Research Article



Shape adaptation of beams (ID) and plates (2D) to maximise eigenfrequencies

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Simone Andresen^{1,2}, Laura M Lottes¹, Selina K Linnemann¹ and Reinhold Kienzler²

Abstract

Finding the optimal structural design to avoid resonance has been a goal for decades. While recent applied methods often result in using additional active systems or higher mass, structural adaptation enables to shift eigenfrequencies without adding weight. The aim of this study is to investigate the influence of the structural adaptation of a beam and a plate on its eigenfrequency change, while varying the height of the structural pre-deformation according to its mode shapes. Besides the maximisation of single eigenfrequencies, also the simultaneous increase of multiple eigenfrequencies is analysed. It is possible to almost exclusively raise the frequency of the targeted *i*-th mode shape (i = 1-5) of a beam, while the increase of the *i*-th plate mode shape frequencies are able to be significantly increased. In conclusion, the investigated, easy applicable method allows a strong eigenfrequency raise of axially constrained 1D and 2D structures by performing only small structural deformations without adding additional weight.

Keywords

Arches, beams, eigenfrequency maximisation, mode shape adaptation, parameter study, plates, shells

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Introduction

Finding the optimal structural design to avoid resonance has been a goal for decades as it is of high interest in many technical areas, for example, in the car industry to minimise the occurring noise and vibration during operation or in the space industry to avoid coupling with the control system.¹ Especially with recent demands for high-function structures including substantial weight reductions, the vibration design in consideration of dynamic characteristics has become extremely important.²

The dynamics of a structure can be modified either actively or passively. As active systems imply the use of additional devices, passive systems are of high interest. By altering the structures geometry or material, changes in the dynamics appear due to alterations of the stiffness or mass.³ In regard to lightweight structures, the change of a structure's geometry to shift its eigenfrequency obtains therefore high potential and has been subject of many studies. Most projects have been focusing on structural optimisations (sizing, topology and

²Bremen Institute for Mechanical Engineering, University of Bremen, Bremen, Germany

Corresponding author:

Simone Andresen, Bionic Lightweight Design and Functional Morphology, Alfred Wegener Institute Helmholtz Centre for Polar and Marine Research, Am Handelshafen 12, Bremerhaven 27570, Germany. Email: simone.andresen@awi.de

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¹Bionic Lightweight Design and Functional Morphology, Alfred Wegener Institute, Helmholtz Centre for Polar and Marine Research, Bremerhaven, Germany

shape optimisation) to find an optimum material distribution with the aim to maximise eigenfrequencies.

Numerical sizing optimisations of 1D and 2D structures under certain static or dynamic loading conditions to optimise specific, single eigenfrequencies have already been studied starting from the 70s, for example, conducting cross-section optimisation of beams^{4,5} or focussing on thickness optimisation of circular and rectangular plates.^{6,7} A regularised mathematical formulation of annular plates with the objective to maximise the 1st eigenfrequency has been presented.⁸ Similar, but taking dual optimisation problems into account, beams, frames and plates have been studied using a generalised steepest descent method to optimise the 1st eigenfrequency with constraints on the deflection while keeping minimum weight and varying the cross-section of beams or frames or the thickness of plates.^{9–11} A summary of literature concerning sizing optimisation to increase frequencies of different 1D and 2D structures was published.¹²

Next to the sizing optimisation, most focus has been put on topology optimisation that was first considered using the homogenisation design method with focus on increasing the 1st eigenfrequency of plates.¹³ The work was extended focusing on single and multiple eigenfrequencies as well as isotropic and composite plates.¹⁴

Other optimisation methods were utilised as well in order to maximise the 1st eigenfrequency or high-order eigenfrequencies, for example, the solid isotropic microstructures with penalisation method,^{15–17} the evolutionary structural optimisation method,^{18,19} the bidirectional evolutionary structural optimisation method²⁰ and the level set method.²¹ Focus points within the topology optimisation were ways to handle and maximise multiple eigenfrequencies^{16,22} as well as to create stiffener layout patterns that are engraved in or added to the plate to maximise eigenfrequencies.^{23,24}

Beside topology optimisation, also shape optimisation has been frequently used. The natural frequency of various plates and shells have been maximised by altering the shape and thickness.²⁵ Also the vibration design of stiffened thin-walled shell structures have been optimised to increase eigenvalues or reach specified eigenvalues.² Parametric investigations regarding topography optimisation have been applied examining the optimal distribution of spherical dimples or cylindrical beads to optimise the natural frequency of plate-like structures.²⁶ Shape optimisations to increase single eigenfrequencies of plate membranes have been performed,²⁷ also considering constraints on the volume.²⁸ An optimum bead orientation was investigated for the maximisation of the first five eigenfrequencies of a cantilever plate²⁹ and continued aiming at a 1st eigenfrequency increase of 3D shell/plate structures by implementing an energy-based method.²³ An optimum distribution of sinusoidal and trapezoidal corrugations was studied in order to shift the fundamental frequency to maximum values.³⁰

It can be concluded that many different approaches have already been studied to maximise the eigenfrequencies of different structures. However, structural optimisations like topology and shape optimisations are often time consuming and require high computational effort. In addition, the manufacturing of the optimised structures is often restricted to additive manufacturing, which has become more promising in the last years, but still lacks on reproducibility, costs and production time. Therefore, all mentioned studies led to good results, however, might not be applied in big scale or for low-cost parts. A current demand for an efficient frequency optimisation method is therefore still present.

A new method of shape optimisation to maximise eigenfrequencies of beams has recently been published, which obtains the potential to be such an efficient optimisation technique.³ Within the study, it was concluded that it is possible to influence the *i*-th mode shape frequency of an axially constrained beam by changing the beam shape according to its *i*-th eigenvectors (mode shapes), while all additional frequencies of lower-order mode shapes remain unaffected. Parametric studies varying the height of the pre-deformation of the beam were performed, which indicated a strong increase of a targeted mode shape eigenfrequency by only applying minor pre-deformations to the beam.³ A specific mode shape frequency increase can therefore be obtained by simply adapting the beam's shape according to its corresponding mode shape. First indications that the method also applies for plates was shown in the application of a cab floor.31

The advantages of this shape adaptation methodology are that no special manufacturing methods are needed as well as that the proposed methodology implies relatively small time and computational efforts compared to other optimisation procedures as solely the mode shapes of the structure have to be known. However, the improved loadbearing capacity (arch or vault effect) due to the pre-deformation, which leads to eigenfrequency increases, requires axial constraints.

This study therefore focuses on an extension of the work of Da Silva and Nicoletti³ by further investigating the potential of the method regarding higher predeformations of simply supported beams. In addition, the parametric studies were applied to squared plates. The study objective was to investigate the potential of the method not only to increase single eigenfrequencies and frequencies of specific mode shapes of beams and plates by adapting the shape according to one single mode shape, but also to study the possibility to use the proposed optimisation method for the maximisation of multiple eigenfrequencies by applying several mode shapes.

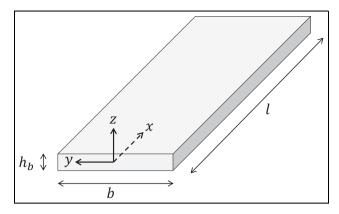


Figure 1. Slender reference beam with coordinate system.

Material and method

A slender beam, which has already been studied,³ and a squared plate were pre-deformed according to their mode shapes conducting parametric studies to investigate the impact of the pre-deformations on the eigenfrequencies.

All parametric and algorithm-based constructions were performed using the software Rhinoceros (version 6 SR10, Robert McNeel & Associates) and its plug-In Grasshopper[®] (version 1.0.0007, Robert McNeel & Associates). The additional Grasshopper[®]-based module ELISE (version 1.0.36, ELISE GmbH, www.elise.de) allowed the construction of the entire design and simulation process and the connection to the solver OptiStruct (Altair[®] HyperWorks[®] Version 2017) used to obtain the numerical results.

Slender beam

The beam geometry and material properties were defined in analogy to Da Silva and Nicoletti³ leading to a 600 \times 30 \times 3 mm beam (Figure 1) made out of aluminium (Young's modulus: 69,000 MPa, density: $2.688 \cdot 10^{-9}$ tmm⁻³, Poisson's ratio: 0.34). Axial constraints are inevitable to increase eigenfrequencies by performing the here analysed shape adaptation.³ Consequently, the beam was simply supported at both ends implying a restriction of all degrees of freedom except for the rotation around the *y* axis.

A modal analysis to calculate the first six eigenfrequencies and eigenmodes was conducted. The adequate mesh was obtained in a mesh convergence study varying the total number of nodes from 3 to 31, based on which CBEAM elements were defined. A sufficient mesh fineness was reached as soon as the results differed less than 5% from the results of the following three finer meshes. The results were evaluated on the 1st and the 6th eigenfrequency.

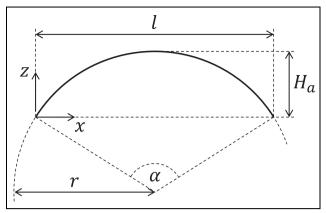


Figure 2. Dimensions of a circular arch.

The numerically obtained results of the straight reference beam and the beam deformed according to the 1st mode shape were compared to analytically obtained eigenfrequencies to confirm the plausibility of the numerical simulations. The eigenfrequencies $f_{i, beam}$ of a simply supported Bernoulli beam characterised by a Young's modulus *E*, a material density ρ , a length *l*, a rectangular cross-section area *A* and a second moment of inertia *I*, both depending on the beam width *b* and the beam height h_b , are defined as follows³²:

$$f_{i, beam} = \frac{1}{2\pi} \frac{(i\pi)^2}{l^2} \sqrt{\frac{EI}{A\rho}}; i = 1, 2, 3, \dots$$
(1)

$$A = b \cdot h_b \tag{2}$$

$$I = \frac{b \cdot h_b^3}{12} \tag{3}$$

The *i*-th bending mode shape in the *xz* plane $W_{i, beam}$ can be described with the following equation³²:

$$W_{i, beam}(x) = \sin\left(\frac{i\pi x}{l}\right); i = 1, 2, 3, ...$$
 (4)

Formulae to calculate the 1st and 2nd eigenfrequency of a circular arch have been published³³ and may serve as approximation for sinusoidal beams with small l/r ratios (Figure 2). The first two eigenfrequencies $f_{1,arch}$ and $f_{2,arch}$ can be calculated as follows:

$$f_{1,arch} = \frac{1}{2\pi} \sqrt{\frac{1}{r^4} \left[0.82 \left(\frac{r}{k}\right)^2 + \left(\frac{\pi^2}{\alpha^2} - 1\right)^2 \right]} \sqrt{\frac{EI}{A\rho}} \quad (5)$$

$$f_{2,arch} = \frac{1}{2\pi} \sqrt{\frac{1}{r^4 \alpha^4} \frac{\alpha^4 - 8\pi^2 \alpha^2 + 16\pi^4}{1 + 0.075 \frac{\alpha^2}{\pi^2}}} \sqrt{\frac{EI}{A\rho}}$$
(6)

with the radius r of the arch curvature depending on the arch span length, which is equal to the undeformed

beam length *l*, the arch height H_a , the central angle α (radian), and the slenderness ratio *k* depending on a rectangular cross-section:

$$r = \frac{1}{2H_a} \left[\left(\frac{l}{2} \right)^2 + H_a^2 \right] \tag{7}$$

$$\alpha = 2 \cdot \sin^{-1} \left(\frac{l}{2r} \right) \tag{8}$$

$$k = \frac{I}{A} = \frac{h_b^2}{12} \tag{9}$$

Maximisation of a specific eigenfrequency or mode shape frequency. A pre-deformation according to the 1st, 2nd, 3rd, 4th and 5th eigenvectors was applied to the analysed beam. In analogy to Da Silva and Nicoletti,³ a maximum relative pre-deformation δ was defined as the quotient of the maximum pre-deformation of the beam δ_{max} and the beam height h_b :

$$\delta = \frac{\delta_{max}}{h_b} \tag{10}$$

In parametric studies, δ_{max} was varied to create maximum relative pre-deformations from 0.0 to 5.0 with step sizes of 0.5 and from 5.0 to 20.0 with step sizes of 5.0. The 5th mode shape showed a first transversal (i.e. out-of-plane) bending mode. Consequently, the height considered for the calculation of the maximum relative pre-deformation δ was the beam width *b*. During the pre-deformation according to mode 5, δ_{max} was varied up to 180 mm, which corresponded to a maximum relative pre-deformation of 6.0.

It may be mentioned that, although the predeformation was quite large, the analyses were still linear. The mechanical system, however, was changed considerably. Instead of beams (and plates), arches (and vaults or shells) were investigated. In linear beam theory, the natural eigenfrequencies do not depend on the axial constraint, whereas the axial constraint (hindered horizontal displacements at the ends) becomes essential for arches and shells as the membrane loadbearing capacity increases and so do the eigenfrequencies. Without the horizontal constraint, the frequency change is marginal.³ Since the span length l was kept constant, the length of the arch changed with increasing pre-deformation. As the aim was to alter eigenfrequencies only due to the structural deformation without changing the mass (and also boundary conditions and material properties), the beam width b was adapted in all calculations to have a constant beam mass of 145 g. If a variation of mass would have been permitted, it would not have been possible to clearly state that the eigenfrequency increases were due to the structural deformations, because mass change would also have strongly manipulated the eigenfrequencies.

During the analyses, the first six eigenfrequencies and the mode shapes were recorded. The resulting eigenfrequencies were always ordered by their value. Consequently, the frequency increase of a specific mode shape can lead to an alternation of the mode shape order. Thus, if the mode shape that was adapted to the beam could not be found within the first six eigenmodes, higher-order mode shapes were searched for the corresponding mode shape.

The results of the parametric studies were compared to the reference beam. The obtained eigenfrequency increase Δf was calculated based on the *i*-th eigenfrequency and the *i*-th eigenfrequency of the reference beam $f_{i, ref}$:

$$\Delta f = \frac{f_i - f_{i,ref}}{f_{i,ref}} \cdot 100\%; i = 1, 2, 3, \dots$$
(11)

Equation (11) was also utilised to obtain the increase of specific mode shape frequencies.

Maximisation of multiple eigenfrequencies. The reference beam was pre-deformed according to linear combinations of different mode shapes to analyse whether a simultaneous maximisation of different eigenfrequencies was possible. The effect of linear combinations of mode 1 and 2, mode 1, 2 and 3, and mode 1, 2, 3 and 4 of the reference beam were tested for maximum relative pre-deformations of 1.0, 3.0, and 5.0. For each beam node, the normalised eigenvectors of the considered mode shapes were added in equal parts. The resulting sum was normalised to obtain a maximum amplitude of 1, which was multiplied by the maximum relative pre-deformation δ that was analysed.

Additionally, weighted linear combinations of mode 1 and mode 2 were investigated considering also maximum relative pre-deformations of 1.0, 3.0, and 5.0. The normalised eigenvectors of mode 1 and mode 2 were multiplied with the factors c and (1 - c), respectively, varying c from 0.0 to 1.0 in step sizes of 0.1. Also here, the sum of the resulting values was normalised and subsequently multiplied with the maximum relative pre-deformation δ .

Squared plate

The considered plate was characterised by an edge length *a* of 100 mm and a constant thickness h_p of 2 mm (Figure 3). The material properties were set analogue to the slender beam resulting into a plate mass of 53.76 g. Concerning the boundary conditions, all degrees of freedom were restricted except for rotations around the *x* axis for the edges parallel to the *x* axis, and for

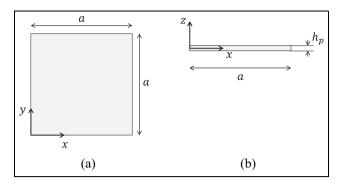


Figure 3. Top view (a) and front view (b) of the squared reference plate with coordinate system.

rotations around the *y* axis for the edges parallel to the *y* axis.

Similar to the investigated beam, a shell mesh convergence study was conducted varying the number of elements per plate edge from 4 to 200 to identify a sufficient element size. A modal analysis was performed in order to obtain the first six eigenfrequencies and the corresponding mode shapes.

As the thickness to length ratio of the reference plate was smaller than 10%, the plate could be considered as thin and the classic plate theory of Kirchhoff was applied to calculate the eigenfrequencies analytically.³⁴ The eigenfrequencies $f_{ij, plate}$ of the simply supported, squared plate depend on the plate edge length *a*, the material density ρ , the plate height h_p and the plate stiffness *K*, which is based on Young's modulus *E*, the plate height and Poisson's ratio ν^{32} :

$$f_{ij,plate} = \frac{\pi}{2a^2} \left(i^2 + j^2 \right) \sqrt{\frac{K}{\rho h_p}}; i, j = 1, 2, 3, \dots$$
(12)

$$K = \frac{Eh_p^3}{12(1-\nu^2)}$$
(13)

The mode shapes $W_{ij, plate}$ of the squared plate were obtained using the following equation:

$$W_{ij,plate}(x,y) = \sin \frac{i\pi x}{a} \sin \frac{j\pi y}{a}; i,j = 1, 2, 3, ...$$
 (14)

The reference plate pre-deformed according to mode 1 was alike a spherical shell. Thus, shallow spherical shell theory was used to validate the numerical calculations. Shallow spherical shell theory implies thin shells.³⁴ These thin shells are characterised by a shell wall length and width of less than 10% of the shell radius, a constant shell thickness and a shell rise of less than about 1/8th of its lateral dimension. During vibration, these shells deform primarily perpendicular to the shell surface. The eigenfrequencies of a shallow spherical shell $f_{ij,SSS}^2$ can be obtained using the following

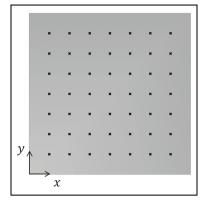


Figure 4. Regular distribution of points on the plate which were used for pre-deforming the plate according to its mode shapes.

equations, in which *R* represents the plate curvature radius and H_{SSS} the height of the spherical shell³⁵:

$$f_{ij,SSS}^{2} = f_{ij,plate}^{2} + \frac{1}{4\pi^{2}} \frac{E}{\rho R^{2}}; i,j = 1, 2, 3, ...$$
 (15)

$$R = \frac{1}{2H_{SSS}} \left[\left(\frac{a}{2}\right)^2 + H_{SSS}^2 \right]$$
(16)

Maximisation of a specific eigenfrequency or mode shape frequency. The plate was pre-deformed according to its first four mode shapes by calculating the z values of 49 regular distributed points on the plate (Figure 4) using formula 14. The surface in between the points was interpolated.

Analogue to the slender beam, the maximum relative pre-deformation δ was defined as the quotient between the maximum pre-deformation of the plate varied from 0.0 to 60 mm and the plate thickness. This time, the plate thickness was adapted to keep a constant mass. The first six eigenfrequencies and the corresponding mode shapes were recorded. If the mode shape that was adapted to the plate could not be found within the first six eigenmodes, higher-order mode shapes were searched for maximum relative pre-deformations of 3.0 and 5.0 to find the corresponding mode shape. The eigenfrequency deviation compared to the reference plate was calculated using formula (11).

To investigate the impact of the boundary conditions on the eigenvector approach, the boundary conditions of the reference plate were varied in three studies involving (1) all four edges clamped, (2) two opposite edges clamped and the other two edges free and (3) one edge clamped and the other three edges free. The plates were pre-deformed according to their corresponding first four mode shapes. The maximum predeformations were varied in the same way as for the

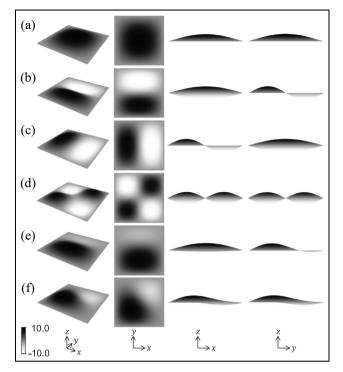


Figure 5. Squared plate pre-deformed according to the 1st (a), 2nd (b), 3rd (c) and 4th (d) single bending mode shape. (e) shows the pre-deformation according to a linear combination of the first two mode shapes (50% mode I and 50% mode 2) and (f) the linear combination of the first three mode shapes (33% mode 1, 33% mode 2, and 33% mode 3). The pre-deformed plates are displayed in four different views for δ = 5.0 and the grey shades represent the maximum pre-deformation (mm).

simply supported plate and their impact on the first six eigenfrequencies was analysed.

Maximisation of multiple eigenfrequencies. According to the investigated beam, linear combinations of mode shapes were applied to the reference plate. As already described in the corresponding paragraph about the beam, mode 1 and 2 and mode 1, 2, and 3 were equally combined. Also here, weighted combinations of mode 1 and mode 2 were investigated. Both analyses were performed for maximum relative pre-deformations of 1.0, 3.0 and 5.0. The mesh properties and boundary conditions coincided with the previous investigations.

Figure 5 exemplarily shows the squared plate predeformed according to the first four mode shapes and a linear combination of the first two and the first three mode shapes. Here, a maximum pre-deformation of 10 mm was considered.

Results

Slender beam

A beam mesh consisting of 24 elements (25 nodes) was chosen to allow the generation of complex beam

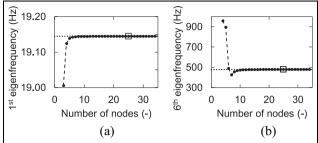


Figure 6. Results of the beam mesh study involving the 1st (a) and the 6th (b) eigenfrequency depending on the number of nodes. The dotted lines represent the analytically obtained results using Bernoulli beam theory. The framed data points indicate the chosen mesh properties.

deformations throughout all simulations, even though the mesh study results showed that the output values had already converged with a coarser mesh and coincided with the analytically obtained values (Figure 6). The first six mode shapes of the reference beam are shown in Figure 7. The 1st, 2nd, 3rd, 4th and 6th mode shapes represented the 1st, 2nd, 3rd, 4th and 5th bending mode shapes in the *xz* plane, while the 5th mode shape was the 1st bending mode in the *xy* plane (1st out-of-plane bending mode). Regarding the eigenfrequencies of the circular arch, the theoretical and numerical frequencies of the 1st and the 2nd mode coincided very well (Figure 8).

Maximisation of a specific eigenfrequency or mode shape frequency. Shaping the beam according to the 1st, 2nd, 3rd and 4th bending mode shape resulted in a strong increase of the eigenfrequency connected to the corresponding mode shape. Figure 9 shows exemplarily the results for shaping the beam according to the 3rd and 4th bending mode shape, while the results for the 1st and 2nd bending mode shape adaptation are attached in the Supplemental Material. Lower-order frequencies remained constant, also for higher maximum relative pre-deformations. However, higher-order frequencies mostly decreased slightly with increasing maximum relative pre-deformation. Especially the frequency of the 5th mode shape (i.e. 1st bending mode in the xz plane) strongly decreased with increasing maximum relative pre-deformation for all beam pre-deformations. Regarding the 5th mode adaptation, however, the corresponding eigenfrequency first increased and later decreased with increasing maximum relative predeformation (Figure 10). The other recorded eigenfrequencies remained constant or decreased slightly.

In summary, a maximum relative pre-deformation of 5.0 already resulted in a frequency increase of a specific mode shape of more than 1000% for a shape adaptation of the beam according to the 1st, 2nd, 3rd

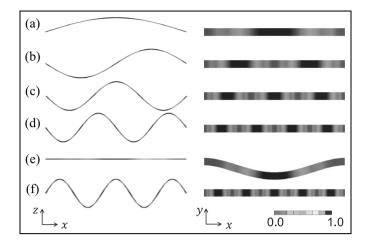


Figure 7. First six bending mode shapes of the analysed beam in the *xz* and the *xy* plane showing the absolute normalised vibration amplitude.

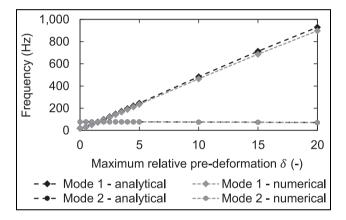


Figure 8. Ist and 2nd eigenfrequency of the beam predeformed according to the bending mode 1 for different maximum relative pre-deformations. The analytically obtained values are based on formulae for circular arches published by Den Hartog.³³

and 4th bending mode shape. Maximum relative predeformation of 20.0 generated frequency increases of more than 4000% for pre-deformations according to mode 1 and 2 and of more than 2000% for mode 3 and 4. However, shaping the beam according to the 5th mode shape increased the corresponding eigenfrequency only by 164% for a maximum relative predeformation of 5.0.

The investigated methods to increase specific eigenfrequencies rose the targeted eigenfrequencies in almost all analyses, as shown in Table 1(a). Only when shaping the beam according to the 4th mode shape, a decrease of all six eigenfrequencies with strongest decrease for the 4th eigenfrequency was present. In addition, the not targeted eigenfrequencies were also altered (i.e. decreased or increased) by 8% to 28% applying maximum relative pre-deformations of 3.0.

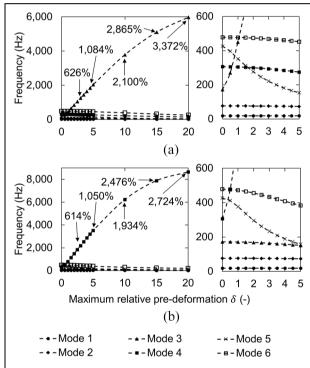


Figure 9. Frequencies of the first six bending mode shapes of the slender beam depending on the maximum relative predeformation according to mode 3 (a) and mode 4 (b). For some data points, the frequency increase of the 3rd and 4th bending mode shape compared to the reference beam is given. A magnified view of the lower right corner of the diagram involving small pre-deformations δ of 0–5 and low frequencies of 0–600 Hz is given on the right-hand side of the figure.

Comparing the obtained frequency changes corresponding to a specific mode, all results showed highest frequency increase for the targeted mode shape (Table 1(b)). The increase varied from 614% for the frequency

Table 1. (a) Eigenfrequency changes compared to the reference beam due to the eigenvector approach, which was applied to the slender beam considering a maximum relative pre-deformation of 3.0. The alterations of the frequencies connected to specific mode shapes (M) are shown in (b). The average deviations of not targeted eigenfrequencies are given. All targeted eigenfrequencies or mode shapes are shown in bold numbers.

(a)		Max. f _I	Max. f ₂	Max. f ₃	Max. f ₄
	fı	299 %	0%	-1%	-1%
	f ₂	85%	122%	-1%	-1%
	f3	0%	76%	28 %	-5%
	f ₄	0%	36%	-4%	-24%
	fs	-5%	11%	10%	-23%
	f6	0%	15%	42%	9%
	Deviation of not targeted eigen frequencies	18%	28%	12%	8%
(b)		Max. f _l	Max. f ₂	Max. f3	Max. f4
	МІ	641%	0%	-1%	-1%
	M2	0%	636%	-1%	-1%
	M3	0%	-1%	626 %	-5%
	M4	0%	-1%	-4%	614%
	M5	-5%	-3%	48%	-45%
	M6	0%	-1%	-2%	9%
	Deviation of not targeted eigenfrequencies	1%	1%	11%	12%

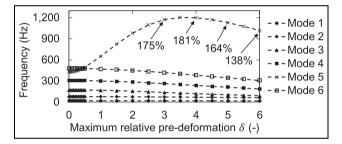


Figure 10. Frequencies of the first six bending mode shapes of the slender beam depending on the maximum relative predeformation according to mode 5. For some data points, the frequency increase of the 1st bending mode shape in the *xy* plane compared to the reference beam is given.

of the 4th bending mode shape to 641% for the frequency of the 1st bending mode shape for the eigenvector approach considering a maximum relative predeformation of 3.0.

The not targeted mode shape frequencies changed in average between 1% and 12%.

Maximisation of multiple eigenfrequencies. Pre-deforming the reference beam according to a single mode shape led to high frequency increases of the considered mode shape, whereas the frequencies of the other mode shapes only differed slightly (Table 2). This trend was also observed for combinations of the different mode shapes. By shaping the reference beam according to a combination of the bending modes 1 and 2, only the frequencies of both considered mode shapes increased, while the frequencies of the other mode shapes remained almost constant. This also applied for the combinations of mode 1, 2, and 3 and of mode 1, 2, 3, and 4, however, the former also resulted in a rise of the 4th bending mode frequency for a maximum relative pre-deformation of 1.0. In all analysed mode combinations, the frequencies of the highest order mode shapes always increased most.

Pre-deforming the reference beam according to weighted combinations of the bending modes 1 and 2 led to different results for varying relative maximum pre-deformations (Figure 11). For $\delta = 1.0$, the highest sum of the 1st and 2nd eigenfrequency was obtained by pre-deforming the beam according to 20% of bending mode 1 and 80% of bending mode 2, for $\delta = 3.0$ by 80% of bending mode 1 and 20% of bending mode 2 and for $\delta = 5.0$ by 100% of bending mode 1 and 0% of bending mode 2.

Squared plate

The mesh size study indicated the sufficiency of an element edge size of 16.6 mm to obtain correct values for the 1st and 6th eigenfrequency (Figure 12). Also in this case, a finer mesh with an element edge size of 1 mm and a total of 10,000 CQUAD4 elements was used to picture the mode shapes correctly. The bending mode shapes of the reference plate are demonstrated in Figure 13. The comparison between the numerical and analytical results for the plate shaped according to the 1st bending mode shape showed good coincidence (Figure 14).

Maximisation of a specific eigenfrequency or mode shape frequency. The pre-deformations of the reference plate

δ	100% Mode	e l			100% Moo	de 2					
	MI	M2	M3	M4	MI	M2	M3	M4			
I	165%	0%	0%	0%	0%	164%	0%	0%			
3	641%	0%	0%	0%	0%	636%	-1%	-1%			
5	1,126%	0%	0%	0%	-1%	1,110%	-3%	-2%			
δ	100% Mode 3			100% Mode 4							
	MI	M2	M3	M4	MI	M2	M3	M4			
I	0%	0%	162%	0%	0%	0%	-1%	159%			
3	-1%	-1%	626%	-4%	-1%	-1%	-5%	614%			
5	-2%	-4%	1,084%	-10%	-3%	-4%	-12%	1,050%			
δ	50% Mode I, 50% Mode 2				33% Mode	e 1, 33% Mode 2, a	nd 33% Mode 3				
	MI	M2	M3	M4	MI	M2	M3	M4			
I	27%	73%	0%	0%	16%	21%	55%	15%			
3	36%	338%	0%	0%	21%	29%	277%	-1%			
	37%	616%	-1%	-1%	20%	29%	509%	-3%			
δ	25% Mode 1, 25% Mode 2, 25% Mode 3, and 25% Mode 4										
	MI		M2		M3		M4				
I	10%		12%		15%		32%				
3	14%		18%		23%		190%				
5	14%		17%		24%		362%				

Table 2. Frequency deviations of the first four bending mode shapes (M1 to M4) compared to the reference beam for different maximum relative pre-deformations δ according to the 1st, 2nd, 3rd and 4th single mode shapes and combinations of these four mode shapes.

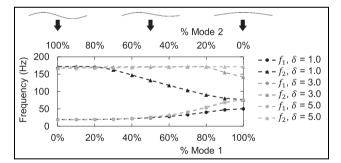


Figure 11. Ist and 2nd eigenfrequency of the beam shaped with a maximum relative pre-deformation δ of 1.0, 3.0 and 5.0 according to weighted combinations of the bending modes 1 and 2. Three beam shapes are exemplarily shown for $\delta = 10.0$. The trends of the 1st eigenfrequencies for $\delta = 3.0$ and $\delta = 5.0$ almost coincided.

implied adaptations of the plate thickness to ensure a constant mass. In the Supplemental Material, a table with the decreased plate thicknesses dependent on the maximum pre-deformations δ_{max} is attached. It has to be noted that due to the plate thickness adaptation, the maximum relative pre-deformations altered.

Pre-deforming the reference plate according to its 1st bending mode shape resulted at first in an increase

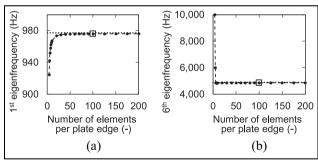


Figure 12. Results of the plate mesh study involving the 1st (a) and the 6th (b) eigenfrequency depending on the number of elements per plate edge. The dotted lines represent the analytically obtained results using the classic plate theory of Kirchhoff. The framed data points indicate the chosen mesh properties.

of all considered eigenfrequencies (Figure 15(a)), which decreased after a maximum relative pre-deformation of 14.7 for the 1st and 6th eigenfrequency and 10.8 for the remaining eigenfrequencies. At its highest peak, the 1st eigenfrequency increased by 606% compared to the reference plate. Figure 15(b) shows how the first six bending mode shapes changed with increasing predeformation of the plate. The 1st bending mode shape

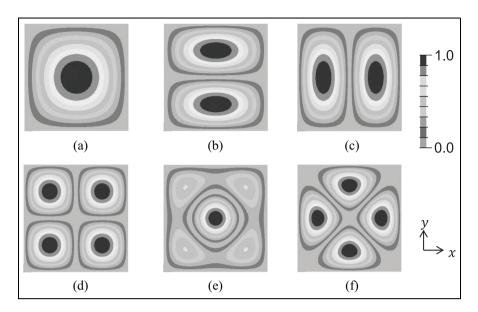


Figure 13. 1st (a), 2nd (b), 3rd (c), 4th (d), 5th (e) and 6th (f) bending mode shapes of the reference plate showing the absolute normalised vibration amplitude.

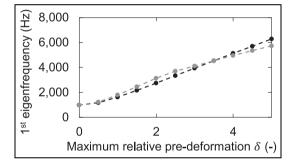


Figure 14. Analytically (black) and numerically (grey) obtained Ist eigenfrequency of the reference plate shaped according to the 1st bending mode shape depending on different maximum relative pre-deformations. The analytical solutions were calculated using classic plate theory of Kirchhoff for $\delta = 0.0$ (equation (12)) and shallow spherical shell theory for $\delta > 0.0$ (equation (15)).

of the reference plate disappeared with maximum relative pre-deformations higher than 4.1 and could also not be found within the first 50 mode shapes. Thus, the frequency of the 1st bending mode shape was increased by 488% at $\delta = 4.1$. For a maximum relative predeformation of 5.1, the 1st mode shape of the predeformed plate coincided with the 6th bending mode shape of the reference plate. The 2nd, 3rd, and 4th bending mode shapes were still present up to a maximum relative pre-deformation of 5.1.

Shaping the reference plate according to its 2nd or 3rd bending mode shape led to the same results (Figure 16). Similar to the plate adaptation according to its 1st bending mode shape, all analysed eigenfrequencies first

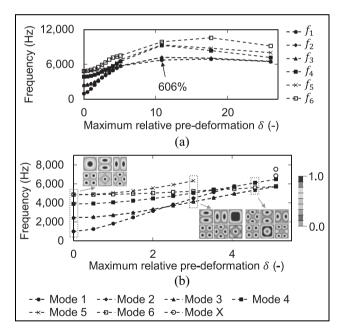


Figure 15. Eigenfrequencies (a) and frequencies of the different mode shapes (b) of the squared plate pre-deformed according to its 1st bending mode shape depending on different maximum relative pre-deformations. The highest 1st eigenfrequency increase is given. In (b), mode shapes which are not among the first six bending mode shapes of the reference plate are named 'Mode X'. For $\delta = 0.0$, 3.0 and 5.1, the mode shapes showing the absolute normalised vibration amplitude are displayed.

increased and then decreased with increasing maximum relative pre-deformation. At their highest peaks at a maximum relative pre-deformation of 11.6 for the 2nd and at 5.2 for the 3rd eigenfrequency, the

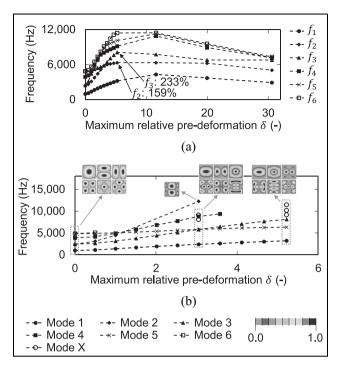


Figure 16. Eigenfrequencies (a) and frequencies of the different mode shapes (b) of the squared plate pre-deformed according to its 2nd or 3rd bending mode shape depending on different maximum relative pre-deformations. For both the 2nd and the 3rd eigenfrequency, the highest increases are given. In (b), mode shapes which are not among the first six bending mode shapes of the reference plate are named 'Mode X'. For $\delta = 0.0, 3.1$ and 5.2, the mode shapes showing the absolute normalised vibration amplitude are displayed.

eigenfrequencies increased by 159% and 233%, respectively. Regarding the mode shapes, only the 1st bending mode shape did not change order until a maximum relative pre-deformation of 5.2, while the frequencies of the other studied mode shapes exceeded each other and thus led to alterations in the mode shape order. The 2nd bending mode shape was found among the first 50 mode shapes up to a maximum relative predeformation of 3.1 with 12,237 Hz, whereas the 3rd bending mode shape could be identified up to a maximum relative pre-deformation of 5.2, however, its shape changed slightly. Thus, the frequency of the 2nd bending mode shape increased by 402% at $\delta = 3.1$ and the frequency of the 3rd bending mode shape by 233% at $\delta = 5.2$.

A plate shape adaptation according to the 4th bending mode shape resulted into similar eigenfrequency plots as for the previous shape adaptations involving first the eigenfrequency increase and with increasing maximum relative pre-deformations a decrease of all eigenfrequencies (Figure 17). The 4th eigenfrequency reached its highest peak at a maximum relative predeformation of 12.1 with an eigenfrequency increase of

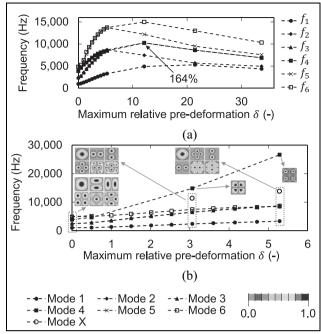


Figure 17. Eigenfrequencies (a) and frequencies of the different mode shapes (b) of the squared plate pre-deformed according to its 4th bending mode shape depending on different maximum relative pre-deformations. The highest 4th eigenfrequency increase is given. In (b), mode shapes which are not among the first six bending mode shapes of the reference plate are named 'Mode X'. For $\delta = 0.0$, 3.1 and 5.3, the mode shapes showing the absolute normalised vibration amplitude are displayed.

164% compared to the reference plate. The 4th bending mode shape was found within the first 50 mode shapes up to a maximum relative pre-deformation of 5.3 and its frequency reached 26,612 Hz resulting in a frequency increase of 584% compared to the reference plate.

The variation of boundary conditions showed that the eigenfrequency increase was dependent on the boundary conditions (Figure 18). While the highest 1st eigenfrequency increase of the fully clamped plate due to pre-deformation according to the bending mode 1 obtained about half of the value as for the simply supported plate, clamping of only two plate sides led to lower eigenfrequency increases. Moreover, the 1st eigenfrequency of the cantilever plate did not even rise. A similar trend could be observed for shaping the plates according to the 2nd, 3rd and 4th bending mode shapes, as seen in the Supplemental Material.

In summary, the eigenvector approach resulted in high eigenfrequency increases by pre-deforming a plate vertically of maximum 10% of the plate edge length (Table 3). The applied plate pre-deformations led not only to an increase of the targeted eigenfrequency, but also to an increase of all other eigenfrequencies of 87% to 224% in average. An exclusive increase of the targeted eigenfrequency was not possible. While the **Figure 18.** Eigenfrequencies of the squared plate predeformed according to its 1st mode shape depending on different maximum relative pre-deformations and the boundary conditions characterised by all sides clamped (a), two opposite sides clamped, the other two free (b) and one side clamped, the other three free (c). The highest 1st eigenfrequency increases are given.

objective to maximise the 1st and 3rd eigenfrequency resulted in highest increase of the targeted eigenfrequency, pre-deforming the plate according to the 2nd or 4th bending mode shape resulted in highest rise of an eigenfrequency, which was not targeted. However, the targeted eigenfrequencies also increased.

Maximisation of multiple eigenfrequencies. The plate predeformation according to linear combinations of the 1st, 2nd and 3rd bending mode shape resulted in frequency rises for all analysed eigenfrequencies (Table 4). However, while a plate adaptation according to the bending mode 1 caused highest increase for the 1st eigenfrequency, the plate pre-deformations according to higher-order mode shapes did not always rise the targeted eigenfrequencies most. Both the linear combination of the bending modes 1 and 2 and of the bending modes 1, 2, and 3 resulted in high eigenfrequency increases for the 1st, 2nd, and 3rd eigenfrequency, while the 4th eigenfrequency rose less. It has to be noted that due to symmetry, a linear combination of the bending modes 1 and 3 caused the same results as the linear combination of the bending modes 1 and 2.

Table 3. Eigenfrequency deviations compared to the reference plate for the eigenvector approach to increase the 1st, 2nd, 3rd and 4th eigenfrequency for a maximum pre-deformation δ_{max} of 10.0 mm. The average deviations of not targeted eigenfrequencies are given. The frequency which increased most is shown in bold numbers.

	Max. f _l	Max. f ₂	Max. f3	Max. f ₄
fı	487%	228%	228%	240%
f2	135%	159%	159%	257%
f ₃	135%	233%	233%	257%
f4	67%	135%	135%	118%
fs	40%	109%	109%	183%
f6	55%	135%	135%	183%
Deviation	87%	168%	153%	224%

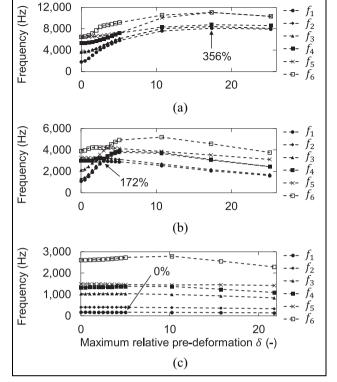
Next to the linear combinations, the weighted combinations of the bending modes 1 and 2 were also considered. The results of Figure 19 indicated that the highest eigenfrequency increase for the 1st eigenfrequency appeared at a full plate adaptation according to the bending mode 1 for all considered maximum predeformations. With increasing influence of the 2nd bending mode, the 1st eigenfrequency decreased, but rose again after a certain ratio. A similar trend could also be observed for the bending mode 2, where highest values of the 2nd eigenfrequency appeared at the complete plate adaptation according to the bending mode 2. The eigenfrequency minima appeared for the 1st eigenfrequency at c values of 0.4, 0.5 and 0.6 for maximum relative predeformations of 1.0, 3.0 and 5.0, respectively. The 2nd eigenfrequency was minimised at c values of 0.8, 0.7 and 0.7 for maximum relative pre-deformations of 1.0, 3.0 and 5.0. The factor c symbolises the ratio of mode 1 as described in the previous chapter.

Discussion

Slender beam

The numerical results obtained by using the defined sufficient mesh fineness coincided very well with the analytical results and can thus be seen as plausible.

Maximisation of a specific eigenfrequency or mode shape frequency. The investigated method to increase specific eigenfrequencies by shaping the beam according to the corresponding mode shape led to high increases of the targeted 1st, 2nd, 3rd, and 4th eigenfrequencies. The results for small maximum relative pre-deformations according to the 1st, 2nd, and 3rd bending mode shape coincided very well with the results obtained by³ indicating an exclusive increase of the targeted mode shape frequency, while especially lower-order eigenfrequencies remained constant. Similar results were obtained



δ	100% Mode I				100% Mode 2			
	fı	f2	f3	f4	fı	f2	fз	f4
I	84%	10%	10%	4%	40%	28%	82 %	25%
3	322%	69%	84%	30%	I 45%	138%	139%	107%
5	487 %	135%	135%	67%	228%	159%	233%	135%
δ	100% Mode 3				50% Mode 1 and 50% Mode 2			
	fı	f2	f3	f4	fı	f2	fз	f4
I	40%	28%	82 %	25%	31%	13%	40%	11%
3	145%	138%	139%	107%	103%	68%	125%	66%
5	228%	159%	223%	135%	181%	132%	141%	96%
δ	33% Mode I, 33% Mode 2, and 33% Mode 3							
	fi		f2		fз		f4	
I	28%		14%		46 %		13%	
3	99 %		76%		I 20%		69%	
5	169 %		137%		141%		115%	

Table 4. Frequency deviations of the first four eigenfrequencies compared to the reference plate for different maximum relative pre-deformations δ according to the 1st, 2nd and 3rd single bending mode shapes and combinations of these mode shapes. The highest increase is shown in bold numbers.

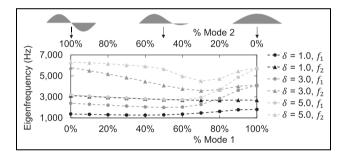


Figure 19. Ist and 2nd eigenfrequency of the plate predeformed with a maximum relative pre-deformation δ of 1.0, 3.0 and 5.0 according to weighted combinations of the bending modes 1 and 2. Three plate shapes are exemplarily shown for $\delta = 10.0$.

by shaping the beam according to the 4th bending mode shape. Among all investigated mode shape adaptations, highest frequency increase of up to 4,589% compared to the reference beam was achieved for the 1st bending mode by shaping the beam according to the targeted mode shape ($\delta = 20.0$). The pre-deformed beam corresponded to an arch, in which uniformly distributed transversal compressive forces lead exclusively to normal forces, as shear forces and bending moments disappear. This causes a high stiffness and thus high eigenfrequencies. However, this was only possible due to the axial (horizontal) constraint of the beam.

It has been stated that as soon as a beam receives a curved shape, the 1st mode shape becomes extensional implying high tangential displacement.³³ The 2nd mode

shape, however, remains as a non-extensional mode, which is characterised by a radial displacement. For small angles α , which imply small maximum relative pre-deformations, the 1st extensional mode shape has a lower frequency than the 2nd mode shape, which is non-extensional.³³ However, the 1st mode shape frequency rises fast with increasing pre-deformation, so that relatively small pre-deformations caused already a mode switch between the 1st and the 2nd mode, which could also be observed in the here analysed beam. The extensional energy is larger than the flexural energy, which is why the 1st mode shape frequency increases strongly with increasing maximum relative predeformation.

Regarding maximum relative pre-deformations larger than 5.0, variations (especially decreases) of not targeted eigenfrequencies were registered, which was not shown by Da Silva and Nicoletti.³ At the same time, higher pre-deformations allowed an additional, very high frequency increase of the targeted mode shape.

A special role played the 5th mode shape, that is, the 1st out-of-plane bending mode (1st bending mode in the *xy* plane). The 5th mode shape frequency decreased strongly with increasing maximum relative pre-deformation, disregarding which mode shapes the beam was pre-deformed according to. Since conservation of the beam mass was obtained by adapting, that is, reducing the beam width, the out-of-plane bending stiffness $Eb^3h_b/12$ was continuously reduced with higher pre-deformations. Thus, the beam became weaker in the *xy* plane and the corresponding frequency was reduced strongly independent of the applied pre-deformation of

the in-plane bending mode. Also the defined boundary conditions had a strong impact on the 5th mode shape frequency. As rotations around the *z* axis were prevented, the 5th mode shape was similar to a 1st in-plane bending mode shape of a clamped-clamped beam. It has been shown that this boundary condition results only in a small increase of the targeted eigenfrequency compared to a beam simply supported at both ends.³ Consequently, regarding the 5th mode adaptation, the corresponding frequency first increased (arch effect) slightly (boundary condition) and later decreased (stiffness reduction) with increasing maximum relative predeformation (Figure 10).

In summary, shaping the beam according to its mode shapes is an efficient way to strongly increase the frequency of a specific mode shape, while the frequencies of the other mode shapes only differ slightly. It has to be noted, however, that the occurrence of mode switching and multiple eigenfrequencies due to the applied pre-deformations is possible.

Maximisation of multiple eigenfrequencies. Beam predeformations according to linear combinations of mode shapes are an efficient method to increase the frequencies of various mode shapes simultaneously. Shaping the beam according to its 3rd bending mode shape, for example, increased the corresponding frequencies strongly, while the frequencies of the other analysed mode shapes remained constant or decreased insignificantly. Pre-deformations according to a linear combination of the 1st, 2nd, and 3rd bending mode shape, however, increased all three corresponding frequencies, while the frequency of the 4th bending mode shape changed only slightly. The weighted combinations of the bending modes 1 and 2 showed that the highest sum of the 1st and the 2nd eigenfrequency was obtained by shaping the beam according to only the 1st or only the 2nd bending mode shape depending on the maximum relative pre-deformation.

Squared plate

As for the analysed beam, the chosen mesh properties led to numerical results that corresponded to the analytical results for maximum relative pre-deformations up to 5.0. Higher pre-deformations were not calculated analytically, as the shell theory only applies for shallow shells.

Maximisation of a specific eigenfrequency or mode shape frequency. By adapting the shape of a simply supported squared plate according to a specific mode shape, an increase of all eigenfrequencies could be stated up to a certain pre-deformation. Applying this method to a fully clamped plate resulted only in a small eigenfrequency increase, while the absence of in-plane constraints hardly altered the eigenfrequencies, which had already been demonstrated for beams.³

It has been shown that by applying structural modifications, the mode shapes change, resulting in an increase of the corresponding eigenfrequency.²⁶ In the case of the 1st bending mode shape, the alterations of the maximum relative pre-deformations led to an increase of the 1st eigenfrequency. It was concluded that the highest eigenfrequency increase of dimpled plates can be achieved by placing the dimples at the plate centre, as it increases the local bending stiffness at the region of high modal strain energy. Modal strain energy was not taken into account in this study, so no statement can be made with regard to the results obtained. However, as the modal strain energy is defined as the product of the element stiffness matrix and the second power of the mode shape component,³⁶ it can be assumed that the mode shape has a high influence and that therefore areas, which show a high modal deflection, also obtain high modal strain energy. This would mean that published statements on dimpled plates can also be applied here to a certain extend.²⁶

The results of this study showed that by altering the shape of a plate, a change in the mode shape frequencies and order always occurred. In addition, a connection between the mode shape change and the eigenfrequency increase could be observed which supports already published statements.²⁶ Considering the results of the eigenfrequency increase connected to their mode shapes, high mode shape changes led to high increases of the corresponding frequencies, while frequencies connected to mode shapes similar to the reference plate obtained low eigenfrequency increases. This can exemplarily be seen for the adaptation of the plate according to the 1st bending mode shape, where the bending modes 4 and 6 showed minor mode shape alterations up to a maximum relative pre-deformation of 5.0 and therefore also low corresponding eigenfrequency increases. High increases were achieved for the frequencies connected to the first three bending modes, while the 1st bending mode frequency showed the highest rise. This 1st bending mode shape was not present within the first 50 mode shapes, leading to the conclusion that a high frequency increase of this mode shape appeared. In analogy to the beam, the 1st extensional mode shape is characterised by a higher energy than for example the 2nd flexural mode shape. Consequently, highest increase of the 1st bending mode frequency was also observed for the plate. Next, the 2nd and 3rd bending mode frequencies changed most. A decrease of the wave dimensions could be found within pre-deformations of 3.0 and 5.0, while the remaining mode shapes seemed similar to the ones occurring for a reference plate.

Similar conclusions could be made for the plate predeformations according to the bending modes 2/3 and

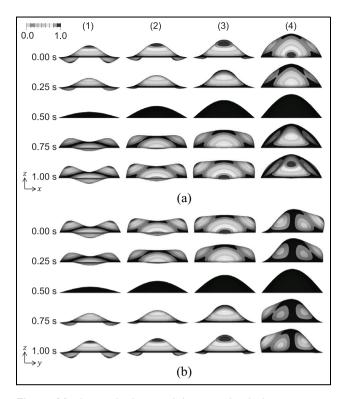


Figure 20. Ist mode shape and the normalised vibration amplitude of the squared plate pre-deformed according to the 1st mode shape at different time frames *t*. The maximum relative pre-deformations 5.0 (1), 10.0 (2), 15.0 (3) and 20.0 (4) are shown in a front view (a) and a side view (b).

4. It has to be noted that the shape adaptation according to the 2nd and 3rd mode shapes led to the same results as the analysed plate was characterised by a constant edge length. Consequently, the 2nd and 3rd eigenfrequency coincided and the 2nd mode shape corresponded to the 90° rotated 3rd mode shape.

Regarding the eigenfrequency decrease for high maximum relative pre-deformations, it has been shown that the 1st plate eigenfrequency increased with rising dimple height up to an optimum height and later decreased with further increasing dimple height.²⁶ It was argued that the eigenfrequency decrease was due to the dimple thickness decrease. Also in the present study, the eigenfrequency decreases for high predeformations might be due to the reduction of the bending stiffness because of the decreasing plate thickness, which was inevitable to have a constant mass (cf., the discussion in connection with the 5th beam mode in a the previous chapter). Moreover, it has to be noted that the formulae used for the eigenfrequency calculations are restricted to small pre-deformations and different dependencies can be expected for plates with high pre-deformations. Figure 20 shows exemplarily the 1st mode shapes of a plate pre-deformed according to the bending mode 1 with maximum relative predeformations of 5.0, 10.0, 15.0 and 20.0, where it can be seen that a flexural mode shape appeared for all predeformations. However, it was shifted from a mainly zmovement to a movement in the v direction with increasing maximum relative pre-deformation due to the reduced bending stiffness because of the smaller plate thickness. As up to a relative pre-deformation of 15.0 an increase of the 1st eigenfrequency can be stated, a connection between the eigenfrequency increase and the change in mode shape direction can be assumed. A comparison of the results from plates shaped according to the bending modes 2/3 and 4 and their mode shapes confirm this assumption. The proposed methodology to increase eigenfrequencies can therefore only be used when deflections of the mode shapes in the same direction as the pre-deformations appear.

In contrast to the results for the slender beam, for which the exclusive increase of a specific mode shape frequency was quite successful, the plate predeformation according to different mode shapes always led to an increase of all considered eigenfrequencies. Consequently, in the case of plates, the mode shapes cannot be seen as independent. It has rather to be noted that all mode shapes deflected certain plate areas, which were also influenced by other mode shapes. Thus, all eigenfrequencies altered when shaping the plate according to certain eigenvectors. In the here obtained results, the frequency of the 1st bending mode showed the highest or second highest frequency increase for any plate adaptations. As the 1st bending mode shape elevates the whole plate area, it is expected that an influence on the corresponding frequency is always present when adapting the plate according to any other mode shape. Though, the highest increase of the 1st eigenfrequency could be achieved when shaping the plate according to the bending mode 1.

Maximisation of multiple eigenfrequencies. In contrast to the investigated beam, the linear combination of several mode shapes cannot be seen as a useful method to simultaneously increase multiple eigenfrequencies as none of the considered combinations led to as high eigenfrequency increases as the adaptation according to the corresponding mode shape. While in the case of the slender beam it was possible to track the frequency changes of the different mode shapes, the mode shapes of the plate altered so strongly that tracking was not possible. Consequently, only the increase of the different eigenfrequencies was evaluated regardless of the changing mode shapes.

The weighted combinations of the bending modes 1 and 2 also indicated that an adaptation according to one single mode shape simultaneously caused an increase of the other eigenfrequencies. Thus, an adaptation according to one mode shape provided better results than a combination, concerning both eigenfrequencies, which had also been observed for the slender beam. The results are consequently in accordance with other findings.^{3,31}

By comparing all plate shapes considered within this report, the highest sum of the first two and three eigenfrequencies were achieved for the eigenvalue adaptation according to the bending mode 4. As described above, highest increase of the first three eigenfrequencies were present, especially influenced by the high values of the 2nd and 3rd eigenfrequency. In summary, it can be concluded from the results that the linear combination itself is not a good method to reach highest increases of several eigenfrequencies.

Comparison between slender beam and squared plate

It has been demonstrated that it is possible to increase the *i*-th mode shape frequency by adapting it according to the *i*-th mode shape while keeping all lower-order mode shape frequencies constant.³ The here described extension of the work demonstrated that this frequency increase is also possible for high-order mode shape frequencies (i = 4, 5). High maximum relative pre-deformations, however, led to small decreases of the not targeted mode shape frequencies.

Contrary to the results obtained for the squared plate, the not targeted beam eigenfrequencies changed only slightly, while a strong change was visible for all plate mode shape frequencies. Thus, applying the eigenvector approach to plates results in strong increases of the targeted eigenfrequencies, but also in high changes of other eigenfrequencies. It can be concluded that the eigenfrequency increase is less specific. This also comes along with intensive alterations in mode shapes, which is why a mode tracking was only partly possible. The beam mode shapes, in contrast, changed less strongly their mode shape order and frequency apart from the targeted mode shape.

Comparison of the different methods to increase eigenfrequencies

Different methods to increase eigenfrequencies and frequencies of specific mode shapes have been published. The main advantage of the eigenvector method to increase eigenfrequencies over other types of stiffening techniques and optimisations is that stiffening and thus high eigenfrequency increase is already achieved with small computational effort. In addition, no mass increase is needed and the method is already efficient for small pre-deformations. To underline this statement, an 1st eigenfrequency increase of 300% for the beam and 221% for the plate was generated by a maximum relative pre-deformation of 2.0 according to the 1st bending mode shape. The 1st bending mode frequency of the beam was even increased by 400%. Maximum relative pre-deformations of 3.0 led to 1st eigenfrequency rises of 300% for the beam and 322% for the plate. The 1st bending mode frequency of the beam increased by 641%. Regarding the plate, both mentioned eigenfrequency increases corresponded to the 1st bending mode shape.

The eigenfrequency increase of simply supported 2D beams by applying topology optimisations have been studied.³⁷ Within this study, a 1st eigenfrequency increase of about 155% and a 2nd eigenfrequency increase of about 135% including mass reduction of 50% was reached. Both values were significantly lower than the beam eigenfrequency increases obtained by the eigenvector approach. The increase of the 1st eigenfrequency of a squared plate with similar boundary conditions has also been investigated by using topology optimisation, where an eigenfrequency increase of 102% with a mass reduction of 50% occurred.37 A study on the optimal distribution of dimples and beading using a genetic algorithm to maximise the 1st eigenfrequency of simply supported rectangular plates revealed maximum frequency increases of 25.4% $(\delta_{max} = 9.0)$ and 160.5% $(\delta_{max} = 7.7)$.²⁶ Comparing these values with the plate eigenfrequency increases that were achieved within this study, higher values occurred at all times, taking similar pre-deformations into account. As the eigenvector method is most similar to beading, it is focused in the following on a comparison between these two techniques. The following advantages and disadvantages apply in both cases: Stiffening is achieved without any additional weight increase and without adding any joints. However, the methods are limited to the material capacity to deform plastically without cracking.²⁶ Beading causes a higher computational effort than the shape adaptation according to eigenvectors, but smaller or similar expenses in the manufacturing are expected as the beads have to be applied only in some parts of the plate. While the embossing is restricted to few places, the overall plate shape needs to be adapted for the proposed eigenvector methodology. On the one hand, this can be seen as an advantage, as the not beaded area of beaded plates should be kept as small as possible as otherwise high stress peaks appear.³⁸ On the other hand, the bulging must also be possible without restraining attached parts or the installation space. Finally, it has to be noted that the eigenvector approach is highly dependent on the boundary conditions and is not applicable to structures which are not axially constrained. Moreover, the applied pre-deformations certainly lead to more complex geometries compared to the un-deformed structures as exemplarily illustrated in Figure 5. But also other mentioned structural methods to rise eigenfrequencies involve an

increase of structural complexity. However, the present results showed that even small pre-deformations, which are supposed to be applied more easily, increased the eigenfrequencies stronlgy. Nevertheless, the presence of axial constraints is inevitable to successfully apply the methodology.

The eigenfrequencies of the here analysed beam and plate have also been increased by applying evolutionary strategic optimisations and topography optimisations to compare these optimisation techniques to the eigenvector approach. The work will soon be presented in a different publication.

Overall, it can be concluded that the comparison of the eigenvector approach to other optimisation techniques indicated higher eigenfrequency increases and frequency increases of specific mode shapes for the former methodology.

Conclusion

We have shown that shaping axially constrained beams and squared plates according to their mode shapes leads to strong eigenfrequency increases. This methodology allows the almost exclusive increase of the frequency of the *i*-th mode shape (i = 1-5) of the beam, while the increase of the *i*-th plate mode shape frequency (i = 1-4) simultaneously alters strongly other eigenfrequencies. However, comparing the eigenfrequency increase approach to other published studies, it is a promising, computationally efficient method to strongly increase eigenfrequencies and frequencies of specific mode shapes by applying only small structural pre-deformations without additional weight increase.

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ORCID iD

Simone Andresen in https://orcid.org/0000-0002-3873-7089

Supplemental material

Supplemental material for this article is available online.

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