too thin surface layers when ice gets very thick by using a rescaled
94 vertical z^* -coordinate [*Campin et al.*, 2008] that distributes the excursion of the free sea
95 surface between all vertical levels to the bottom. The same model set-up has been used
96 by *Itkin et al.* [2013] except that now the vertical background diffusivity has been lowered
97 to 10^{-6} m²/s as recommended by *Nguyen et al.* [2009] to achieve a better defined Arctic
98 halocline.

99 The model is initialized by the PHC climatology [*Steele et al.*, 2001] and has initially
100 no sea ice cover. For a spin-up we run the model for 30 years forced by the atmospheric
101 climatology of the Coordinated Ocean Research Experiment (CORE) version 2 based on
102 a reanalysis of the National Center for Atmospheric Research/National Centers for Envi-
103 ronmental Prediction (NCAR/NCEP) [*Large and Yeager*, 2009]. Subsequently the model
104 is driven from 1948 to 1978 by daily atmospheric data also provided by CORE. Our model
105 experiments start in 1979 and continue until 2010. They are forced by the atmospheric

106 reanalysis – The Climate Forecast System Reanalysis (NCEP–CFSR) [*Saha et al.*, 2010].
107 Surface salinity in ice free regions is restored to a mean salinity field (PHC climatology)
108 with a time scale of 180 days to avoid model drift. River runoff was treated like a surface
109 volume flux and it was prescribed for the main Arctic rivers according to the Arctic Ocean
110 Model Intercomparison Project (AOMIP, <http://www.whoi.edu/projects/AOMIP/>) pro-
111 tocol. Open boundaries are formulated following *Stevens* [1991]; they are located at 50°N
112 in the Atlantic and just south of the Bering Strait. Temperature and salinity at the open
113 boundaries are taken from the PHC climatology. The Barents Strait inflow is prescribed
114 as 0.8 Sv and the stream function at open boundary in the Atlantic Ocean derived from
115 a North Atlantic simulation [*Gerdes and Köberle*, 1995] and modified to balance the net
116 volume flow in the model domain.

2.1. Landfast ice Parameterization

117 The mechanisms that determine the landfast ice formation and extent depend of the
118 specific region and are not fully understood. There have been many attempts to model or
119 parameterize landfast ice. *Lieser* [2004] developed a parameterization based on the ratio
120 of sea ice thickness and total water column depth. For a specified threshold ratio the grid
121 cells assigned as landfast ice remained at rest and surface momentum flux into the ocean
122 was set to zero. His approach in a $1/4^\circ$ (28 km) model resulted in too thick landfast ice that
123 even survived the summer. A similar procedure was used in higher horizontal resolution
124 models (3-12 km) by *Johnson et al.* [2012] and *Rozman et al.* [2011]. Both studies use
125 prescribed landfast ice areas obtained from bathymetrical limits or observations. In the
126 Kara Sea, where the landfast ice forms over deep waters behind a row of coastal islands,

127 *Olason* [2012] successfully modeled landfast ice by adjusting the internal sea ice strength
 128 parameters in the viscous-plastic rheology of *Hibler* [1979].

129 Our model grid does not resolve small islands and shallow topographical features in
 130 the Siberian Seas, where the landfast ice might get grounded. Therefore we designed a
 131 simplified and uniform parametrization based on *König Beatty and Holland* [2010] that
 132 takes into account water column depth and landfast ice internal strength. The latter is
 133 justified by a sharp salinity gradient between the shallow shelf waters of the Kara, Laptev
 134 and East Siberian Sea and the deep ocean. In 1992, *Dethleff et al.* [1993] documented
 135 freshwater ice up to 100 km seaward off the Lena Delta. In 1999 the freshwater and
 136 brakish sea ice with salinity below 1 was confined to the coastal waters adjacent to the
 137 eastern Lena Delta with water depth less than 10 m, while river water on average still
 138 contributed 62% of the landfast ice further offshore in the southeastern Laptev Sea [*Eicken*
 139 *et al.*, 2005]. .

140 In the widely used sea ice strength parametrization (e.g. *Hibler* [1979]; *Zhang and*
 141 *Hibler III* [1997]) compressive strength P depends just on the sea ice thickness h and
 142 concentration A :

$$P = P^* h \exp(-C^*(1 - A)), \quad (1)$$

143 where the empirical sea ice strength parameters, $P^* = 2750N/m^2$ and $C^* = 20$ are
 144 constants.

145 The landfast ice parametrization used in this study takes into account the unresolved
 146 shallow topographical features inside the maximal landfast ice edge mark (25 m) and
 147 increased sea ice internal strength attributed to the lower sea ice salinity in the same

148 area by setting the P^* to the double of the drift ice. Such landfast ice would still fail
 149 under strong offshore wind if the distance between the coastline and 25-m-bathymetrical
 150 boundary is large. To prevent this we amended the sea ice rheology in the regions shallower
 151 than 25 m with sea ice tensile strength T following *König Beatty and Holland* [2010]:

$$\zeta = \frac{P + T}{2\Delta}, \quad (2)$$

$$\eta = \frac{P + T}{2\Delta e^2} = \zeta/e^2, \quad (3)$$

and

$$p = P - T, \quad (4)$$

152 where ζ is bulk viscosity, η is shear viscosity, $e = 2$ is eccentricity constant and p is the
 153 pressure term. This moves the elliptical yield curve in the principal stress space into the I.
 154 quadrant, when the water column is shallower than 25 m and leaves the curve unmodified
 155 otherwise. $T = \frac{P}{2}$, which is consistent with the estimates by *Tremblay and Hakakian*
 156 [2006].

157 For the sensitivity study we compare a control run (CTRL) and a landfast ice run (LF)
 158 that differ only in this additional landfast ice parameterization.

3. Impact of the landfast parameterization on the sea ice

159 In the Laptev and East Siberian Sea the landfast ice cover forms in December, reaches
 160 its maximal extent in April, breaks up into fully drifting ice in May or June and melts in
 161 summer. Starting in May, the atmospheric temperatures are too warm to have a significant
 162 amount of brine produced in the polynya. Hence, in this study we define wintertime
 163 as the months December to April. In our model landfast ice breaks up completely in

164 early summer. At atmospheric temperatures above zero and when A drops below 1, sea
165 ice loses its internal strength exponentially making P^* irrelevant, because P depends
166 exponentially on A .

167 If we define coastal sea ice with drift speeds below 1 mm/s as landfast ice, then the
168 parametrization in LF comparing to CTRL leads to substantially larger areas of landfast
169 ice in April, during the annual maximal extent (Figure 3 panels a and b). The simulated
170 landfast ice is slow enough for polynyas to form at its seaward edge (Figure 3 c and d).
171 The 1 mm/s contour and polynyas in LF agree better with the 1997-2006 landfast ice
172 edge produced by the Arctic and Antarctic Research Institute, Saint Petersburg, Russia
173 (*Smolyanitsky et al.* [2007], hereafter AARI dataset). In LF, however, the Laptev Sea
174 polynya is located too far away from the coast. This might be a consequence of the coarse
175 model resolution and the high P^* value. In the Kara Sea our parametrization has only
176 a minor effect because the shelf sea ice is much deeper than the 25 m threshold we use
177 in the parametrization and the small offshore islands there are located in areas where the
178 surrounding seas exceed the depth of 150 m.

179 The effect of the parametrization is clearly visible in the sea ice thickness maps (Fig. 4).
180 LF sea ice is by up to 30 cm thinner under the regions affected by the landfast ice
181 parametrization in the Kara, Laptev, East Siberian, Chukchi and Beaufort Seas. The
182 difference pattern accurately matches the landfast ice extent from the AARI dataset
183 (Fig. 3). The largest negative differences are in the polynya at the landfast ice edge
184 in the East Siberian and Laptev Seas, while in the coastal regions of the Laptev Sea
185 differences are positive to reflect the erroneous polynya location in the CTRL.

186 The thinner sea ice in the landfast ice area in LF is a consequence of exclusive thermo-
187 dynamic sea ice growth. The thermodynamical growth itself is also different between the
188 runs (Fig. 4): in both cases the highest growth rates occur in the polynyas. In LF this is
189 further offshore than in CTRL.

190 As a volume of water is thermodynamically turned into sea ice, salt is expelled from this
191 volume and remains in the ocean. Because this process increases the salinity in the surface
192 ocean, one can define a 'virtual salt flux' $S_{flux} = FW_{flux}S$, where FW_{flux} is freshwater
193 flux and S is the salinity of the surface water minus 4 where the salinity is equal or higher
194 than 4. This approximates the constant sea ice salinity of 4 in our model. The resulting
195 S_{flux} differences resemble the difference pattern in the thermodynamic growth (Fig. 4).
196 Because S_{flux} depends also on the sea surface salinity its is relatively low in the East
197 Siberian Sea although the sea ice production there is high. S_{flux} is an order of magnitude
198 higher and opposite in sign compared to the virtual salt flux generated by the sea surface
199 salinity restoring (not shown). Salinity is only restored to climatology in the ice free part
200 of the model grid cell.

201 The largest differences in the winter sea ice cover between the runs are found in the
202 Laptev and East Siberian Seas. This is the area with the greatest landfast ice extent
203 over the shallow shelf and hence our parametrization has the largest effect there. In time
204 series, we compare the contributions of wintertime sea ice concentration, thickness, sea ice
205 production, ocean surface salinity and S_{flux} to wintertime dense shelf water production in
206 both simulations (Fig. 5). The mean sea ice concentration over the Siberian Seas is high
207 and very similar in both simulations. This is not surprising as the differences between
208 the runs are mainly in the positioning of the polynyas within the Siberian Seas. The sea

ice concentrations are a little lower in CTRL as the drifting sea ice allows more small
leads. Sea ice is about 10 cm thinner in LF than in CTRL. This is a consequence of less
ridging through convergent motion in the landfast ice area and less advection of sea ice
into the region. In contrast, there is more thermodynamically grown sea ice in LF. While
the polynyas in LF are shifted offshore into more saline sea water, a higher fraction of
fresh river water reduces the salinity in the shelf seas. Towards the end of the runs these
counteracting effect leads to a similar surface salinity in polynyas in both runs. Still,
higher sea ice production leads to higher S_{flux} in LF.

The differences in the sea ice production and S_{flux} between CTRL and LF are geograph-
ically distributed as expected so that we conclude that this sensitivity study can be used
to estimate the importance of the landfast ice for the Arctic Ocean halocline stability.

4. Impact of the landfast parameterization on the Arctic halocline

In this section we compare the end of winter (April) climatologies (2000-2010) because
of the high seasonality of the mixed layer depth in the Arctic. First we examine the
buoyancy, $b = -g \frac{\rho - \rho_0}{\rho_0}$ at the top of the halocline (25 – 30 m). The differences between
LF and CTRL show regional effects of the landfast ice parametrization with less buoyancy
in LF in the Canadian Basin and more buoyancy centered in Makarov and Eurasian Basins
(Fig. 6a and b).

To understand the relevance of these differences to the water column stability we ex-
amine three oceanographic sections; each section runs through one of the major Arctic
basins: Canadian, Makarov and Eurasian as shown on Fig. 1 and 6. In the first section
running from the Eurasian coast over the Wrangel Island to the Ellesmere Island (Fig. 7)
there is increased salinity by locally up to 0.6 at the halocline layer in LF (25 – 30 m),

231 while the surface layer tends to be fresher than in CTRL. Especially in the surface layer,
232 the differences are not homogeneous and locally have alternating signs. Along with the in-
233 creased temperature at the top halocline depth this suggests this layer is fed by relatively
234 warm shelf bottom water mixed with brine originating from the sea ice production. This
235 influence is supported by the Hovmöller diagrams of differences at the section across the
236 pathway of the water masses from the East Siberian Sea to the Canadian Basin (Fig. 8
237 a and b). In LF, this shelf water has continuously larger salinity at the halocline depth
238 (panel a), while the surface layer is typically fresher which generally is associated with
239 higher river water content (panel b). The salinity and river water fraction in the surface
240 layer differences are not constant in time. Notably the salinity difference is positive in the
241 early 1990s and in the mid 2000s. Both events are accompanied with a negative difference
242 in river water fraction. The differences in temperature and salinity described above have
243 consequences for the stratification. The buoyancy frequency N^2 is larger in LF than in
244 CTRL pointing to a higher stability at the winter mixed layer depth.

245 In contrast, the second section which runs from East Siberian Sea to the Northern
246 Greenland (Fig. 9) shows a different pattern with higher salinities by up to 0.8 in the
247 surface layer that reflects a decreased river water content in LF compared to CTRL. The
248 salinity at the depth of the halocline only increases only by about 0.1, such that the overall
249 stability of the upper halocline is reduced. The third section runs from the Laptev Sea
250 to the Fram Strait (Fig. 10). Here the differences between the runs are similar as in the
251 second section, but not as pronounced. While there is still less river water in the surface
252 layer, the halocline layer has more river water in LF than in CTRL. Here, the landfast
253 ice parametrization has no clear effect on the halocline stability. This is again supported

254 by the Hovmöller diagrams of differences at the section across the pathway of shelf water
255 from the Laptev Sea (Fig. 8 c and d). In LF, the surface layer is mostly more saline,
256 while the time series of the river water fraction differences show a strong seasonal cycle
257 with typically less river water in the surface layer in the winter and more in the summer
258 (not shown).

5. Discussion

259 Landfast ice simulated in LF grows almost exclusively thermodynamically and allows
260 very little ice production due to convergent motion (dynamical sea ice growth). Therefore
261 the coastal regions covered by landfast ice in LF are thinner than in CTRL. To some
262 extent, the arbitrary choice of the 25 m isobath for turning on the parameterization place
263 the polynyas in the approximately observed positions. In contrast, in CTRL polynyas are
264 located directly on the coast. The shift of the polynyas from the coast in CTRL toward
265 a realistic landfast ice edge location in LF moves the brine production closer to the
266 shelfbreak and into more saline ocean where more brine is produced per sea ice volume.

267 To determine if the amount of brine produced in our simulations is realistic we first
268 compare the available sea ice production estimates for the Laptev Sea. In LF during
269 2000s a mean of 144 km^3 of sea ice is produced per winter (not shown, but Fig. 5a
270 shows winter sea ice production for the East Siberian and Laptev Sea, roughly half of
271 which is produced in that Laptev Sea. That is above *Rabenstein et al.* [2013]'s estimate
272 of $94 \pm 27 \text{ km}^3$ for the southeastern Laptev Sea for winter (late December till mid-April)
273 2007/2008. *Willmes et al.* [2011] estimate a lower $55 \pm 15 \text{ km}^3$ for the entire Laptev Sea
274 during winter, but these authors only took into account ice production in areas with sea
275 ice thinner than 20 cm.

276 Our results indicate that with an appropriate representation of landfast ice more river
277 water is stored in the Siberian Seas. While in CTRL more river water is dispersed by the
278 wind acting through the drifting sea ice, in LF this river water is protected from the wind
279 by the immobile shield of the landfast ice and remains on the shelf. From there the river
280 water plume is driven northeastward by geostrophic currents (black arrows on Fig. 1)
281 into the Canadian Basin and less river water reaches other parts of the Arctic. In LF,
282 in summers with strong offshore winds the increased amounts of the river water stored
283 on the shelf during the winter are driven northwards into the European and Makarov
284 Basins. Still, these river water pulses do not change the sign of the salinity differences in
285 the surface layer and do not affect the halocline stability.

286 The combined effect of increased brine export to the Central Arctic and redistribution of
287 the river water from the winter shelf in LF compared to the CTRL changes the buoyancy
288 at the Arctic halocline. In the Canadian Basin the surface layer freshens locally due to
289 increased river water content as a consequence of the landfast ice parametrization, while
290 the halocline layer becomes saltier due to the difference in the brine production in the
291 Siberian polynyas. Both surface and halocline depth differences make the halocline in LF
292 stronger. Our simulation supports *Nguyen et al.* [2012]’s conclusions that the Canadian
293 Basin halocline is maintained by the dense water formed on the Eurasian shelf.

294 In the Makarov and Eurasian Basin the difference signal is dominated by the increase
295 in the surface salinity that is again the consequence of the redirection of the river water to
296 the Canadian Basin. Consequently the already weak halocline in this sector of the Arctic
297 [*Boyd et al.*, 2002; *Rudels et al.*, 2004] is further eroded in LF. The effect of the landfast
298 ice resembles that of the Great salinity anomaly [*Steele and Boyd*, 1998; *Johnson and*

299 *Polyakov*, 2001] observed in the early 1990s, where the changes in the wind circulation
300 enhanced the river water export to the Canadian Basin and increased brine production
301 in the Siberian Seas polynyas that eroded the Arctic halocline over substantial parts
302 of the Arctic Ocean. During the Great salinity anomaly and similarly also in the mid
303 2000s landfast ice in LF has an effect outweighing the anomalous offshore winds and the
304 differences in surface salinity and river water content of the shelf waters originating from
305 the East Siberian Sea during these two periods change the sign (Fig. 8a and b).

306 *Yu et al.* [2014] detected a general decrease of the Arctic landfast ice extent of 7 ± 1.5 %
307 per decade in the period from 1976 till 2007. In the Chukchi, Laptev and East Siberian
308 Seas also the landfast ice season is getting shorter, reflecting the general negative trend
309 in the summer sea ice extent in the Siberian Seas [*Comiso et al.*, 2008; *Stroeve et al.*,
310 2012]. According to our results this trend should have implications for the halocline
311 stability. The reduction in the extent and duration in the Siberian Seas would lead to
312 a weaker halocline in the Canadian Basin, but conversely also to a stronger halocline in
313 the Makarov and Eurasian Basins. Especially for the latter two basins the Atlantic
314 water layer has been reported to be warming [*Polyakov et al.*, 2010]. In such a case, the
315 loss of the landfast ice would inhibit the heat transport from the Atlantic Water layer to
316 the surface and delay the further rapid sea ice loss in the Arctic.

6. Summary and conclusions

317 An accurate representation of the landfast ice in a sea ice-ocean coupled model has an
318 impact not only on the winter sea ice and brine production but also on the river water
319 distribution. The landfast ice in the LF shields the river water that remains on the shelf
320 in winter from the winds and promotes northeastward flow of the river water with the

321 coastal currents. We have been able to show this in a sensitivity study with a simple-to-
322 implement landfast ice parametrization that generated landfast ice over extensive areas
323 of the Laptev and East Siberian Seas. In LF more river water reaches the Canadian
324 Basin and less the Makarov and Eurasian Basin. Also the polynyas are located closer
325 to the shelf break. From there more brine reaches the halocline depth of the Canadian
326 Basin. The surface salinity decrease due to the river water and the halocline salinity
327 increase due to the brine both strengthen the halocline in the Canadian Basin, whereas
328 in the Makarov and Eurasian Basins the surface salinity increase erodes the already low
329 halocline stability. Consequently in the latter two basins where the Atlantic water layer
330 is still relatively warm and the halocline is already eroded, extensive landfast ice further
331 weakens the halocline and promotes entraining Atlantic water into the ocean mixed layer
332 that might lead to sea ice melt.

333 Based on our simulations we recommend to include our landfast ice parametrization to
334 those regional numerical models that address the halocline stability and shelf - deep basin
335 exchanges. The stratification modified by brine produced in the polynyas and the river
336 water that supplies the Arctic Ocean with sediments, nutrients, and pollutants are key
337 factors for biogeochemical processes.

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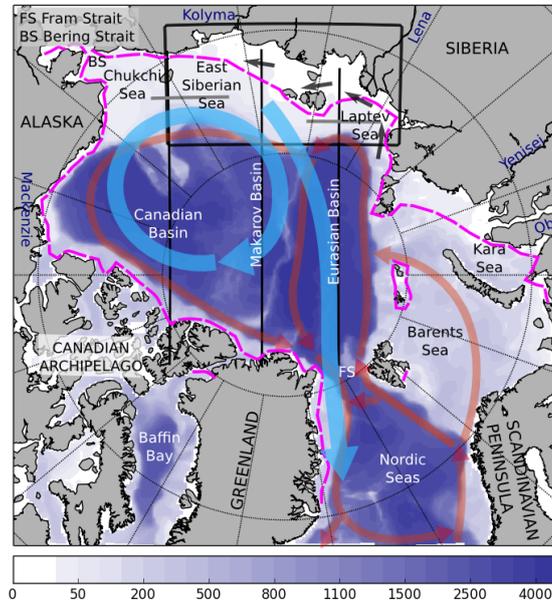


Figure 1. The Arctic Ocean and its marginal seas. The sea ice and surface circulation is schematically represented by light blue and mid-depth circulation by red arrows (simplified from *Rudels et al.* [1994]), respectively. The landfast ice edge is depicted by the magenta dash line. The coastal current in the Laptev and East Siberian Seas is schematically represented by the black arrows. The black lines and box mark the sections and regions used in the model analysis.

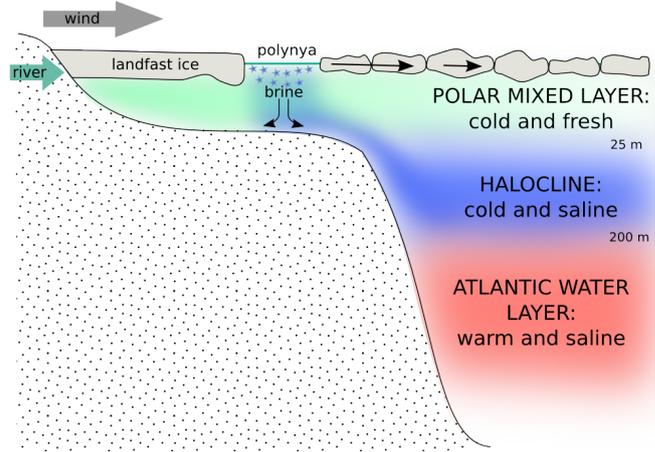


Figure 2. Arctic Ocean vertical stratification and winter processes maintaining the Arctic halocline.

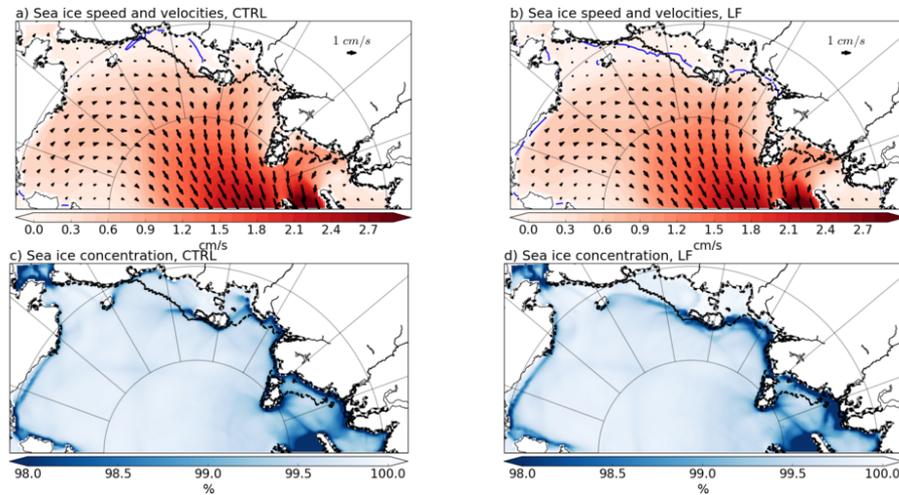


Figure 3. The effect of the landfast ice parametrization on the mean April (2000-2010) sea ice concentration (a,b) and motion (c,d). a,c - CTRL, b,d - LF. Speed 1 mm/s is contoured by blue line. Mean April landfast edge from the AARI dataset (1997-2006) is depicted by gray dash line.

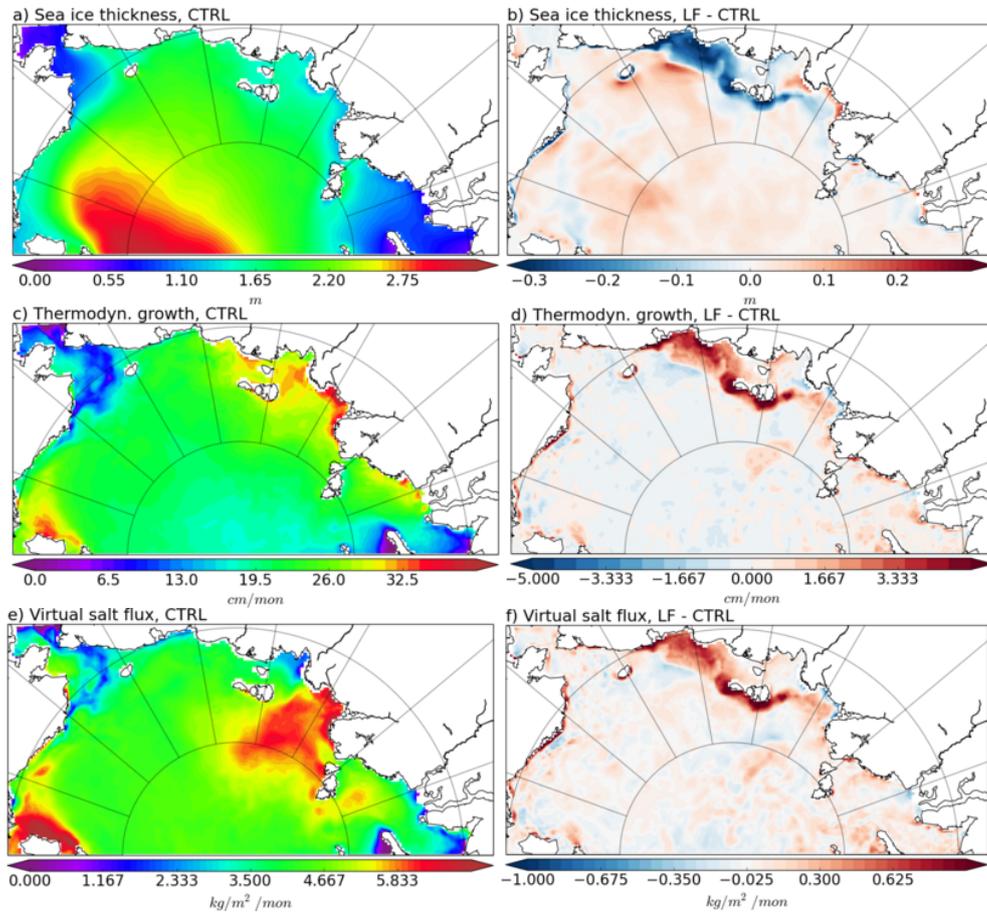


Figure 4. Mean winter (2000-2010) sea ice thickness, thermodynamical growth and salt flux from the sea ice thermodynamical growth: a - CTRL, b - LF minus CTRL.

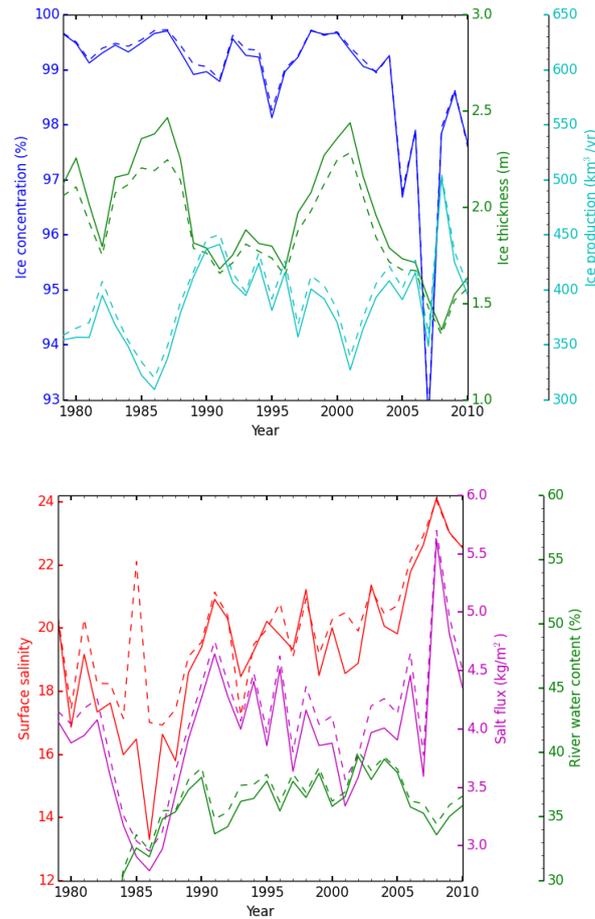


Figure 5. Mean winter (December-April) sea ice time series for the Laptev Sea and East Siberian Sea. CTRL and LF are represented by solid and dash lines, respectively: a - mean sea ice concentration and thickness with total sea ice production; b - sea surface salinity of the areas with the production higher than 30 cm, salt flux resulting from the production and river water content.

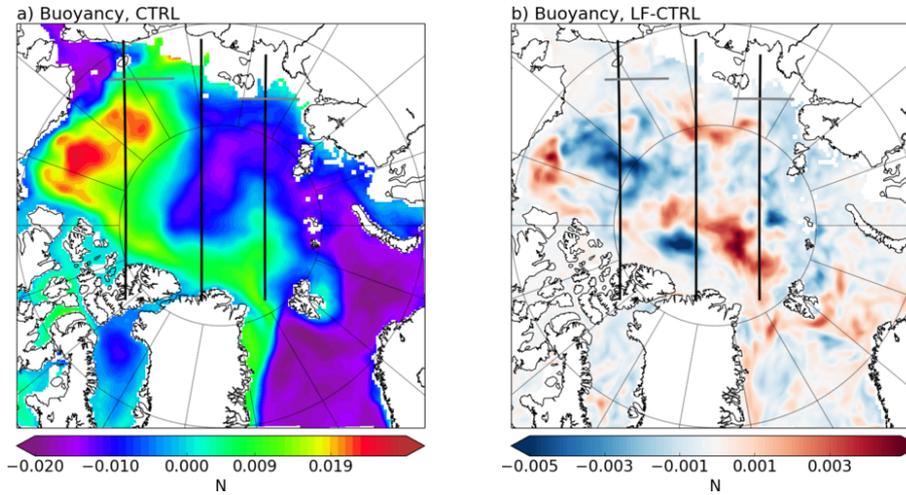


Figure 6. Mean April (2000-2010) buoyancy at the top halocline (25-30 m): a - CTRL, b - LF minus CTRL. Black lines mark the profile sections across Canadian, Makarov and Eurasian Basins (Figs. 7, 9, and 10). Gray line marks a section for the time series at the Eurasian shelf break (Fig. 8).

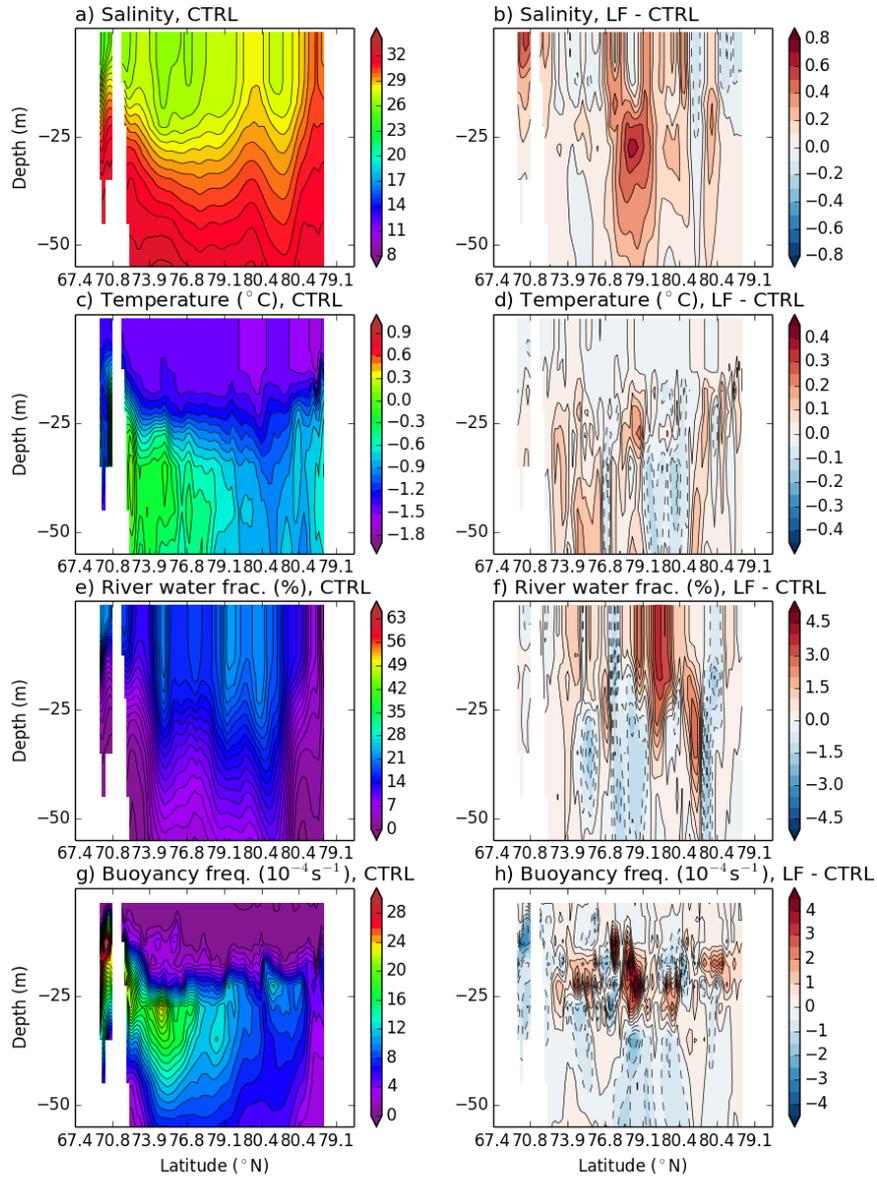


Figure 7. Mean April (2000-2010) salinity, temperature, river water fraction and buoyancy frequency along the section across the Canadian Basin.

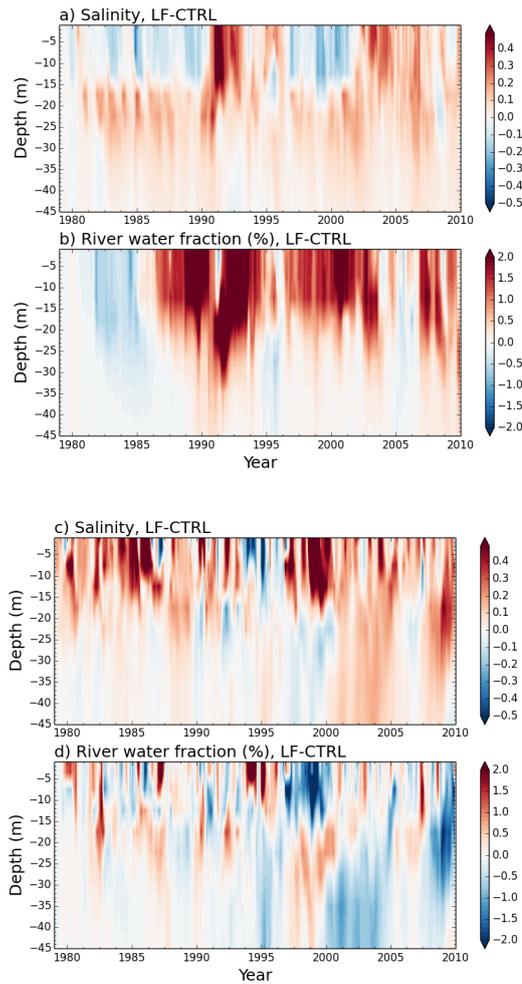


Figure 8. Hovmöller diagrams of monthly mean salinity and river water fraction differences between LF and CTRL: average over an oceanographic section at the Eurasian shelf break north of the East Siberian Sea (a,b) and at the Laptev Sea shelf break (c,d). Sections are marked by gray lines on Fig. 1 and 6.)

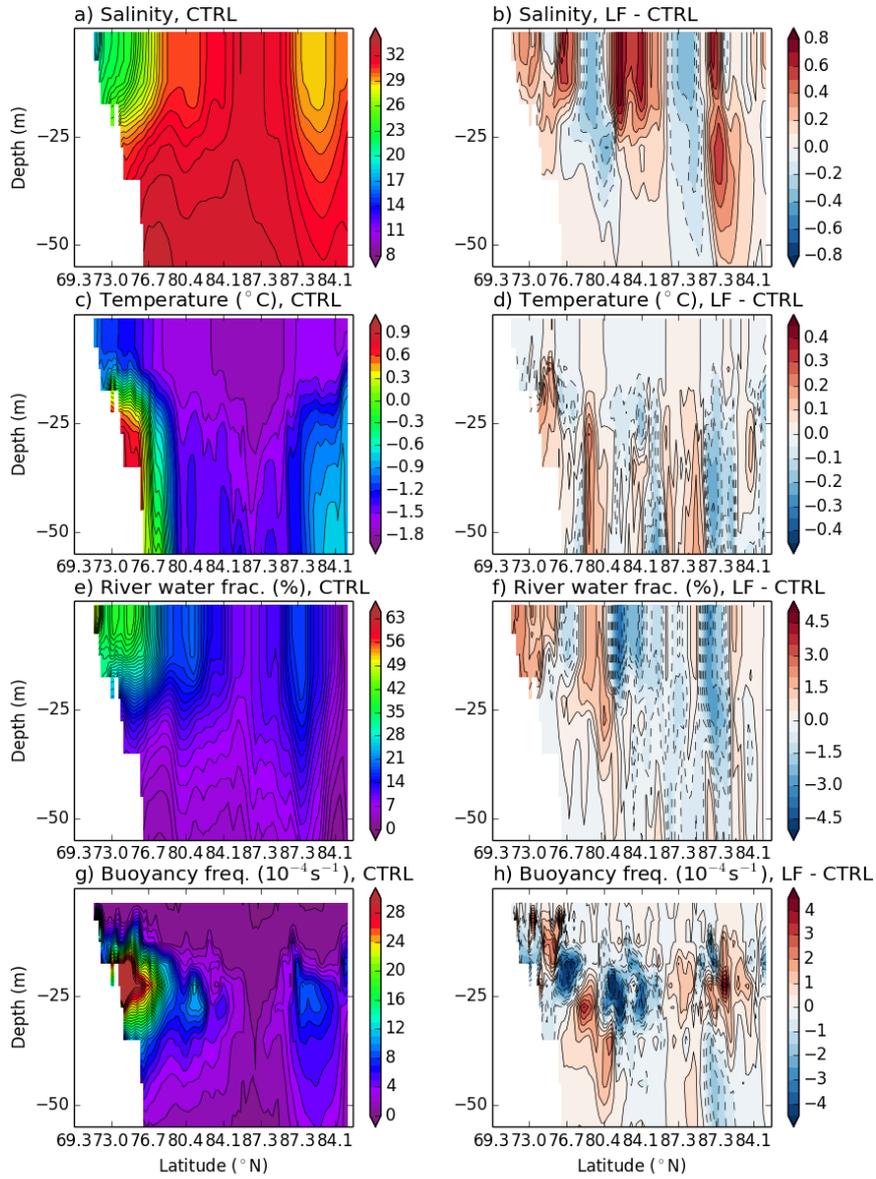


Figure 9. Mean April (2000-2010) salinity, temperature, river water fraction and buoyancy frequency along the section across the Makarov Basin

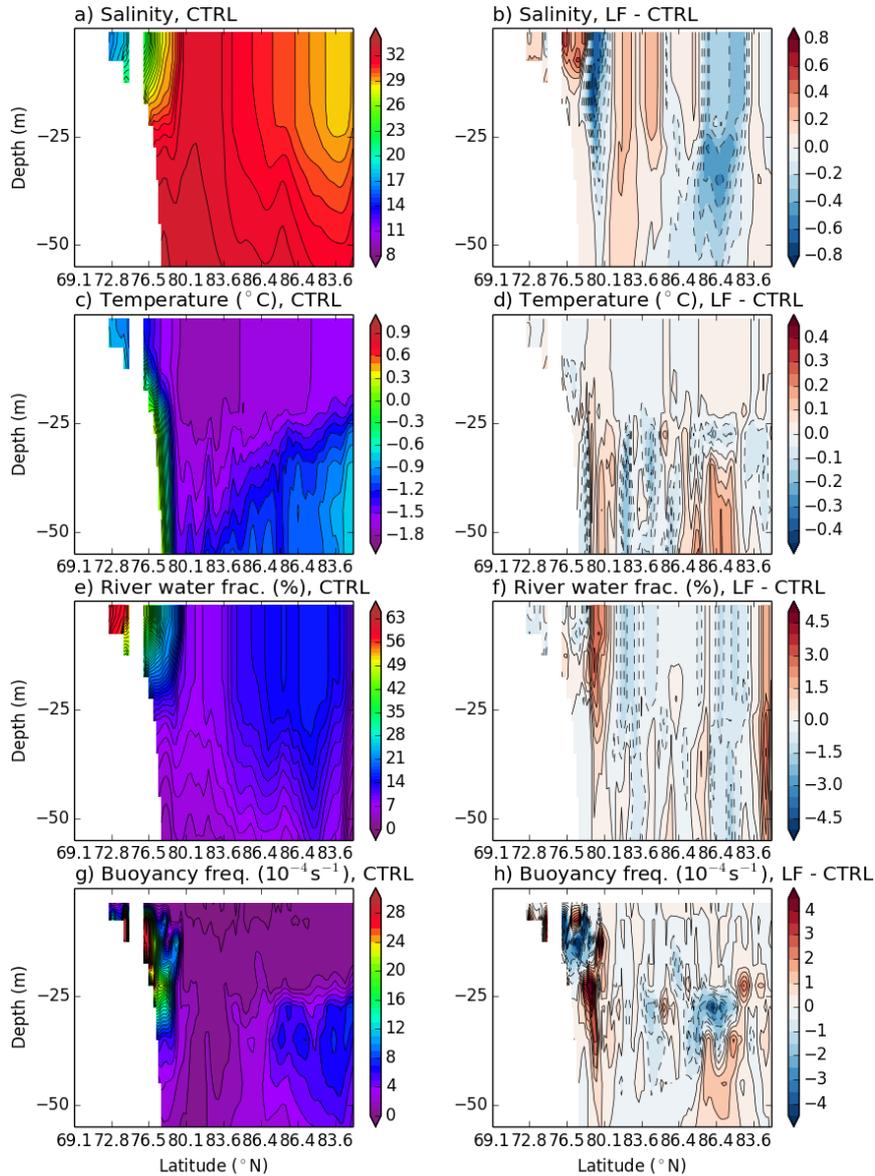


Figure 10. Mean April (2000-2010) salinity, temperature, river water fraction and buoyancy frequency along the section across the Eurasian Basin.