

Sea ice dynamics solvers in the MITgcm

Martin Losch (Alfred-Wegener-Institut, Bremerhaven)

Jean-Michel Campin (MIT, Cambridge, MA)

Complicated dynamics



which satellite?

Complicated dynamics



MITgcm (Menemenlis, Hill)

Outline: Sea-ice solvers in MITgcm



- Picard solvers (LSR, Krylov)
 - JFNK solver
 - EVP solvers: mEVP, aEVP
-
- new MEB rheology in the pipeline

sea ice dynamics are very non-linear

$$m \frac{\partial \mathbf{u}}{\partial t} = \nabla \cdot \boldsymbol{\sigma} + R, \quad R = \text{other terms}$$

$$\text{with } \sigma_{ij} = \frac{P}{2\Delta} \left\{ 2\dot{\epsilon}_{ij} e^{-2} + [(1 - e^{-2})(\dot{\epsilon}_{11} + \dot{\epsilon}_{22}) - \Delta] \delta_{ij} \right\}$$

with abbreviations

$$\Delta = \sqrt{(\dot{\epsilon}_{11} + \dot{\epsilon}_{22})^2 + e^{-2} [(\dot{\epsilon}_{11} - \dot{\epsilon}_{22})^2 + 4\dot{\epsilon}_{12}^2]}$$

$$\dot{\epsilon}_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) \quad (\text{strain rates})$$

$$\longrightarrow m \frac{\partial \mathbf{u}}{\partial t} \propto \frac{\partial}{\partial x_i} \left(\frac{P}{\Delta} \frac{\partial u_i}{\partial x_j} \right) + \text{similar terms}$$

$$\mathbf{A}(\mathbf{u}) \cdot \mathbf{u} = \mathbf{b}$$

$$\Rightarrow \text{solve} \quad \mathbf{A}(\mathbf{u}_{n-1}) \cdot \mathbf{u}_n = \mathbf{b}$$

- traditional method, e.g., PSOR, Hübner (1979), LSOR, Zhang and Hübner (1997), (Gauss-Seidel) for linear solver
- Krylov method for linear solver (Lemieux and Tremblay, 2009), requires preconditioner
- stable, but slow

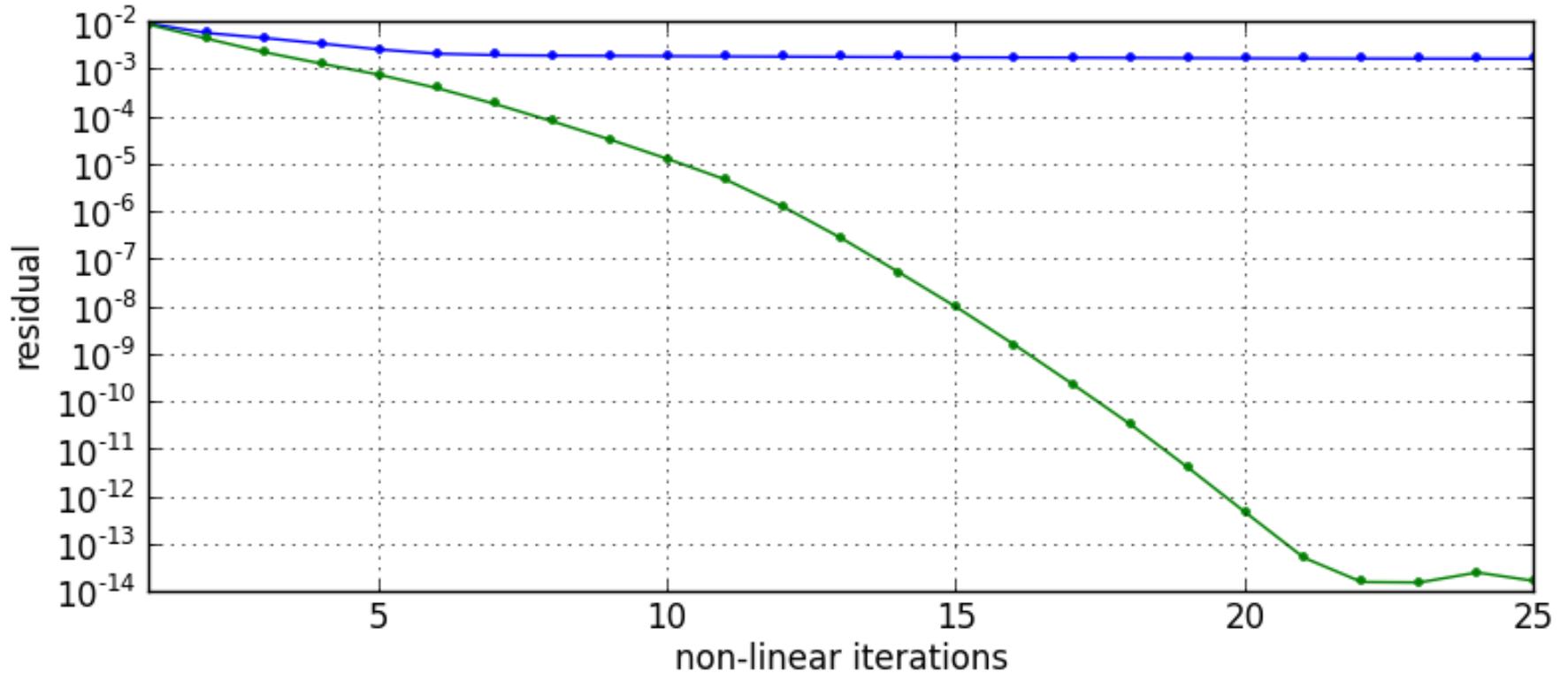
$$\mathbf{F}(\mathbf{u}) = \mathbf{A}(\mathbf{u}) \cdot \mathbf{u} - \mathbf{b}$$

$$\mathbf{F}(\mathbf{u}_n) = \mathbf{F}(\mathbf{u}_{n-1}) + \mathbf{F}' \Big|_{\mathbf{u}_{n-1}} \delta \mathbf{u} \stackrel{!}{=} 0$$

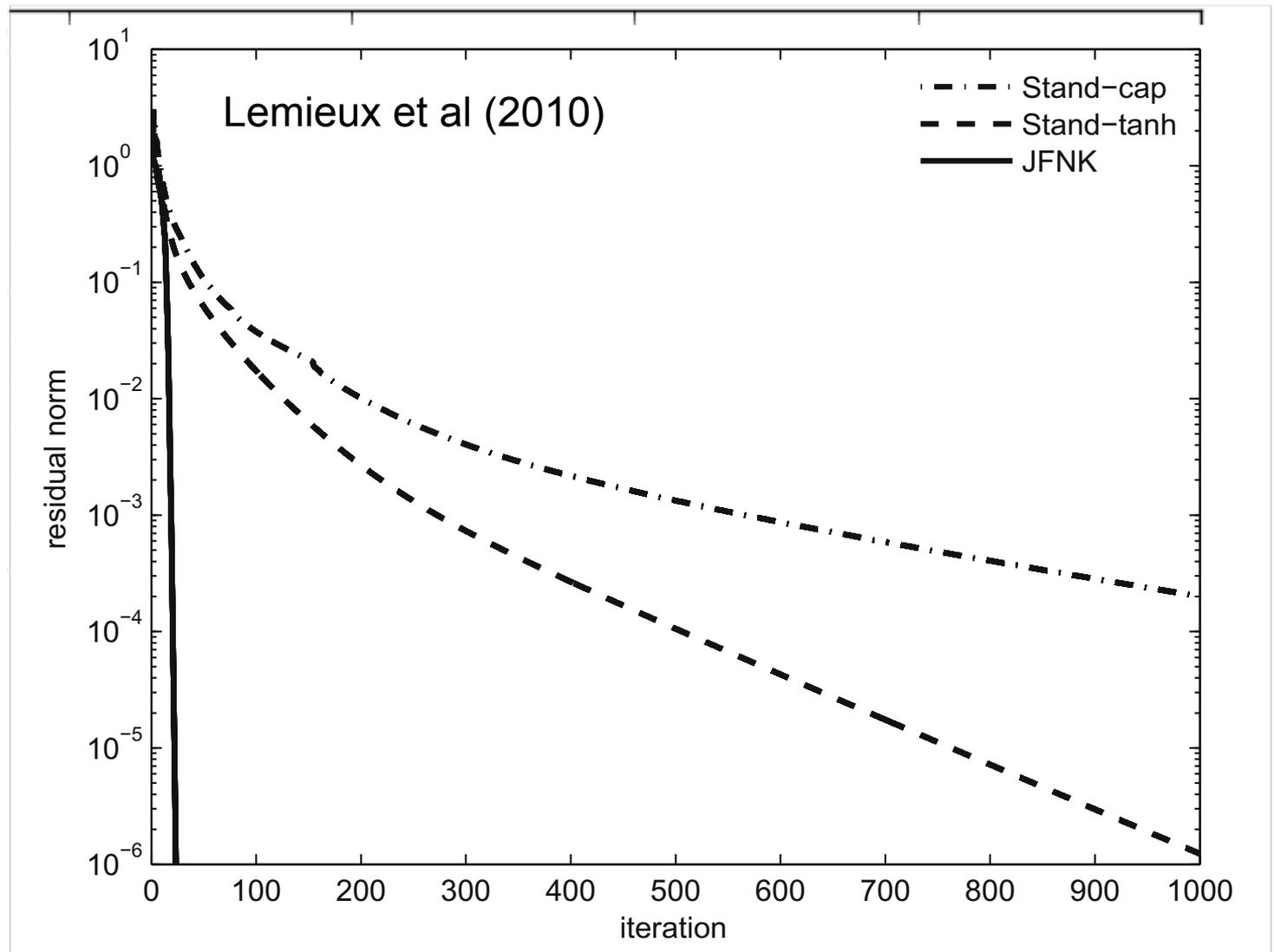
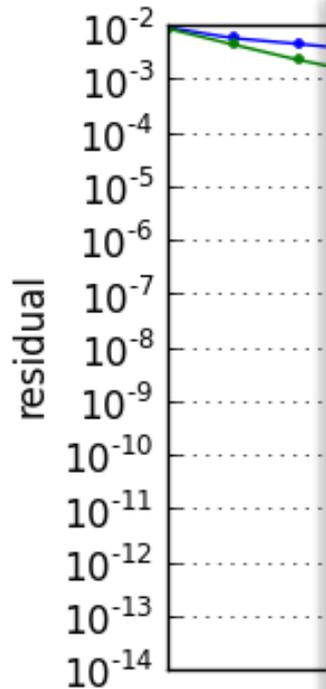
$$\Rightarrow \text{solve} \quad \mathbf{F}'_{n-1} \delta \mathbf{u} = -\mathbf{F}(\mathbf{u}_{n-1}) \quad \Rightarrow \quad \mathbf{u}_n = \mathbf{u}_{n-1} + \delta \mathbf{u}$$

- better (quadratic) convergence near minimum (Lemieux et al. 2010, 2012, Losch et al 2014)
- preconditioner for Krylov solver necessary
- expensive
- unstable, especially at high resolution
- stabilization (e.g. Mehlmann and Richter 2017, involves mixing JFNK and Picard methods)

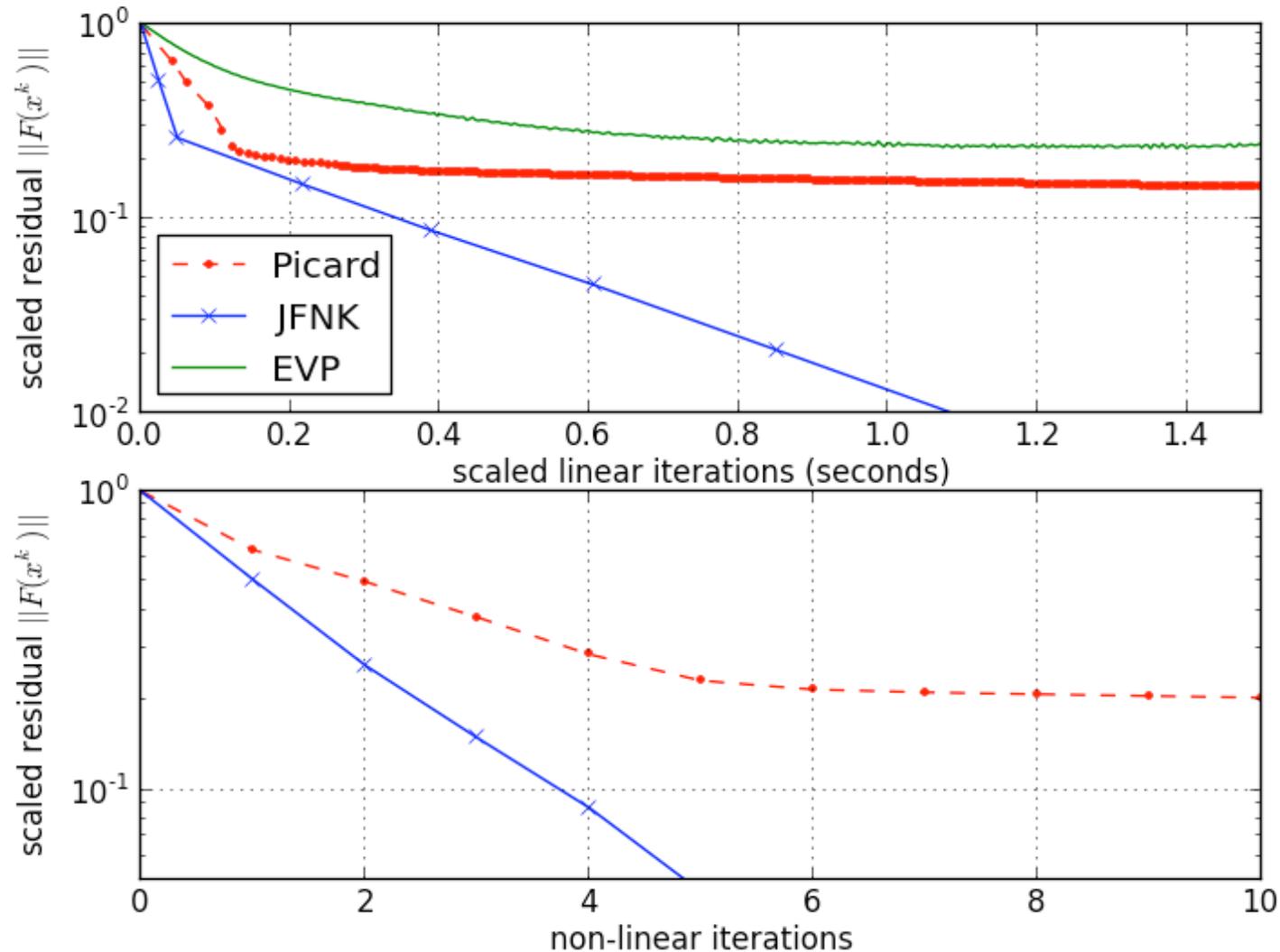
Picard vs. JFNK



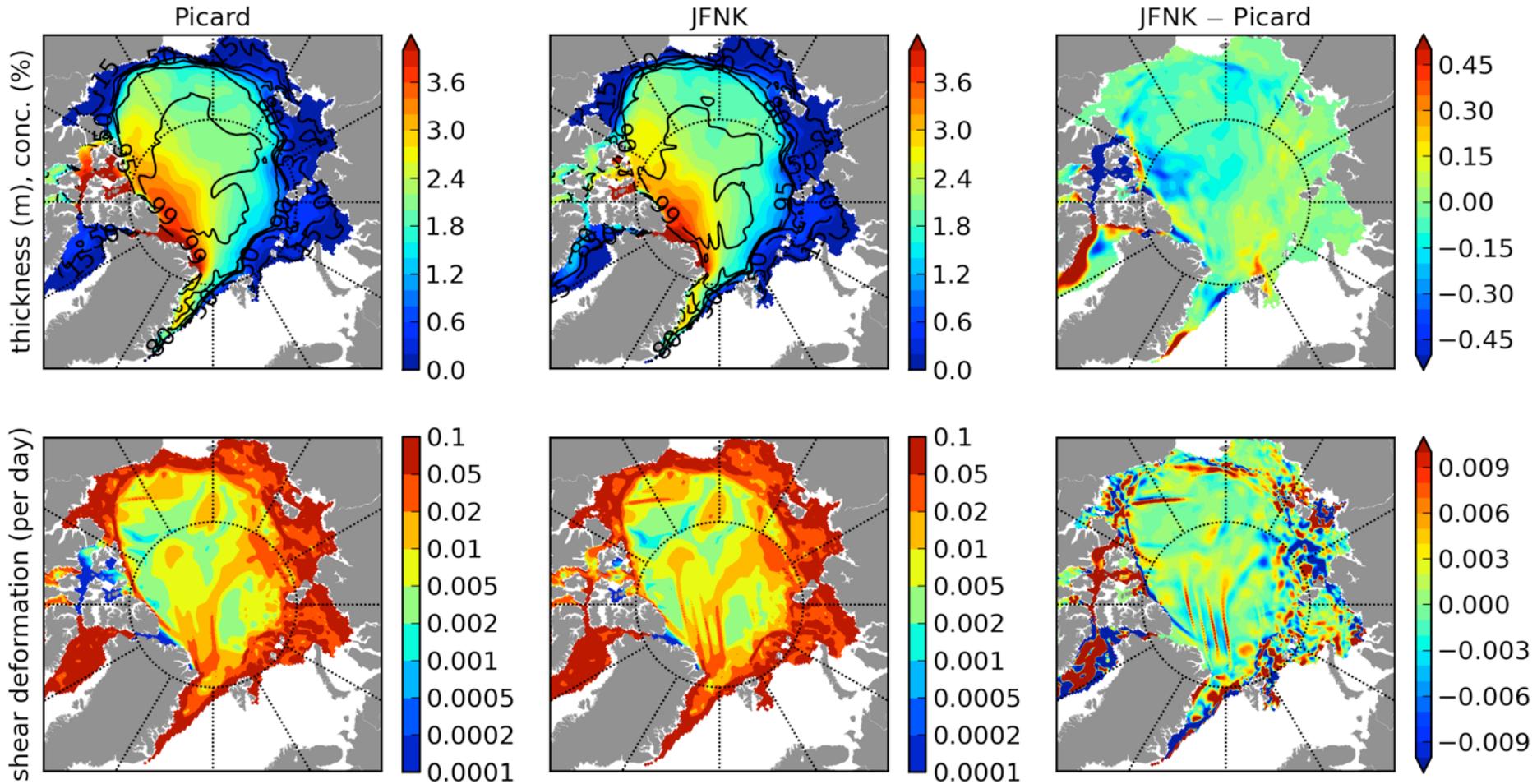
Picard vs. JFNK



“Timing” of solvers



Does it matter?



$$\sigma_{ij} = \frac{P}{2\Delta} \left\{ 2\dot{\epsilon}_{ij} e^{-2} + [(1 - e^{-2})(\dot{\epsilon}_{11} + \dot{\epsilon}_{22}) - \Delta] \delta_{ij} \right\}$$

$$\Leftrightarrow \frac{\Delta e^2}{P} \sigma_{ij} + \left[\frac{\Delta(1 - e^2)}{2P} (\sigma_{11} + \sigma_{22}) + \frac{\Delta}{2} \right] \delta_{ij} = \dot{\epsilon}_{ij}$$

- Hunke and Dukowicz (1997)
- does not converge (definitely not to VP, Lemieux et al. 2012, Losch and Danilov 2012)
- adding inertial term to momentum equations fixes convergence (Lemieux et al. 2012, Bouillon et al 2013)
- m(odified)EVP, a(daptive)EVP (Kimmritz et al 2015, 2016, 2017)

$$\sigma_{ij} = \frac{P}{2\Delta} \left\{ 2\dot{\epsilon}_{ij} e^{-2} + [(1 - e^{-2})(\dot{\epsilon}_{11} + \dot{\epsilon}_{22}) - \Delta] \delta_{ij} \right\}$$

$$\Leftrightarrow \left(\frac{1}{E} \frac{\partial \sigma_{ij}}{\partial t} + \right) \frac{\Delta e^2}{P} \sigma_{ij} + \left[\frac{\Delta(1 - e^2)}{2P} (\sigma_{11} + \sigma_{22}) + \frac{\Delta}{2} \right] \delta_{ij} = \dot{\epsilon}_{ij}$$

- Hunke and Dukowicz (1997)
- does not converge (definitely not to VP, Lemieux et al. 2012, Losch and Danilov 2012)
- adding inertial term to momentum equations fixes convergence (Lemieux et al. 2012, Bouillon et al 2013)
- m(odified)EVP, a(daptive)EVP (Kimmritz et al 2015, 2016, 2017)

New EVP equations

based on Lemieux et al. (2012), Bouillon et al. (2013),
add “inertial-like” term to momentum equations

$$\boldsymbol{\sigma}^{p+1} - \boldsymbol{\sigma}^p = \frac{1}{\alpha} \left(\boldsymbol{\sigma}(\mathbf{u}^p) - \boldsymbol{\sigma}^p \right),$$

$$\mathbf{u}^{p+1} - \mathbf{u}^p = \frac{1}{\beta} \left(\frac{\Delta t}{m} \nabla \cdot \boldsymbol{\sigma}^{p+1} + \frac{\Delta t}{m} \mathbf{R}^{p+1/2} \right)$$

New EVP equations

based on Lemieux et al. (2012), Bouillon et al. (2013),
add “inertial-like” term to momentum equations

$$\boldsymbol{\sigma}^{p+1} - \boldsymbol{\sigma}^p = \frac{1}{\alpha} \left(\boldsymbol{\sigma}(\mathbf{u}^p) - \boldsymbol{\sigma}^p \right),$$

$$\mathbf{u}^{p+1} - \mathbf{u}^p = \frac{1}{\beta} \left(\frac{\Delta t}{m} \nabla \cdot \boldsymbol{\sigma}^{p+1} + \frac{\Delta t}{m} \mathbf{R}^{p+1/2} + \mathbf{u}_n - \mathbf{u}^p \right)$$

New EVP equations

based on Lemieux et al. (2012), Bouillon et al. (2013),
add “inertial-like” term to momentum equations

$$\boldsymbol{\sigma}^{p+1} - \boldsymbol{\sigma}^p = \frac{1}{\alpha} \left(\boldsymbol{\sigma}(\mathbf{u}^p) - \boldsymbol{\sigma}^p \right),$$

$$\mathbf{u}^{p+1} - \mathbf{u}^p = \frac{1}{\beta} \left(\frac{\Delta t}{m} \nabla \cdot \boldsymbol{\sigma}^{p+1} + \frac{\Delta t}{m} \mathbf{R}^{p+1/2} + \mathbf{u}_n - \mathbf{u}^p \right)$$

now, with $\boldsymbol{\sigma}^{p+1} = \lim_{p \rightarrow \infty} \boldsymbol{\sigma}^p$ and $\mathbf{u}_{n+1} := \lim_{p \rightarrow \infty} \mathbf{u}^p$

the discretized equations converge to true VP

$$\frac{m}{\Delta t} \left(\mathbf{u}_{n+1} - \mathbf{u}_n \right) = \nabla \cdot \boldsymbol{\sigma}(\mathbf{u}_{n+1}) + \mathbf{R}^*$$

with $\mathbf{R}^* := \lim_{p \rightarrow \infty} \mathbf{R}^{p+1/2}$

New EVP equations

New momentum equations

$$\boldsymbol{\sigma}^{p+1} - \boldsymbol{\sigma}^p = \frac{1}{\alpha} \left(\boldsymbol{\sigma}(\mathbf{u}^p) - \boldsymbol{\sigma}^p \right),$$

$$\mathbf{u}^{p+1} - \mathbf{u}^p = \frac{1}{\beta} \left(\frac{\Delta t}{m} \nabla \cdot \boldsymbol{\sigma}^{p+1} + \frac{\Delta t}{m} \mathbf{R}^{p+1/2} + \mathbf{u}_n - \mathbf{u}^p \right)$$

with $\alpha\beta \gg \gamma = \frac{P}{2\Delta} \frac{(c\pi)^2}{A} \frac{\Delta t}{m}$

from stability analysis (Kimmritz et al, 2015, 2016).

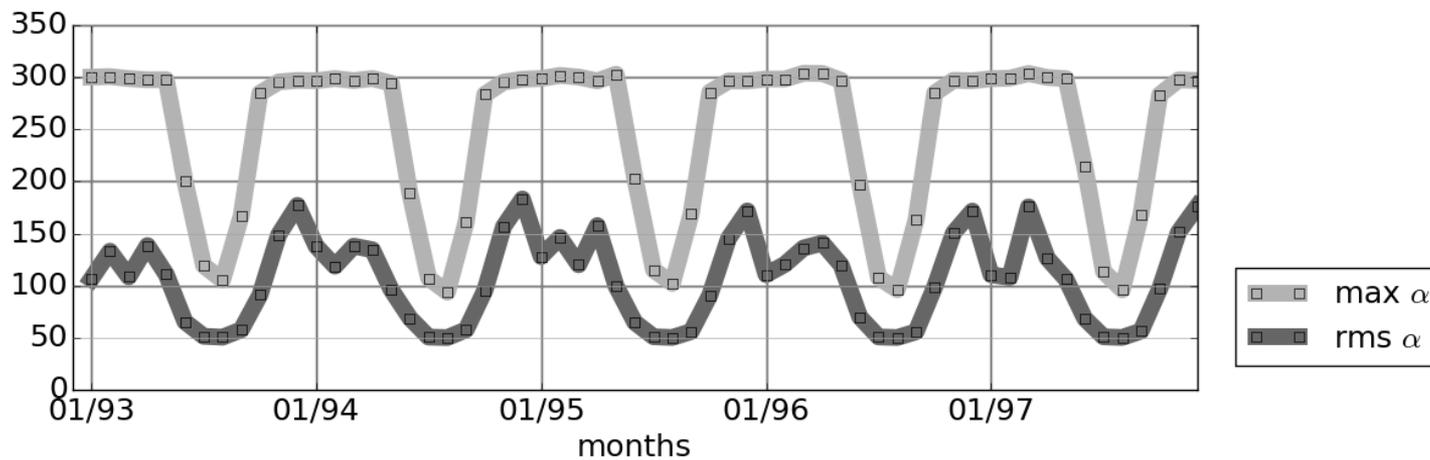
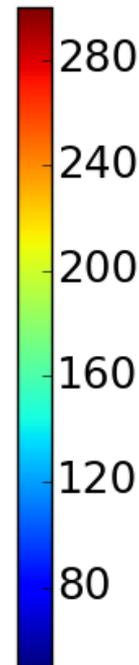
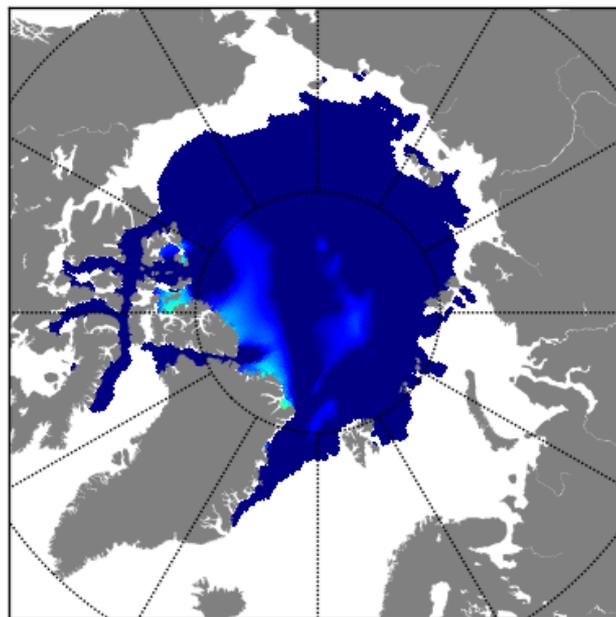
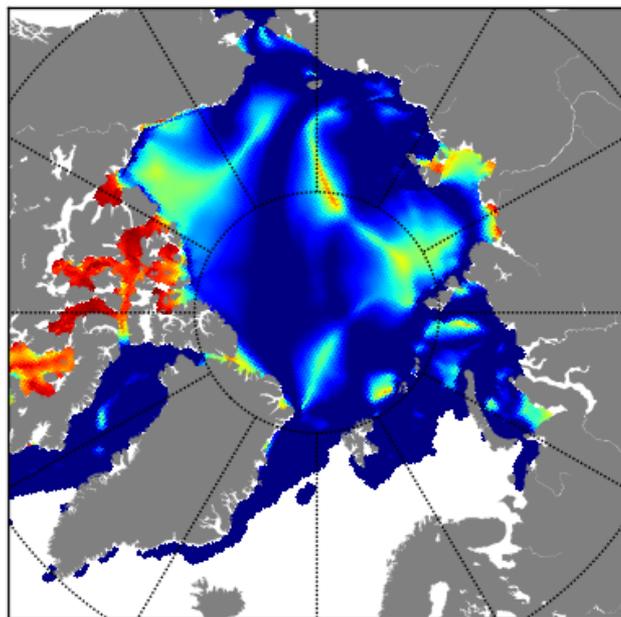
modified EVP: $\alpha, \beta = \text{constant, order}(300)$

adaptive EVP: $\alpha = \beta = (4\gamma)^{1/2}$

Parameter α (aEVP, N = 500)

31/03/93

30/09/93

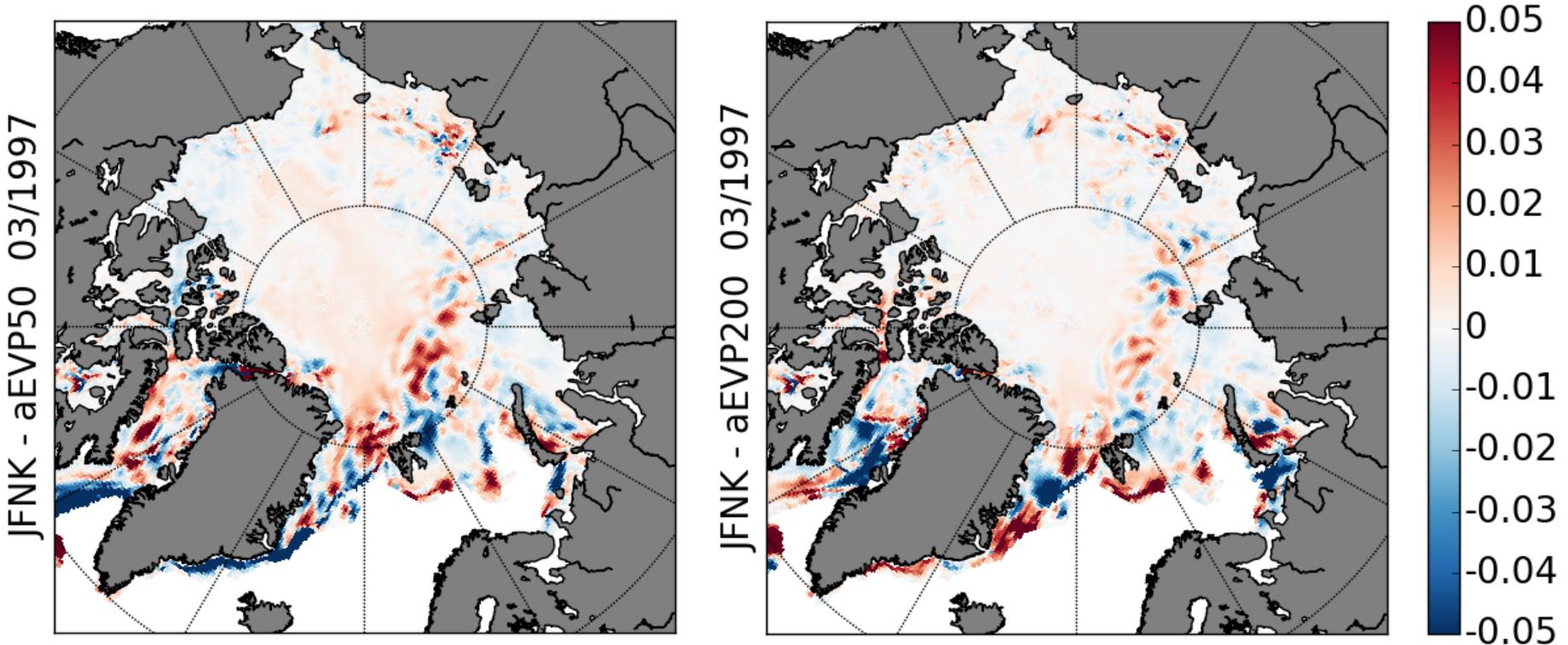


Kimmitz, Losch, Danilov (2017)

Convergence to VP solution: JFNK — aEVP difference in ice thickness (m) at 27 km resolution

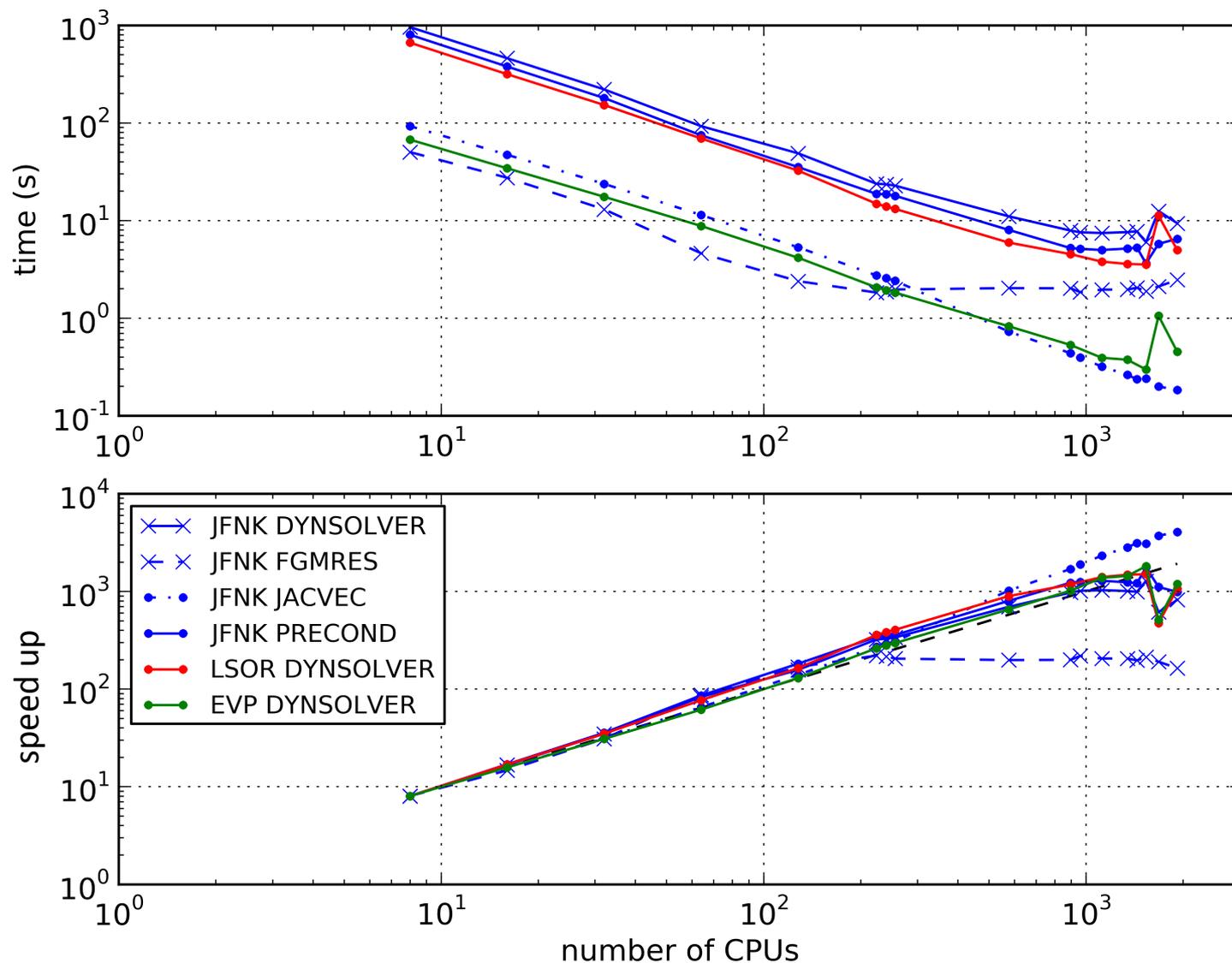
50 EVP iterations

200 EVP iterations



Kimmitz et al. (2017)

MITgcm as a testbed: scalability



high resolution simulations

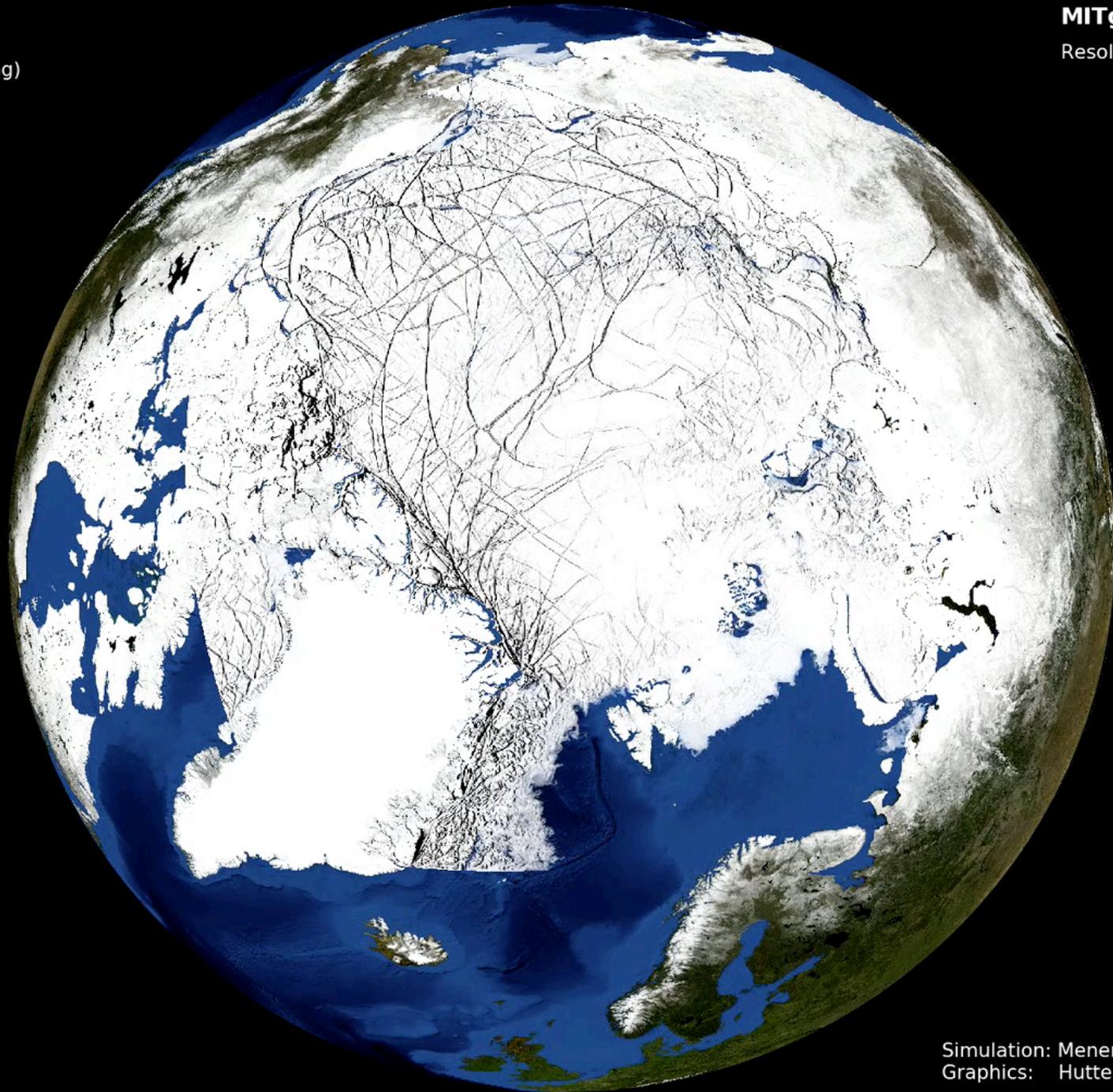
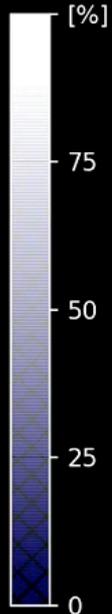


Sea Ice

Concentration (Opacity)
and Thickness (Shadowing)

MITgcm

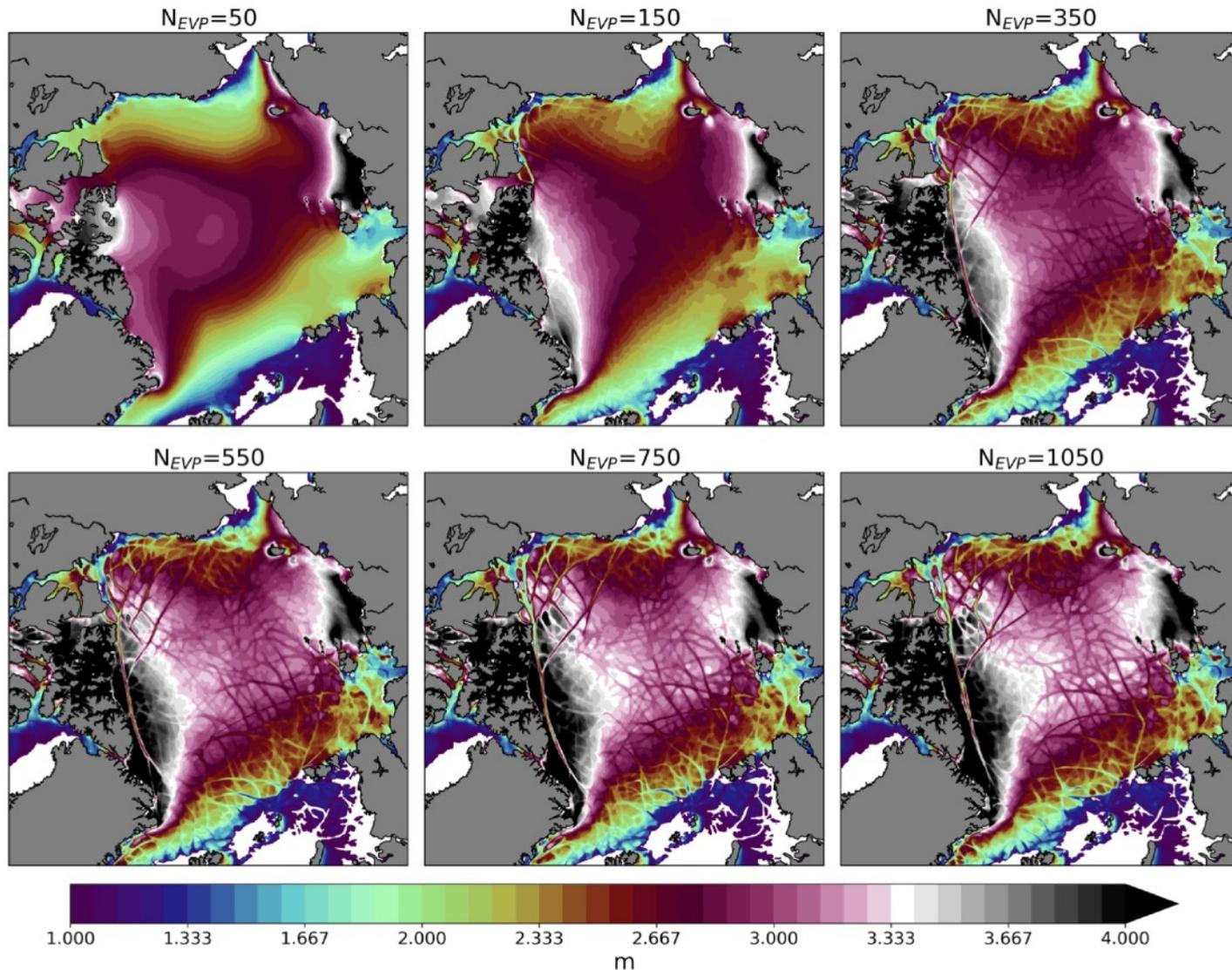
Resolution (1km)



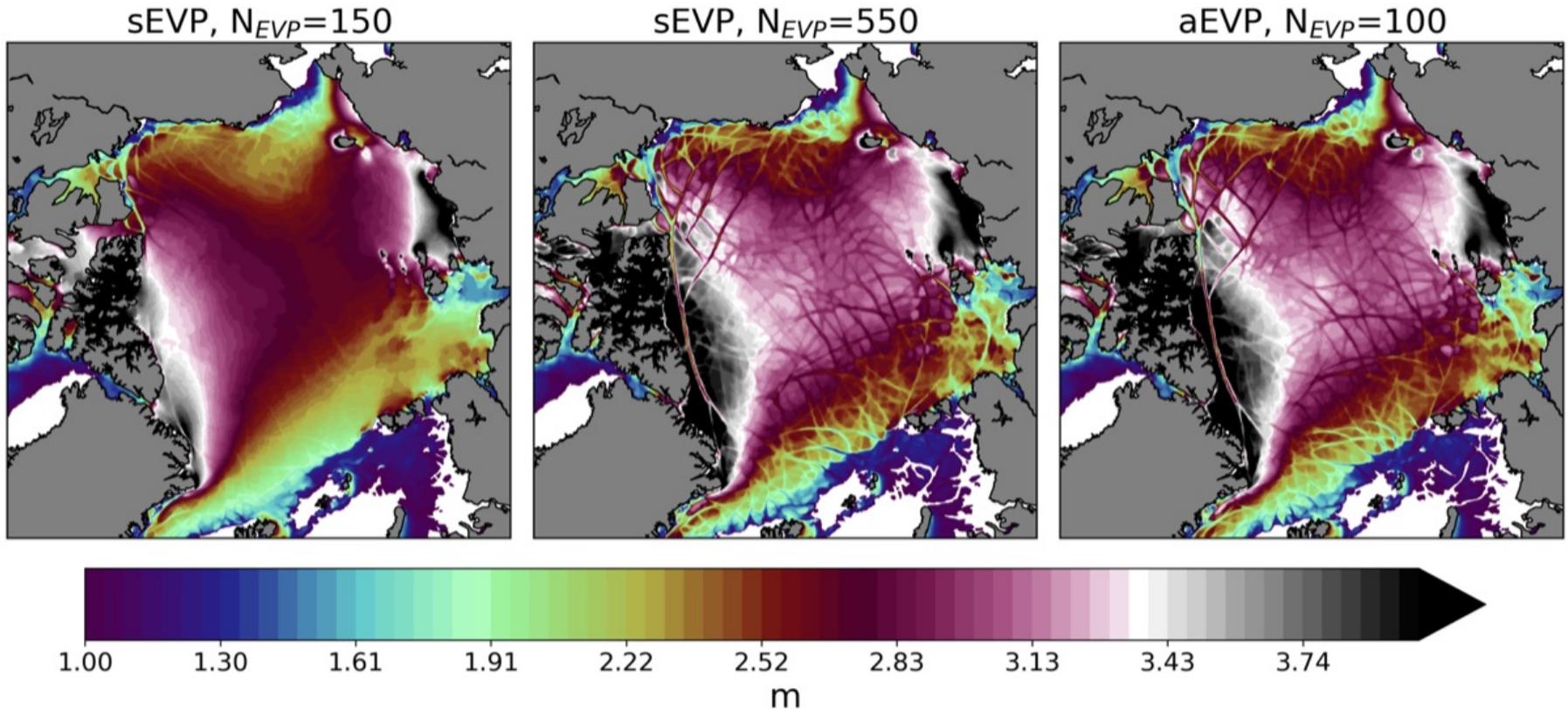
2012/05/27

Simulation: Menemenlis (JPL)
Graphics: Hutter (AWI)

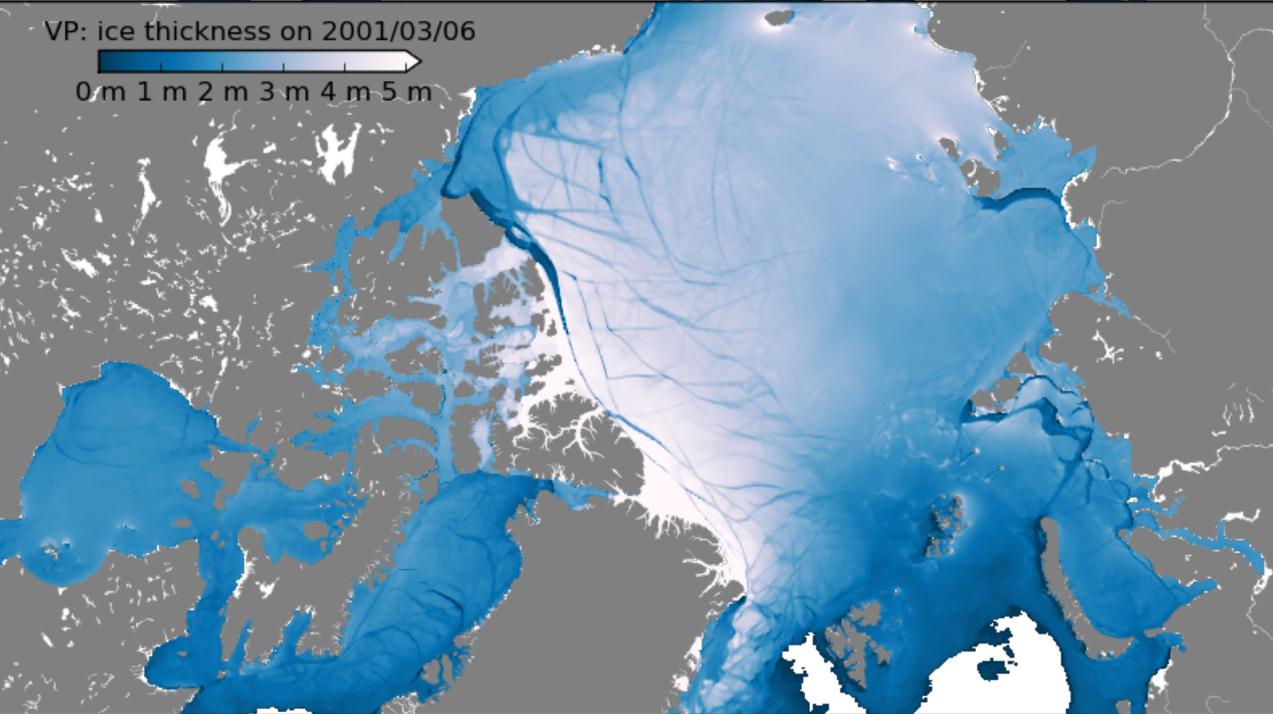
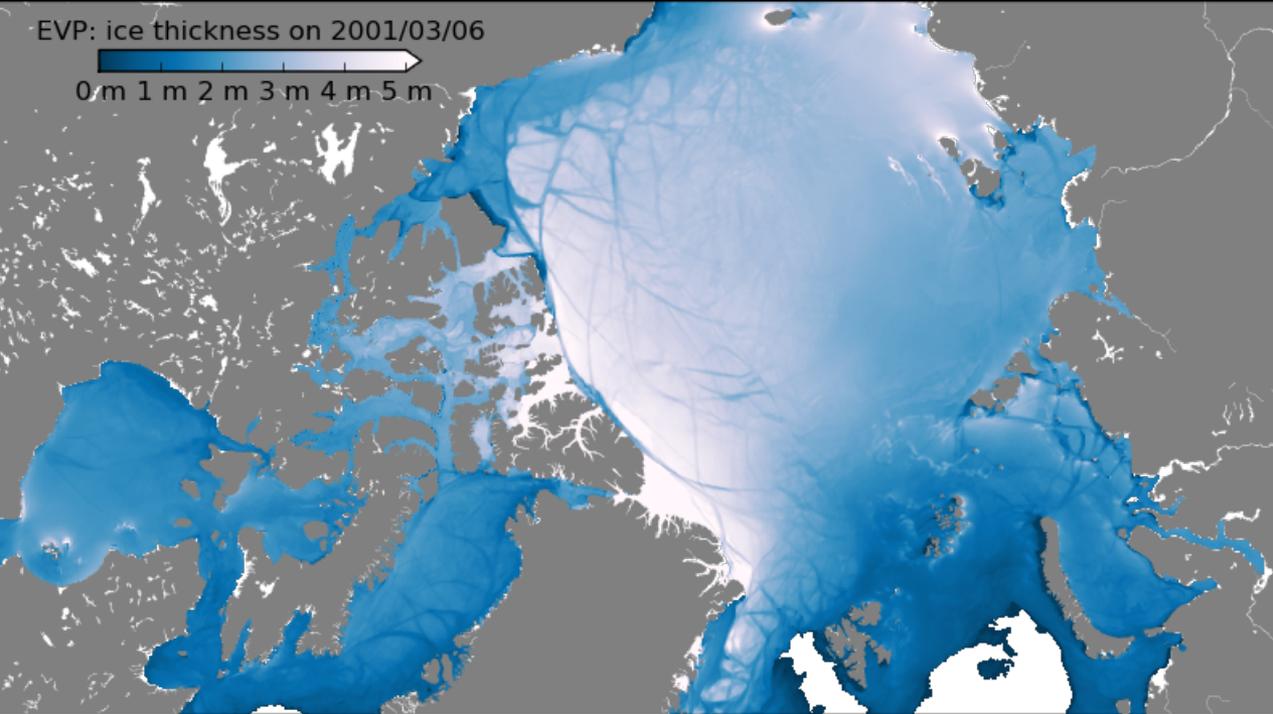
EVP “convergence” (in FESOM)



EVP “convergence” (in FESOM)



Koldunov et al., submitted manuscript, grid resolution ~ 4 km



Convergence to
VP solution:
ice thickness (m)
at 4.5 km grid
spacing

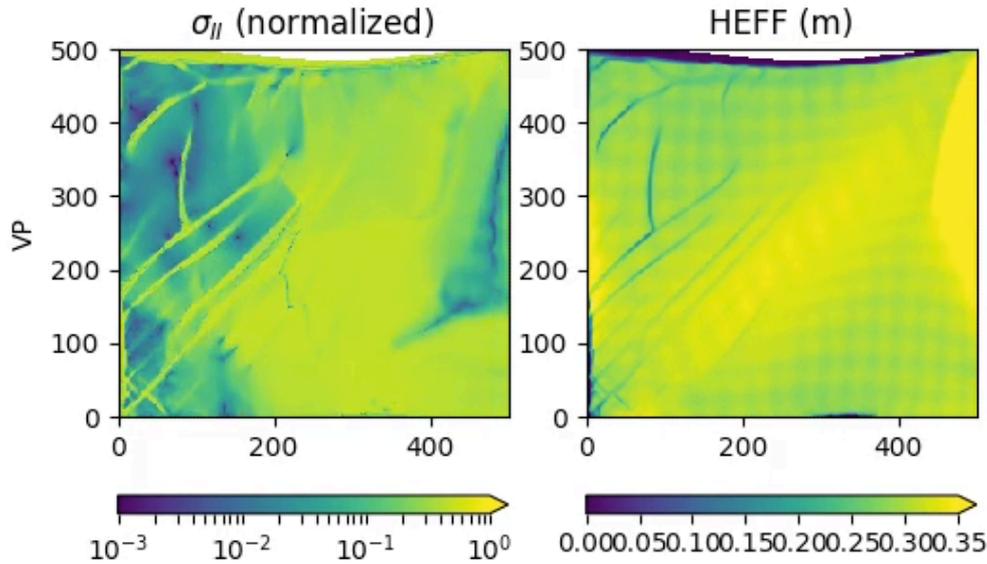
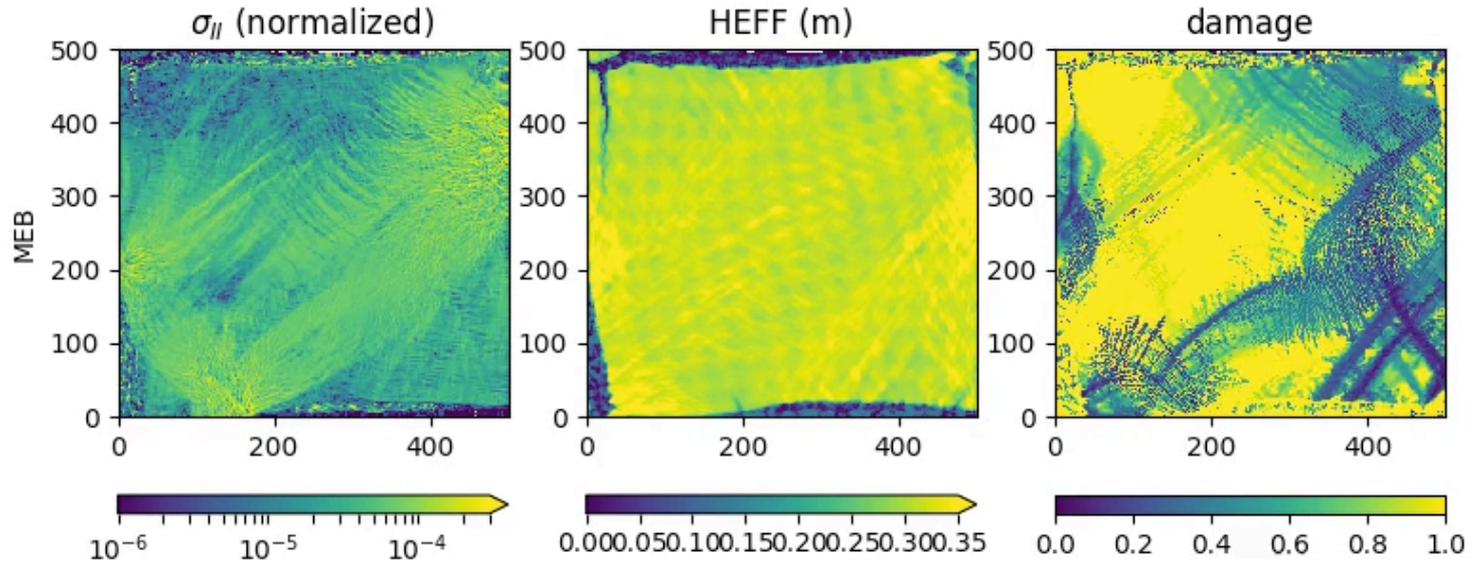
$$\alpha\beta \gg \gamma = \zeta \frac{(c\pi)^2}{A_c} \frac{\Delta t}{m}$$

stability
parameter
depends on grid
spacing and local
ice viscosity

$$\frac{1}{E} \frac{\partial \sigma}{\partial t} + \frac{1}{\lambda} \sigma = K : \dot{\epsilon}$$

- violation of a Mohr-Coulomb failure criterion determines a damage parameter
- damage parameter affects ice strength, elasticity
- (Girard et al 2011, Rampal et al 2015, Dansereau et al 2016)
- In the pipeline for the MITgcm

0001/01/08 09:00



not sure if I should show this

Summary

- Viscous Plastic rheology:
 - Picard solvers with LSR and Krylov solvers
 - JFNK solver
 - many EVP variants, especially stable EVP algorithms (Kimmritz et al. 2015, 2016)
- Maxwell Elasto-Brittle rheology (not quite, yet)
- all in the same code framework (no confounders in comparisons)
- main issues remain, especially at high resolution: convergence, stability vs. geophysical plausibility vs. time to solution