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# A salty deep ocean as a prerequisite for glacial termination

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**Deglacial transitions of the middle to late Pleistocene (terminations) are linked to gradual changes in insolation accompanied by abrupt shifts in ocean circulation. However, the reason these deglacial abrupt events are so special compared with their sub-glacial-maximum analogues, in particular with respect to the exaggerated warming observed across Antarctica, remains unclear. Here we show that an increase in the relative importance of salt versus temperature stratification in the glacial deep South Atlantic decreases the potential cooling effect of waters that may be upwelled in response to abrupt perturbations in ocean circulation, as compared with sub-glacial-maximum conditions. Using a comprehensive coupled atmosphere–ocean general circulation model, we then demonstrate that an increase in deep-ocean salinity stratification stabilizes relatively warm waters in the glacial deep ocean, which amplifies the high southern latitude surface ocean temperature response to an abrupt weakening of the Atlantic meridional overturning circulation during deglaciation. The mechanism can produce a doubling in the net rate of warming across Antarctica on a multicentennial timescale when starting from full glacial conditions (as compared with interglacial or subglacial conditions) and therefore helps to explain the large magnitude and rapidity of glacial terminations during the late Quaternary.**

A multitude of ice-core and marine records show that the past 800,000 years of Earth's history were dominated by global-scale alternations between glacial and interglacial conditions with a high degree of nonlinearity in the climatic response to gradual external forcing<sup>1,2</sup>. This is particularly relevant for glacial terminations, which are thought to be driven by relatively weak insolation forcing due to changes in Earth's orbital configuration<sup>3,4</sup>. In this context, it has been proposed that disruptions to circulation patterns in the coupled atmosphere–ocean system might have acted to accelerate Antarctic warming and the outgassing of CO<sub>2</sub> from the Southern Ocean (SO) to cross a critical level necessary to maintain interglacial conditions at a global scale<sup>5–7</sup>. A possible mechanism to reach such a point is an abrupt weakening of the Atlantic meridional overturning circulation (AMOC) and a corresponding reduction of northward oceanic heat transport associated with SO and Antarctic warming, according to the bipolar-seesaw concept<sup>8–10</sup>, which has been detected in a wide spectrum of models<sup>11,12</sup>. However, since abrupt oscillations between relatively weak and strong AMOC states and corresponding SO changes are thought to be a ubiquitous feature of millennial-scale climate variability throughout most glacial periods<sup>13,14</sup>, it is open to question what facilitates such a strong temperature response during deglaciation (Fig. 1).

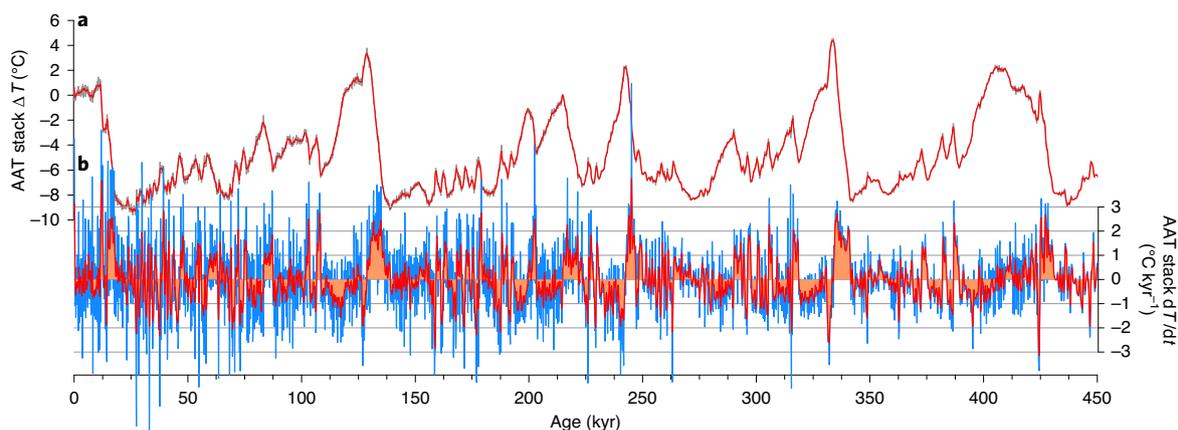
## Interglacial and glacial water-mass configurations

The large-scale meridional water-mass distribution in the modern ocean is characterized by deep waters originating in the North Atlantic (North Atlantic Deep Water; NADW) and the SO (Antarctic Bottom Water; AABW). These water masses are connected by mixing

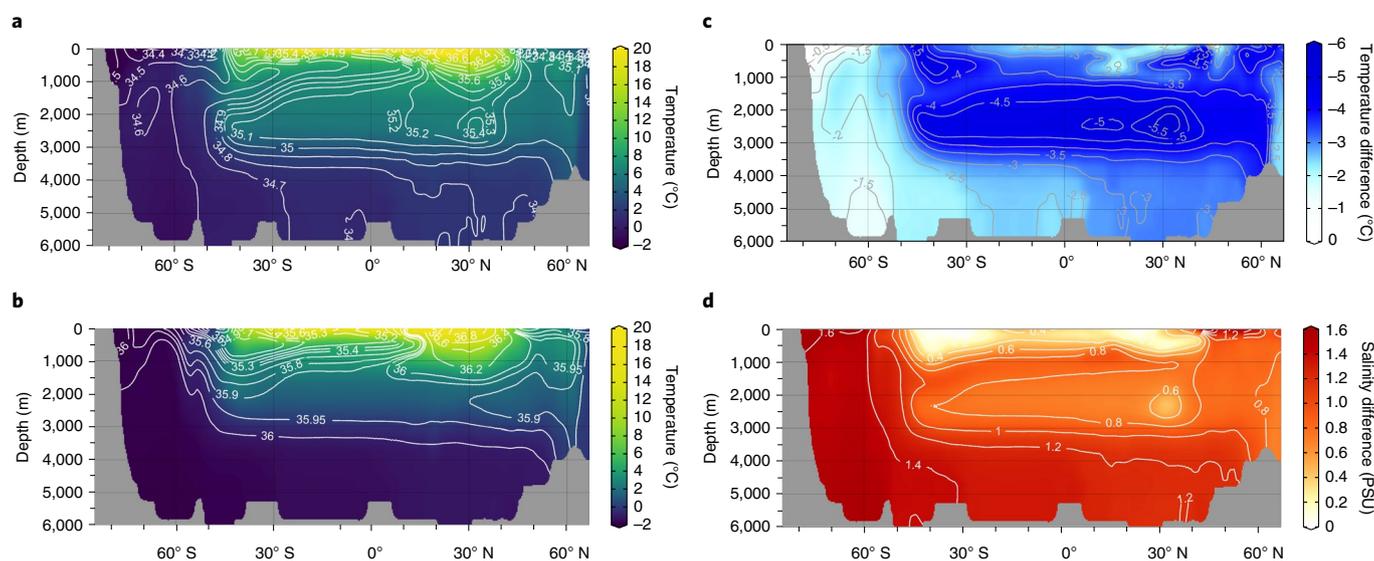
processes and wind-driven upwelling within the SO<sup>15–18</sup>. AABW represents the densest water mass in the modern ocean due to its low temperature (AABW is relatively fresh and cold compared with the overlying NADW)<sup>18,19</sup>. Only at great depth does the pressure dependence of the thermal expansion coefficient in the equation of state (the thermobaric effect) mean that the warmth of NADW makes it less dense than AABW<sup>20</sup>. Thus, a reduction in the temperature contrast between the two water masses will tend to destabilize the water column, ultimately resulting in NADW replacing AABW as the deepest water mass.

With respect to ocean temperature, the most accurate/robust evidence for glacial ocean cooling on a global scale stems from noble gases trapped in ice cores. Bereiter et al.<sup>21</sup> show that mean ocean cooling ( $2.5 \pm 0.24$  °C) was greater than the maximum possible cooling of AABW, since modern AABW has a temperature of  $-0.88$  °C<sup>22</sup> and seawater has a freezing point of about  $-2$  °C. Of particular relevance to our study is the implication that 'warmer' water masses such as NADW and intermediate water masses within the Southern Atlantic must have cooled to a greater extent than AABW<sup>23</sup>. This implies that thermal deep-ocean stratification in the glacial ocean was weaker than it is today. In a hypothetical scenario with no change to the salinity distribution, we can see that such heterogeneous cooling of the glacial ocean could lead to an unstable water column (Extended Data Fig. 1), and yet proxy reconstructions suggest that the glacial version of AABW maintained its position as the deepest water mass at least in the South Atlantic region<sup>24,25</sup>. Therefore, something probably acted to compensate for the loss of stability from thermal stratification in the glacial ocean.

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**Fig. 1 | Records of Antarctic temperature and the rate of temperature change. a**, Time series of Antarctic temperature anomaly<sup>57</sup> (AAT  $\Delta T$ ). **b**, Rate of change ( $dT/dt$ ) calculated after a temporal smooth of 500 years (red curve) and 200 years (blue curve). Deglacial events display higher rates of change on average (see Methods for discussion).



**Fig. 2 | Inter-glacial and glacial temperature and salinity distribution and changes between them as zonal averages in the Atlantic sector. a, b**, Inter-glacial (pre-industrial,  $PID_{CTRL}$ ) (**a**) and glacial (Last Glacial Maximum,  $LGM_{CTRL}$ ) (**b**) temperature and salinity (practical salinity units (PSU), white contours) distribution. **c, d**, Glacial-interglacial difference ( $LGM_{CTRL} - PID_{CTRL}$ ) for temperature (**c**) and salinity (**d**). All fields represent 100 yr averages. Plotted using Ocean Data View (<http://odv.awi.de>).

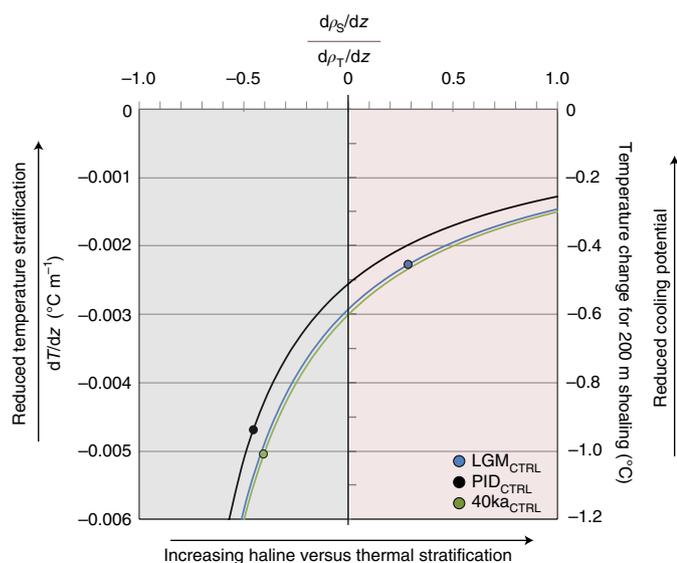
Indeed, other proxy evidence<sup>26,27</sup> suggests that the salinity of glacial AABW could have been much greater than today, even to the point that deep-ocean salinity stratification may have been reversed with respect to the modern<sup>26</sup>. Such salinity changes could be brought about by enhanced equatorward sea-ice export in the glacial SO<sup>28</sup> in combination with reduced melting of Antarctic ice as a consequence of colder NADW<sup>20</sup>. Analyses of different boundary conditions in atmosphere–ocean general circulation model (AOGCM) equilibrium simulations suggest that especially the interplay of low  $CO_2$ , low obliquity<sup>29</sup> and large Northern Hemisphere ice sheets<sup>30</sup> provides favourable conditions for the formation of salty deep waters around Antarctica.

Using the fully coupled Earth system model (COSMOS)<sup>29,31</sup> (Methods), Zhang et al. have previously shown<sup>29</sup> that the application of glacial boundary conditions can lead to particularly salty deep waters in the SO region compared with pre-industrial conditions. In this study, we augment the glacial ( $LGM_{CTRL}$ ) and pre-industrial ( $PID_{CTRL}$ ) control experiments of Zhang et al.<sup>29</sup> to

include an intermediate climate state ( $40ka_{CTRL}$ , Methods, Figs. 3 and 4 and Extended Data Figs. 2–9) to investigate any differences in the response to ‘terminal’ and non-terminal AMOC perturbations (which we describe later).

In the glacial ocean simulation  $LGM_{CTRL}$ , the strongest temperature decrease (of more than  $4^\circ C$  with respect to both  $PID_{CTRL}$  and  $40ka_{CTRL}$ ) occurs within levels of the southward NADW return flow (Fig. 2 and Extended Data Figs. 2 and 3). The local  $LGM_{CTRL}$  cooling maximum at  $\sim 2.5$  km depth (Fig. 2c) can be linked to a combination of a shoaling ( $\sim 200$  m) of the glacial NADW cell and production of colder NADW (although we note that a model–data comparison for our glacial state with marine calcite  $\delta^{18}O$  values includes the possibility that NADW formation might actually have been weaker and even shallower<sup>32–34</sup> than in our simulated LGM state<sup>35</sup>).

In stark contrast to the changes at intermediate depths, the weakest temperature and strongest salinity differences between the pre-industrial and glacial states are detected in the deep SO (Fig. 2c,d). Here we observe a shift from a purely temperature-controlled



**Fig. 3 | Temperature changes with depth in dependence of haline to thermal density ratio and background climate.** Temperature gradients ( $dT/dz$ ) along lines of constant density stratification ( $d\rho/dz$ ) for the individual climate states with reference conditions from the zonally averaged Atlantic at  $30^\circ\text{S}$  and 3,000 m depth. The pressure reference level is 3,000 dbar to account for the impact of pressure conditions on temperature and stability properties at the chosen reference level of 3,000 m depth. The circles show the position of the simulated interglacial ( $\text{PID}_{\text{CTRL}}$ , black), 40 ka ( $40\text{ka}_{\text{CTRL}}$ , green) and Last Glacial Maximum ( $\text{LGM}_{\text{CTRL}}$ , blue) states (compare Extended Data Fig. 6). The curves show solutions for stable density stratification ( $d\rho/dz > 0$ ; that is, density increases with depth) with a stable thermal stratification ( $d\rho_T/dz > 0$ ; that is, temperature decreases with depth). Hence, positive values on the x axis indicate a stable haline stratification ( $d\rho_S/dz > 0$ ; that is, salinity increases with depth, pink shaded phase space) and negative values indicate an unstable haline stratification ( $d\rho_S/dz < 0$ ; that is, salinity decreases with depth, grey shaded phase space).

stratification to one that is also characterized by a stable haline density gradient (at least within the SO). A similar contrast is observed between glacial and intermediate (40 thousand years ago (ka)) conditions (Extended Data Fig. 2).

We note that the mechanism of deep-ocean salinification proposed by Adkins<sup>20</sup> (which calls on the cooling of NADW as a means for reducing the melting of Antarctic land ice and related input of freshwater to the SO) could work in tandem with increased sea-ice export (as observed in our model simulation) to produce highly saline bottom waters.

An overview of simulated salinity structure in LGM states (unperturbed) encompassing the last three Paleoclimate Model Intercomparison Project (PMIP)/Coupled Model Intercomparison Project (CMIP) model generations<sup>36–40</sup> shows that a simulation of a salty glacial deep ocean (Extended Data Fig. 10) by AABW on a global scale is a persistent challenge, which partly reflects limits of current models to simulate key processes<sup>41,42</sup> that influence AABW characteristics. This also includes the mechanism described by ref. 20, and while the current generation of coupled AOGCMs cannot resolve that mechanism explicitly, we can at least demonstrate its viability using a coupled ice-sheet/ice-shelf cavity model (Methods and Extended Data Fig. 4). However, the precise mechanism by which a salty deep ocean is generated is not important for the following discussion, in which we investigate the theoretical implications of a relative increase in the importance of salinity versus temperature stratification in the glacial ocean. For further discussion, see Methods.

### Salinity gradient and temperature changes with depth

Deep-ocean stratification depends on the interplay between salinity and temperature changes with depth. We therefore derive an analytical expression (which is independent of the model simulations) for the thermal gradient ( $dT/dz$ ) that accounts for the interplay between thermal and haline contributions to the net density change with depth:

$$\frac{dT}{dz} = -\frac{d\rho/dz}{\alpha_T \rho_0} \left( 1 + \frac{d\rho_S/dz}{d\rho_T/dz} \right)^{-1}$$

Here,  $d\rho_S/dz$  and  $d\rho_T/dz$  represent the haline and thermally driven density changes with depth and  $d\rho/dz$  is the combined (net) change in density with depth (which must remain positive to maintain stability). Density and the thermal expansion coefficient are indicated by  $\rho_0$  and  $\alpha_T$ . We now explore the implications of a glacial ocean state in which an increase in salinity stratification accommodates a decrease in the thermal contrast between northern- and southern-sourced deep-water masses. In particular, we are interested in the surface SO temperature response that would result from the upwelling of deeper waters (for example, due to a perturbation of the AMOC) that are relatively warm (not as cold) compared with a non-glacial state.

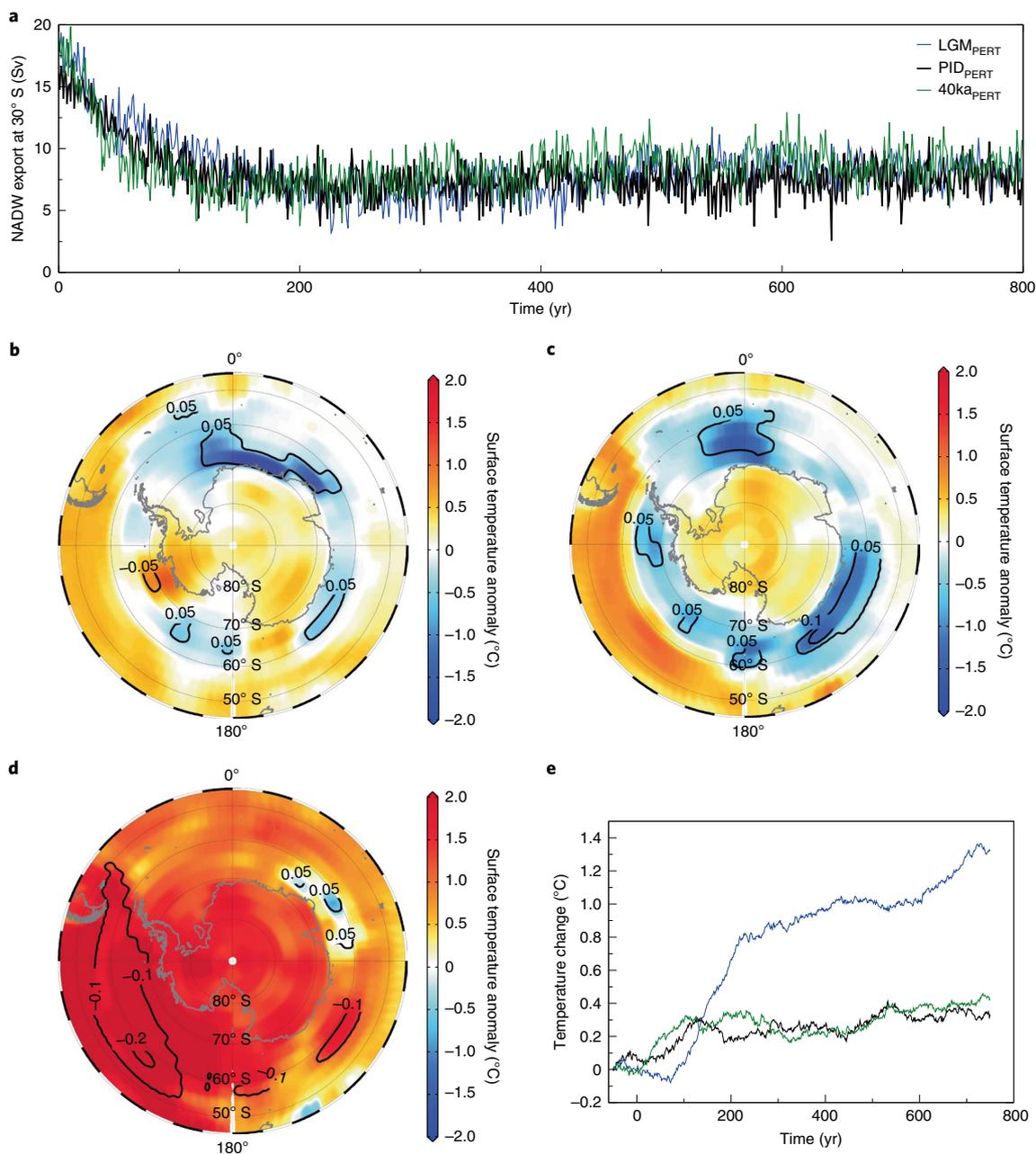
With the aid of the different terms in the equation, we can evaluate the importance of the stable glacial salinity stratification for the reduced thermal stratification in  $\text{LGM}_{\text{CTRL}}$ . We focus our attention on the deep South Atlantic (zonally averaged at  $30^\circ\text{S}$ ) to assess the export characteristics out of the Atlantic basin and to avoid effects of intra-basin scale upwelling on Southern Hemisphere communication. In particular, the choice of the depth level influences the thermohaline characteristics of the reference water masses. Therefore, we choose a level at 3,000 m depth, which is within a broader core of waters with similar stratification characteristics that represent an unstable ( $\text{PID}_{\text{CTRL}}$ ,  $40\text{ka}_{\text{CTRL}}$ ) or stable ( $\text{LGM}_{\text{CTRL}}$ ) salinity stratification. The first term on the right-hand side of the equation can then be diagnosed from our modelled control runs, enabling us to plot  $dT/dz$  versus the ratio of haline to thermally driven density changes

( $\frac{d\rho_S/dz}{d\rho_T/dz}$ ) for a variety of background states (Fig. 3).

For any of the three different background states with a given density structure ( $d\rho/dz$ ,  $\rho_0$  and  $\alpha_T$ ), it can be seen that for constant  $d\rho/dz$  the ratio of haline to thermally driven density stratification must increase to accommodate any decrease in the vertical temperature contrast between northern- and southern-sourced deep-water components in the South Atlantic (coloured curves in Fig. 3). The key difference between our glacial simulation and a pre-industrial or intermediate (40 ka) state is the large increase in salinity stratification (which is actually reversed with respect to the non-LGM states) relative to the thermal stratification. Together with the palaeo-evidence described earlier, this suggests that the cooling potential of upwelled water to the surface SO in a full glacial is relatively much less than for either the pre-industrial or intermediate state. Naturally, this has implications for the oceanic and atmospheric response to a perturbation of the AMOC that might be expected as a necessary ingredient of glacial termination.

### Southern high-latitude response to AMOC weakening

We next perform perturbation simulations to investigate the implications of different glacial–interglacial water-mass configurations with respect to SO and Antarctic temperature responses to an abrupt weakening of the AMOC (Fig. 4 and Extended Data Fig. 5). We use the same model set-up as in the control simulations and apply a freshwater perturbation of 0.15 Sv ( $1\text{ Sv} = 10^6\text{ m}^3\text{ s}^{-1}$ ) to the North Atlantic ice-rafted debris belt for 800 years (further details are provided in Methods). In summary, we find that warming sea surface conditions and declining sea-ice cover occur as the dominant



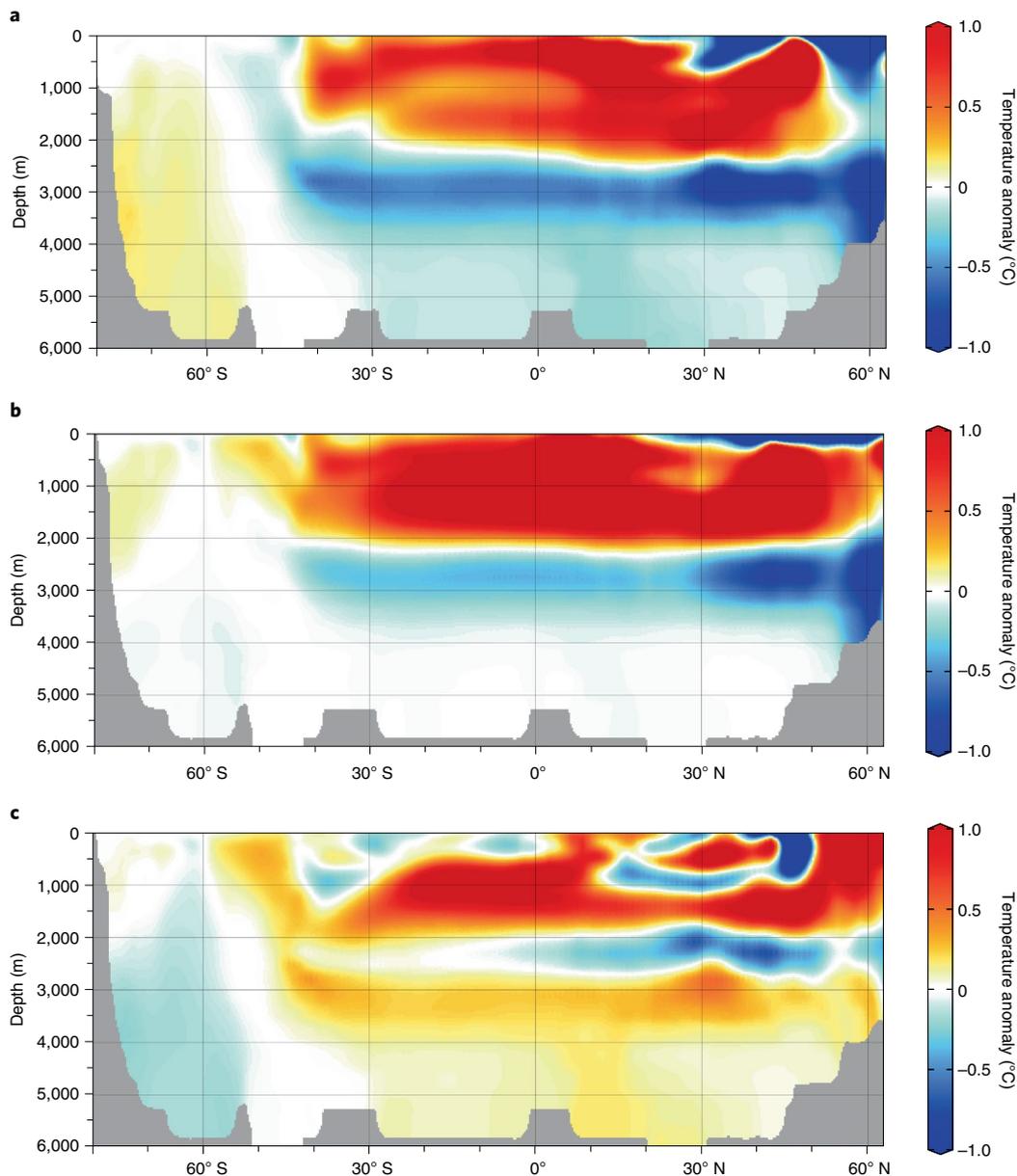
**Fig. 4 | Changes in circulation strength, southern high-latitude temperature and sea-ice area fraction in response to North Atlantic freshwater perturbation. a**, NADW export (Sv) at 30° S. **b–d**, Surface temperature (°C) and sea-ice area fraction changes (iso-lines) for interglacial ( $\text{PID}_{\text{PERT}} - \text{PID}_{\text{CTRL}}$ ) (**b**), 40 ka ( $40\text{ka}_{\text{PERT}} - 40\text{ka}_{\text{CTRL}}$ ) (**c**) and Last Glacial Maximum ( $\text{LGM}_{\text{PERT}} - \text{LGM}_{\text{CTRL}}$ ) (**d**) conditions, averaged between model years 700 and 800. **e**, The averaged Antarctic surface temperature changes as a 100 yr running mean with the same colour code as in **a**. Panels **b–d** plotted using Ocean Data View (<http://odv.awi.de>).

high-latitude Southern Hemisphere response to an AMOC weakening only under LGM conditions (Fig. 4d,e).

In the following discussion, we focus on the comparison of the perturbation simulations  $\text{LGM}_{\text{PERT}}$  and  $\text{PID}_{\text{PERT}}$  (the  $40\text{ka}_{\text{PERT}}$  responses are similar to those of  $\text{PID}_{\text{PERT}}$ ) (Fig. 4). On weakening of AMOC, surface warming at the southern boundary of the Atlantic basin (~30° S) penetrates to a depth of ~2 km irrespective of the background conditions (Fig. 5a,b). This warming response is a common characteristic among NADW perturbation simulations<sup>43,44</sup>, which can be linked to heat storage<sup>45</sup> in the low latitudes and South Atlantic and reduced ventilation of intermediate and deep waters, combined with the downward mixing of heat from

the thermocline<sup>46</sup>. This ‘advection–diffusion balance’ concept is also consistent with the relatively strong  $\text{LGM}_{\text{PERT}}$  warming (Fig. 5) along the perturbed glacial AMOC-return pathway (Extended Data Fig. 7), reflecting the more pronounced vertical temperature contrast across the NADW cell in the unperturbed  $\text{LGM}_{\text{CTRL}}$  state (Fig. 2 and Extended Data Fig. 3).

In all experiments, the Atlantic basin temperature response below ~2 km is one of cooling (Fig. 5 and Extended Data Fig. 8) as the NADW cell weakens (Fig. 4a) and shoals (Extended Data Fig. 7) during the first ~200 years of hosing. Compared with  $\text{PID}_{\text{PERT}}$ , the more pronounced  $\text{LGM}_{\text{PERT}}$  cooling between 2.0 and 2.5 km (Fig. 5c) reflects the initially shallower NADW cell (Extended Data Fig. 3) as



**Fig. 5 | Temperature anomalies and the difference between them. a, b**, Temperature anomalies by AMOC weakening for interglacial ( $PID_{PERT} - PID_{CTRL}$ ) (**a**) and Last Glacial Maximum ( $LGM_{PERT} - LGM_{CTRL}$ ) (**b**) conditions between model years 100 and 200 as a 100 yr mean, zonally averaged in the Atlantic sector. **c**, The difference between the glacial and interglacial anomalies ( $LGM - PID$ ). Plotted using Ocean Data View (<http://odv.awi.de>).

the AMOC is perturbed (Extended Data Fig. 7). Below this depth, however, cooling in  $LGM_{PERT}$  is weaker than in the  $PID_{PERT}$  (Fig. 5c). At first glance, this is perplexing since shoaling of the overturning cell (Extended Data Fig. 7) and isopycnals at these depth levels is more pronounced in  $LGM_{PERT}$  (the density levels at 3,000 m depth in  $PID_{CTRL}$  and  $LGM_{CTRL}$  shoal by  $\sim 160$  m and  $\sim 230$  m, not shown). However, the phenomenon is a direct consequence of the build-up of relatively warmer waters in the glacial deep ocean (with respect to intermediate depths) permitted by the increase in salt stratification (Fig. 3 and Extended Data Fig. 6).

As predicted, this fundamental change in deep-ocean stratification (from temperature controlled in  $PID_{CTRL}$  to one that is dependent on a stable haline density gradient in  $LGM_{CTRL}$ ) leads to a very different SO surface temperature response to North Atlantic freshwater forcing and AMOC perturbation. In  $PID_{PERT}$ , cooling is observed across the surface SO due to the upwelling of cold deep

waters when the AMOC is in a perturbed state. By contrast, for  $LGM_{PERT}$  between  $\sim 100$  and 200 years we observe a continuous surface warming across Antarctic circumpolar latitudes south of  $\sim 45^\circ$  S and extending into the Drake Passage (south of  $\sim 55^\circ$  S), reflecting the upwelling of relatively warmer waters that characterize the glacial deep ocean in our simulations (Fig. 5 and Extended Data Figs. 8 and 9). It is important to emphasize that cooling does occur in the deep ocean during  $LGM_{PERT}$  (Fig. 5b), but this cooling is reduced relative to  $PID_{PERT}$  (Fig. 5c) and, in combination with warming at intermediate depths, produces a net warming at the SO surface.

The contrasting SO surface temperature responses lead to very different temperature responses over Antarctica (Fig. 4e). In fact, the integrated surface heat flux anomaly south of  $55^\circ$  S is of opposite signs in  $LGM_{PERT}$  and  $PID_{PERT}$ . Only in the glacial state does the AMOC perturbation result in atmospheric warming in the southernmost ocean region. When combined with atmospheric

heat transport, the positive ocean–atmosphere heat flux observed in LGM<sub>PERT</sub> produces by far the largest and fastest temperature rise across Antarctica within the first ~200 years when compared with PID<sub>PERT</sub> and 40ka<sub>PERT</sub> (Fig. 4e). After ~200–300 years, NADW strength and geometry have largely adjusted to the persistent freshwater perturbation, and Antarctic warming proceeds at a reduced rate.

### Salty glacial deep ocean and deglaciation

We have demonstrated that a fundamental change in deep-ocean stratification can lead to a greatly enhanced climatic response to AMOC weakening while in a glacial state. The simulated Antarctic seesaw response is amplified by a factor of ~2 on a multicentennial timescale compared with interglacial or subglacial conditions. The surface cooling (as opposed to warming) that we observe in the SO high latitudes under pre-industrial and subglacial conditions might be a model-specific characteristic<sup>12</sup>, but it is important to note that the key to an amplified seesaw warming is a more pronounced haline to thermal density stratification ratio (which is predicted from proxy reconstructions, as described earlier). Increasing this ratio has the same effect irrespective of the climate background conditions, as shown in the phase space analysis in Fig. 3. The underlying mechanism might also therefore represent a key to understanding the spread of SO high-latitude responses in previous glacial AMOC perturbation experiments using a range of models<sup>47</sup>. For the specific model used in our study, a damped SO surface temperature response is obtained for LGM experiments using a modern-like salinity stratification<sup>29</sup> as employed by Kageyama et al.<sup>47</sup>.

Our simulations do not include a representation of the global carbon cycle, but several previous studies have called on AMOC-induced stratification changes within the SO as a key player in glacial–interglacial CO<sub>2</sub> variability<sup>48–52</sup>. Furthermore, it has also been suggested that AMOC-induced CO<sub>2</sub> changes<sup>53</sup> might provide an internal feedback with the potential to promote transitions between different climate states<sup>54</sup>. Abrupt changes in the mode of AMOC occurred throughout the last glacial period<sup>13</sup> and were linked to temperature changes across Antarctica and (on occasion) variations in CO<sub>2</sub><sup>55</sup>, and yet only those changes associated with the last termination were sufficient to mark the transition back to an interglacial state. Previous studies have suggested various combinations of ice-sheet size<sup>2</sup> and specific orbital configurations<sup>1,56</sup> in an attempt to explain why glacial termination occurs only when it does. Our findings provide an alternative explanation: that global-scale water-mass salinity characteristics represent a key player to promote deglaciation once a salty deep-ocean reservoir has been formed at full glacial conditions.

### Online content

Any methods, additional references, Nature Research reporting summaries, source data, extended data, supplementary information, acknowledgements, peer review information; details of author contributions and competing interests; and statements of data and code availability are available at <https://doi.org/10.1038/s41561-021-00857-3>.

Received: 15 August 2019; Accepted: 15 October 2021;  
Published online: 3 December 2021

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## Methods

**Rate of change in Antarctic temperature records.** In Fig. 1, the relatively large rates of change observed in the upper sections of the record (particularly the Holocene, which has not experienced high-amplitude millennial-scale variability) using a 200 yr pre-smooth reflects the much higher temporal resolution of the upper sections of ice cores relative to deeper sections. Using a pre-smoothing of 500 yr alleviates this problem while (in our opinion) preserving an accurate (not overly smoothed) representation of the relative rates of change associated with millennial-scale events.

**Model description.** We use the comprehensive fully coupled Earth system model COSMOS (ECHAM5-JSBACH-MPIOM) in this study. The atmospheric model ECHAM5<sup>38</sup>, complemented by a land surface component with dynamic vegetation JSBACH<sup>39</sup>, is used at T31 resolution (~3.75°), with 19 vertical hybrid sigma-pressure levels. The ocean model MPIOM<sup>40</sup>, including sea-ice dynamics that are formulated using viscous-plastic rheology<sup>61</sup>, runs on a bipolar curvilinear Arakawa-C grid. The grid size (120 × 101) corresponds to a formal longitudinal resolution of ~3° (360/120) and a latitudinal resolution of ~1.8° (180/101). The vertical domain is represented by 40 unevenly spaced depth layers. The climate model has already been used to simulate, for example, the past millennium<sup>31</sup>, the Miocene warm climate<sup>62,63</sup>, the Pliocene<sup>64</sup>, the internal variability of the climate system<sup>65</sup>, the Last Glacial Maximum climate<sup>29,47</sup> and glacial millennial-scale variability at intermediate climate conditions<sup>54,66</sup>. The glacial (LGM<sub>CTRL</sub>) and pre-industrial (PID<sub>CTRL</sub>) states used in our study are experiments LGMW and PI from Zhang et al.<sup>29</sup>. The glacial state is characterized by a salty deep-ocean reservoir in reasonable agreement with available salinity reconstructions. Note that the generation of an interglacial-to-glacial salinity reversal is expected to be dependent on, for example, the model itself, the model configuration and the experimental set-up (Extended Data Fig. 10). However, for the mechanism described here, a salinity reversal per se is not necessary. The data-based demand that the temperature gradient between northern and southern sources of deep water within the Atlantic was reduced during the LGM compared with today<sup>20</sup> implies that a reduction in the thermal deep-ocean stratification was accompanied by a stronger glacial salinity increase of AABW relative to NADW. Were this not the case, then the density stratification of the glacial South Atlantic/SO would necessarily have been weaker than today. A simple evaluation for the importance of a glacial salinity reversal for glacial–interglacial density changes is provided in the corresponding section.

**Experimental design.** Besides the pre-industrial (PID<sub>CTRL</sub>) and the Last Glacial Maximum (LGM<sub>CTRL</sub>) control climates from Zhang et al.<sup>29</sup>, we conduct an equilibrated MIS3 control simulation (40ka<sub>CTRL</sub>) by imposing fixed boundary conditions for 40 ka. In detail, the three orbital parameters in experiment 40ka<sub>CTRL</sub> are 0.013146 (eccentricity), 358.17° (precession of the equinoxes, the angle between the Earth's position during the Northern Hemisphere vernal equinox and the orbit perihelion) and 23.61° (obliquity)<sup>67</sup>. The greenhouse gas concentrations are 195 ppm (CO<sub>2</sub>), 413 ppb (CH<sub>4</sub>) and 231 ppb (N<sub>2</sub>O)<sup>68,69</sup>. The ice-sheet configuration is a combination of ice-sheet reconstructions from ICE-5G<sup>70</sup> and PMIP3. That is, the topography anomaly between 40 ka and the LGM in ICE-5G is first calculated and then is added to the PMIP3 LGM topography<sup>29</sup>. Global mean sea level is ~80 m lower at 40 ka than the pre-industrial level. Experiment 40ka<sub>CTRL</sub> is integrated for 5,000 years to an equilibrated state that serves as the starting point for the perturbation experiment.

Identical hosing experiments are then conducted under the three different climate backgrounds in experiments PID<sub>PERT</sub>, LGM<sub>PERT</sub> and 40ka<sub>PERT</sub>. A freshwater flux of 0.15 Sv is added to the Ruddiman Belt (40°–55°N, 45°–20°W) for 800 years to ensure a quasi-equilibrated response in the SO. A freshwater perturbation of 0.15 Sv agrees with estimates of deglacial meltwater release during Heinrich Stadial 1 as assessed in different studies<sup>71–73</sup>. Global mean sea-level height is corrected for the input of additional freshwater. Hence, coastlines remain the same during a freshwater perturbation experiment. Averages of 100 years are shown in our results.

**Vertical temperature and salinity changes.** On the basis of the linear approximation of the density equation and neglected pressure variations, we can calculate density changes with depth by:

$$\frac{d\rho}{dz} = \rho_0 \left( -\alpha_T \frac{dT}{dz} + \beta_S \frac{dS}{dz} \right) \quad (1)$$

where  $\rho_0$  is density,  $T$  is temperature,  $S$  is salinity,  $\alpha_T$  is the thermal expansion coefficient and  $\beta_S$  is the haline contraction coefficient. So that a temperature contribution to density changes with depth equates to:

$$\frac{d\rho_T}{dz} = -\rho_0 \alpha_T \frac{dT}{dz} \quad (2)$$

and a salinity contribution equates to:

$$\frac{d\rho_S}{dz} = \rho_0 \beta_S \frac{dS}{dz} \quad (3)$$

Rearrangement of equation (2) leads to:

$$\frac{dT}{dz} = -\frac{d\rho_T/dz}{\alpha_T \rho_0} \quad (4)$$

Further substituting  $d\rho_T/dz$  in equation (4) by:

$$\frac{d\rho_T}{dz} = \frac{d\rho}{dz} - \frac{d\rho_S}{dz} \quad (5)$$

temperature changes with depth can be expressed as a function of the ratio between the haline and thermal density changes with depth:

$$\frac{dT}{dz} = -\frac{d\rho/dz}{\alpha_T \rho_0} \left( 1 + \frac{d\rho_S/dz}{d\rho_T/dz} \right)^{-1} \quad (6)$$

**Formation of a salty glacial deep ocean: model limitations.** State-of-the-art AOGCMs lack fully interactive ice-sheet components and do not allow to simulate increased glacial melt in response to ocean warming. Furthermore, essential processes such as ice shelf–ice cavity ocean interactions are not included in the current generation of fully coupled AOGCMs<sup>74</sup>, which limits the ability to investigate a fuller spectrum of relevant processes and interactions such as the ‘pre-freshening’<sup>20</sup> of AABW via NADW. Furthermore, the ocean model component in our study and other global climate models is too coarse to represent the northward transport of dense shelf waters across the Antarctic shelf edge and their subsequent sinking to depth<sup>74</sup>, which represents a fundamental source of abyssal ocean circulation in the modern ocean. However, there is another mode associated with open-ocean convection<sup>75</sup> that has been identified as an additional surface to deep-ocean exchange process based on satellite data<sup>76,77</sup>. This modern mode with notable deep (>2,000 m) open-ocean convection between 90°S and 55°S under pre-industrial conditions has been diagnosed in more than two-thirds of the CMIP5 models. While this polynya-induced mode is not dominant today, it could have been the main source of AABW formation during full glacial conditions with a widely ice-covered continental shelf and a lower sea level than today. This situation might have favoured a glacial open-ocean type convection mode that is free to reach the seafloor<sup>75</sup>.

In that sense, a dominant glacial production regime of AABW by open-ocean convection, predominantly in the centres of the Weddell and Ross Sea gyres, as modelled in our model set-up<sup>78</sup> might capture the actual glacial water-mass conversion much better than under modern conditions.

To test the viability of the mechanism for the generation of a salty glacial deep ocean as proposed by Adkins<sup>20</sup>, we have evaluated how far our simulated control states are consistent with less melting of Antarctic ice by colder NADW during glacial conditions. In this context, we have applied the Potsdam Ice-Shelf Cavity Model (PICO)<sup>79</sup> to quantify changes in the subsurface melt rates. PICO is developed from the ocean box model of Olbers and Hellmer<sup>80</sup> and is implemented as a module of the Parallel Ice Sheet Model (PISM)<sup>81</sup> to simulate the vertical overturning circulation in ice-shelf cavities and thus enables the computation of subsurface melt rates (Extended Data Fig. 4) consistent with this circulation. We run the simulations at a spatial resolution of 16 × 16 km. The PICO settings are the standard settings that were used by Reese et al.<sup>79</sup> to obtain a realistic pre-industrial ice sheet. The PISM settings were taken from Sutter et al.<sup>82</sup>, and the atmospheric forcing is from regridded yearly RACMO2.3.2 m air temperature and precipitation fields<sup>83</sup>.

Driving this model set-up with the thermohaline ocean fields of experiment PID<sub>CTRL</sub> in equilibrium simulations of 200 kyr we end up with an aggregated melt flux of ~1,400 Gt yr<sup>-1</sup> (Extended Data Fig. 4a), which is not far from the observed value 1,500 ± 237 Gt yr<sup>-1</sup> for present-day conditions<sup>84</sup>. Keeping the same experimental set-up but replacing the thermohaline fields by the 40ka<sub>CTRL</sub> and LGM<sub>CTRL</sub> fields, results in a reduction of the aggregated melt flux by ~26% and 50%, respectively. This suggests that a reduction of Antarctic ice melt by colder NADW could indeed contribute to the generation and maintenance of a salty deep-ocean reservoir, especially since surface melt rates are negligible in the contribution to Antarctic ice-mass loss for modern climate conditions and colder climates. Hence, attempts should be made to incorporate this mechanism in future Earth system model approaches to explore essential interactions in particular at the ice–ocean interface, and it is open to question how far a fuller representation of the AABW formation mechanism would lead to a different model response. These investigations should be combined with the analysis of, for example, benthic foraminiferal  $\delta^{18}\text{O}$  and Mg/Ca records to discriminate sea-ice-derived versus Antarctic-ice-derived salinity stratification.

**Glacial deep-ocean salinity structure in models of different PMIP/CMIP generations.** Currently, there is no collective framework of perturbation experiments for glacial and interglacial conditions that would allow us to test our perturbation experiments within a model intercomparison. Nevertheless, we have generated an overview of simulated glacial deep-ocean salinity structure in LGM states (unperturbed) encompassing the last three PMIP/CMIP model generations (Extended Data Fig. 10).

Using the same model as employed in our study, the glacial deep-ocean salinity structure (including a PMIP2/CMIP4<sup>36,37</sup> model intercomparison) was investigated in the study of Zhang et al.<sup>29</sup>. From the six other models in that study, two models (CCSM3<sup>85</sup> and HadCM3M2<sup>86</sup>) also generated a stable glacial salt stratification (Extended Data Fig. 10). Note that PMIP utilized no specific protocol concerning the initial ocean condition for LGM simulations<sup>36,37</sup>, and notably the three models (including COSMOS) that were initialized from a previous glacial ocean state were the only ones to simulate a stable glacial deep-ocean salt stratification<sup>29</sup>. We suggest that this difference is a key to the different outcomes.

In model simulations at hand for the PMIP3/CMIP5<sup>38,39</sup> and PMIP4/CMIP6<sup>40</sup> generations, the north–south deep-ocean salinity contrast is dominated by a relatively salty North Atlantic, except within CCSM4<sup>97</sup>. In that model, saltier deep SO conditions are simulated, but these salinity characteristics do not extend into the Atlantic basin. Hence, the simulation of a salty glacial deep ocean via AABW at a global scale is not a coherent characteristic in LGM simulations with current coupled AOGCMs.

Where do we stand now? We think at least three possibilities should be taken into account for further research:

- The inference of a salinity-stratified glacial ocean via AABW has been questioned, and the scarce pore-water chlorinity data availability necessitates an increase in data constraints<sup>88</sup>.
- The simulation of a salty deep ocean seems to depend on conditions before the LGM, which therefore need to be incorporated into transient simulations rather than initiating models with LGM boundary conditions<sup>29</sup>. In particular, orbital configurations incorporating low obliquity and low atmospheric CO<sub>2</sub> concentrations in combination with Northern Hemisphere ice sheets in a later stage of a glacial cycle (but before the glacial maxima) should be considered as an alternative framework for simulating a salty glacial deep ocean<sup>29,30</sup>.
- Models need to be improved; accurate formation of AABW via shelf processes is an ongoing challenge<sup>41</sup>. Furthermore, important processes such as interactive ice sheets and shelf–ice cavity ocean interactions need to be included for detailed studies, and we note that the mechanism of deep-ocean salinification proposed by Adkins<sup>20</sup> (Extended Data Fig. 4) could work in tandem with increased sea-ice export to produce highly saline bottom waters. Recently, it has also been suggested that the export of icebergs from the SO to the Atlantic basin might modify interglacial-to-glacial salt balance between NADW and AABW<sup>42</sup>. As yet, these key processes and mechanisms are missing in coupled climate models, and therefore the conjecture that models that are simulating the salty reservoir might do it for the wrong reason cannot be excluded.

**Glacial deep-ocean salinity constraints.** The observation/reconstruction of extremely high glacial SO salinities and abyssal waters is very scarce, and a salinity value of  $37.1 \pm 0.2$  PSU at Shona Rise<sup>26</sup> represents a singularity that has been questioned by the analyses of pore fluid chlorinity/salinity data from deep-sea cores using estimation methods deriving from linear control theory<sup>88</sup>. Similarly to the ‘glacial’ state in Galbraith and Lavergne<sup>30</sup>, we simulate a lower salinity at this core location (36.04 PSU). Aside from this extreme salinity value, our simulated salinity in the deep Pacific between 3,500 m and 5,500 m is 36.03. This value is within the uncertainty of the reconstructed salinity ( $36.1 \pm 0.1$ ) based on chloride concentrations of deep Pacific bottom waters at five different core locations containing the same depth range<sup>27</sup>.

**Estimates for the importance of salinity changes.** To estimate the importance of interglacial-to-glacial salinity changes for the maintenance of deep-ocean stratification, we suggest two simple criteria. One criterion can help to evaluate the necessity of salinity changes in the evolution of the simulated glacial temperature structure. For example, in the PID<sub>CTRL</sub> and LGM<sub>CTRL</sub> experiment pair (Fig. 2), an unstable glacial situation applies if:

$$\frac{d\rho_{T\_LGM}}{dz} + \frac{d\rho_{S\_PID}}{dz} < 0 \quad (7)$$

That is, the glacial temperature decrease with depth is too weak to balance the effect of pre-industrial salinity decrease with depth on density (compare Extended Data Fig. 1).

While the latter criterion can help to assess the necessity of salinity changes, it does not answer the importance of a salinity reversal and the transition from an unstable to a stable salinity structure. Therefore, we additionally suggest to compare the LGM<sub>CTRL</sub> with a PID<sub>CTRL</sub> temperature gradient that would result from the shift of the unstable to a neutral salt stratification. This can be depicted in Fig. 3. Here the glacial cooling with depth (blue dot) is weaker than the intercept of the PID<sub>CTRL</sub> solution with the  $y$  axis. The corresponding criterion can be expressed as:

$$\frac{dT_{LGM}}{dz} > -\frac{d\rho/dz}{\alpha_{T\_PID} \rho_{0\_PID}} \quad (8)$$

That is, even if the unstable haline stratification contribution to  $d\rho_{-PID}/dz$  would be replaced by a neutral haline density stratification ( $d\rho_{s-PID}/dz=0$ ), the reduced

thermal stratification along solutions with constant  $d\rho_{-PID}/dz$ ,  $\alpha_{T\_PID}$  and  $\rho_{0\_PID}$  still requires a cooling with depth, which is stronger than  $dT_{LGM}/dz$ .

## Data availability

The ice-core data used in this paper are available via the referenced sources. The results from the PMIP3–CMIP5 and PMIP4–CMIP6 models can be found on the ESGF (Earth System Grid Federation) website. The PMIP2 Data Archive is available via <http://pmip2.lscce.ipsl.fr/>. Ice-sheet model data of this study are available at Pangaea (<https://doi.org/10.1594/PANGAEA.927126>). COSMOS model data that provide the basis for figures are archived on Zenodo (<https://doi.org/10.5281/zenodo.5137637>).

## Code availability

The standard model code of the comprehensive fully coupled atmosphere–ocean general circulation model (AOGCM), COSMOS (version COSMOS-landveg r2413, 2009) is available upon request from the Max Planck Institute for Meteorology in Hamburg (<https://www.mpimet.mpg.de>). Post-processing of model output and model data analysis has been performed with CDO (Climate Data Operators, versions 1.7.1, 1.9.5 and 1.9.10, <https://code.mpimet.mpg.de/projects/cdo>). We thank the developers of PISM and PICO for making their code publicly available at <https://github.com/pism/pism/tree/v1.1.1-gb8899b3>.

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## Acknowledgements

We thank colleagues in the Paleoclimate Dynamics group at the Alfred Wegener Institute Helmholtz Center for Polar and Marine Research (AWI) in Bremerhaven for general support and the AWI Computing Centre for keeping the supercomputer running. M. Fahrenberg is acknowledged for downloading PMIP4 data, and we thank X. Shi for providing LGM salinity data for AWI-ESM. We acknowledge financial support by PACES and REKLIM through the Helmholtz association, as well as from PalMod (01LP1504A and 01LP1915A) through the German Federal Ministry of Education and Research to G.K. We also acknowledge financial support from the National Science Foundation of China (41988101 and 42075047) to X.Z. and UK NERC (grants NE/J008133/1 and NE/L006405/1) to S.B. Development of PISM is supported by NSF grants PLR-1603799 and PLR-1644277 and NASA grant NNX17AG65G. Computational resources were made available by the infrastructure and support of the computing centre of the Alfred Wegener Institute in Bremerhaven and the DKRZ in Hamburg, Germany. The authors also thank all the modelling groups who provided the PMIP2, PMIP3 and PMIP4 output for analysis, WCRP, CMIP panel, PCMDI, ESGF infrastructures for sharing data, WCRP and CLIVAR for supporting the PMIP project. The Laboratoire des Sciences du Climat et de l'Environnement (LSCE) is acknowledged for collecting and archiving the PMIP2 model data. The PMIP2 Data Archive is supported by CEA, CNRS and the Programme National d'Etude de la Dynamique du Climat (PNEDC).

## Author contributions

G.K. conceived and developed the research. X.Z. designed and performed the COSMOS experiments. L.B.S. performed the ice-sheet model simulations. G.K. wrote the manuscript with help from S.B. All authors contributed to the interpretation and discussion of the results and to the preparation of the manuscript.

## Funding

Open access funding provided by Alfred-Wegener-Institute, Helmholtz Centre for Polar and Marine Research.

## Competing interests

The authors declare no competing interests.

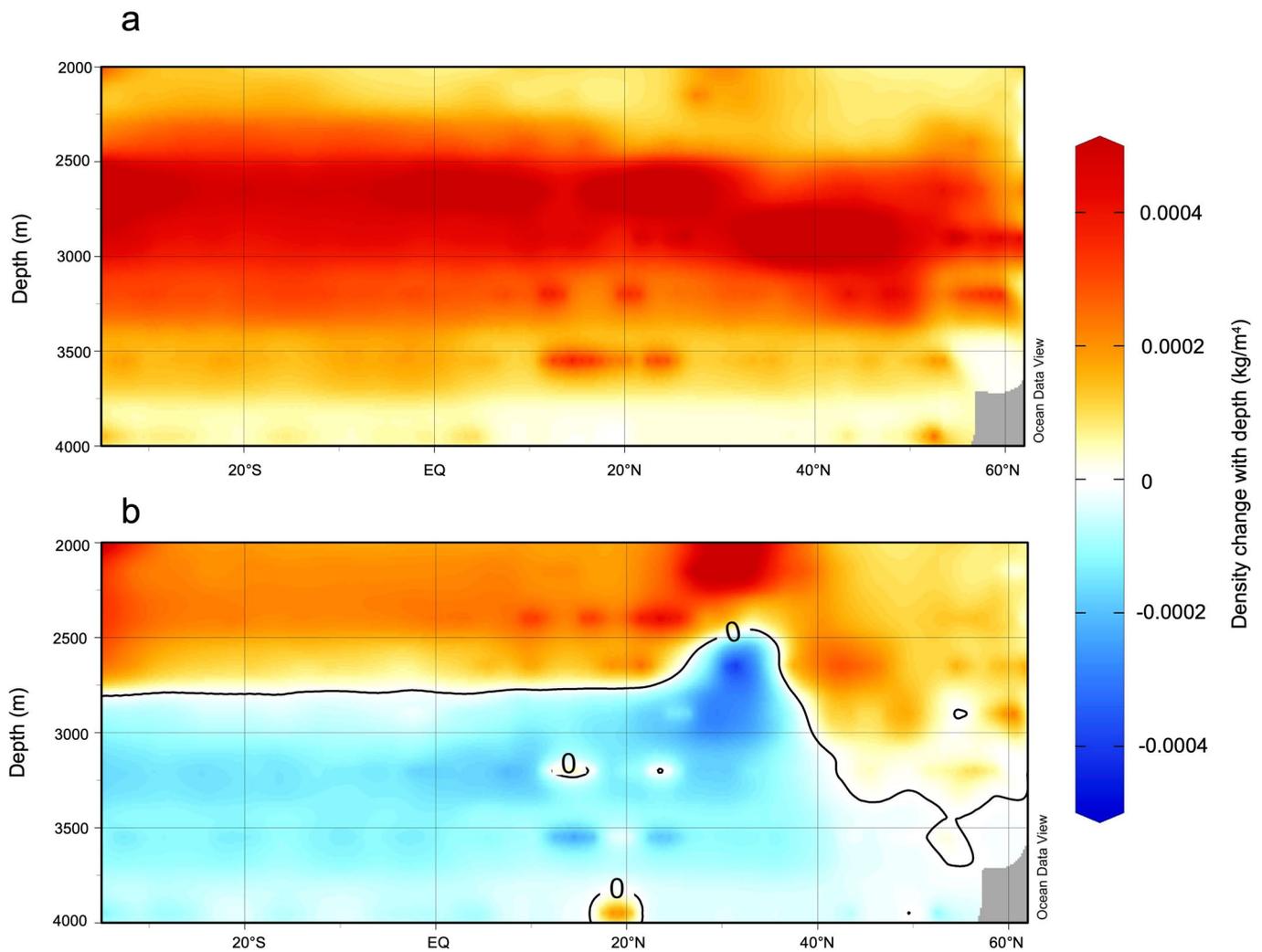
## Additional information

**Extended data** is available for this paper at <https://doi.org/10.1038/s41561-021-00857-3>.

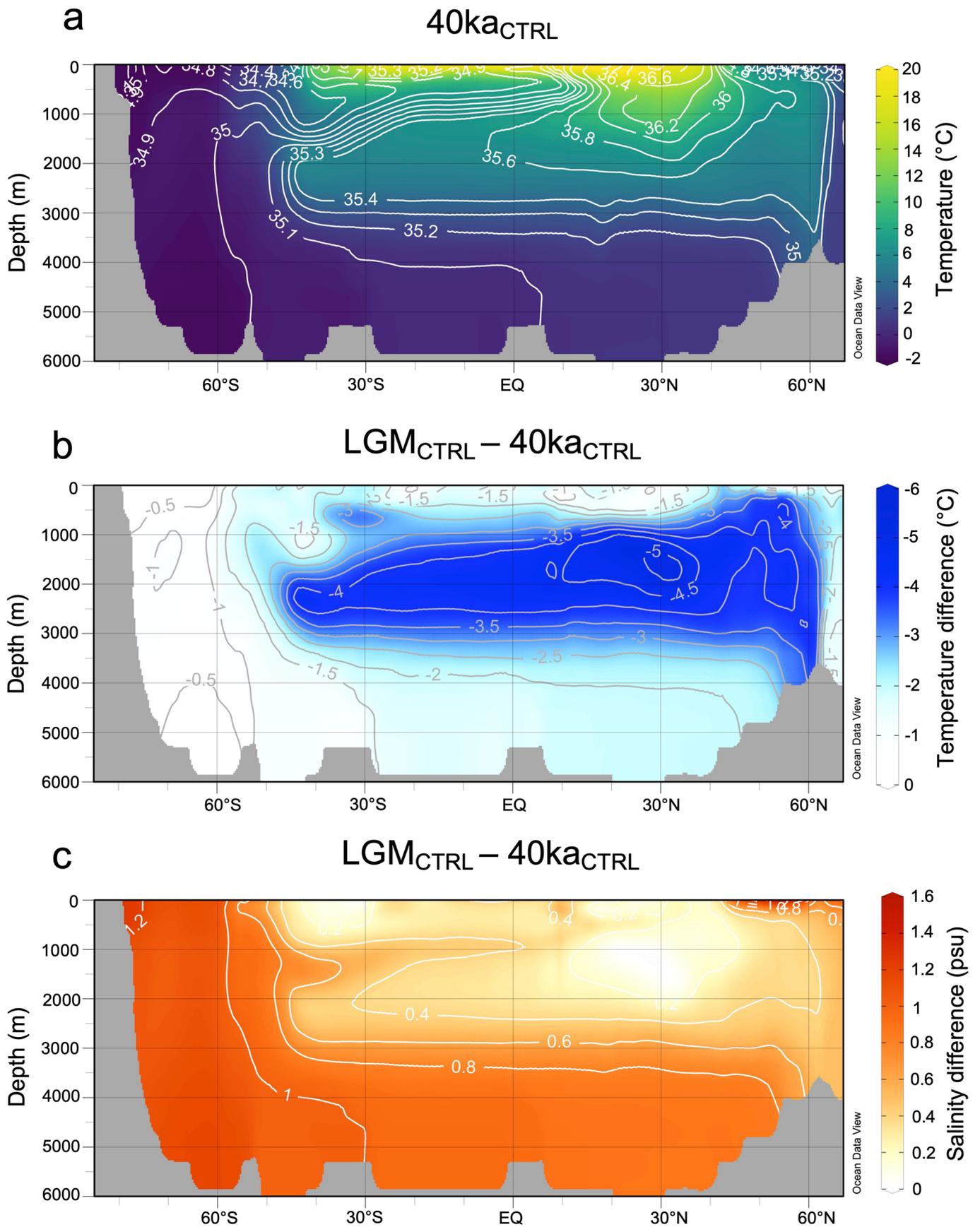
**Correspondence and requests for materials** should be addressed to Gregor Knorr.

**Peer review information** *Nature Geoscience* thanks the anonymous reviewers for their contribution to the peer review of this work. Primary Handling Editor: James Super.

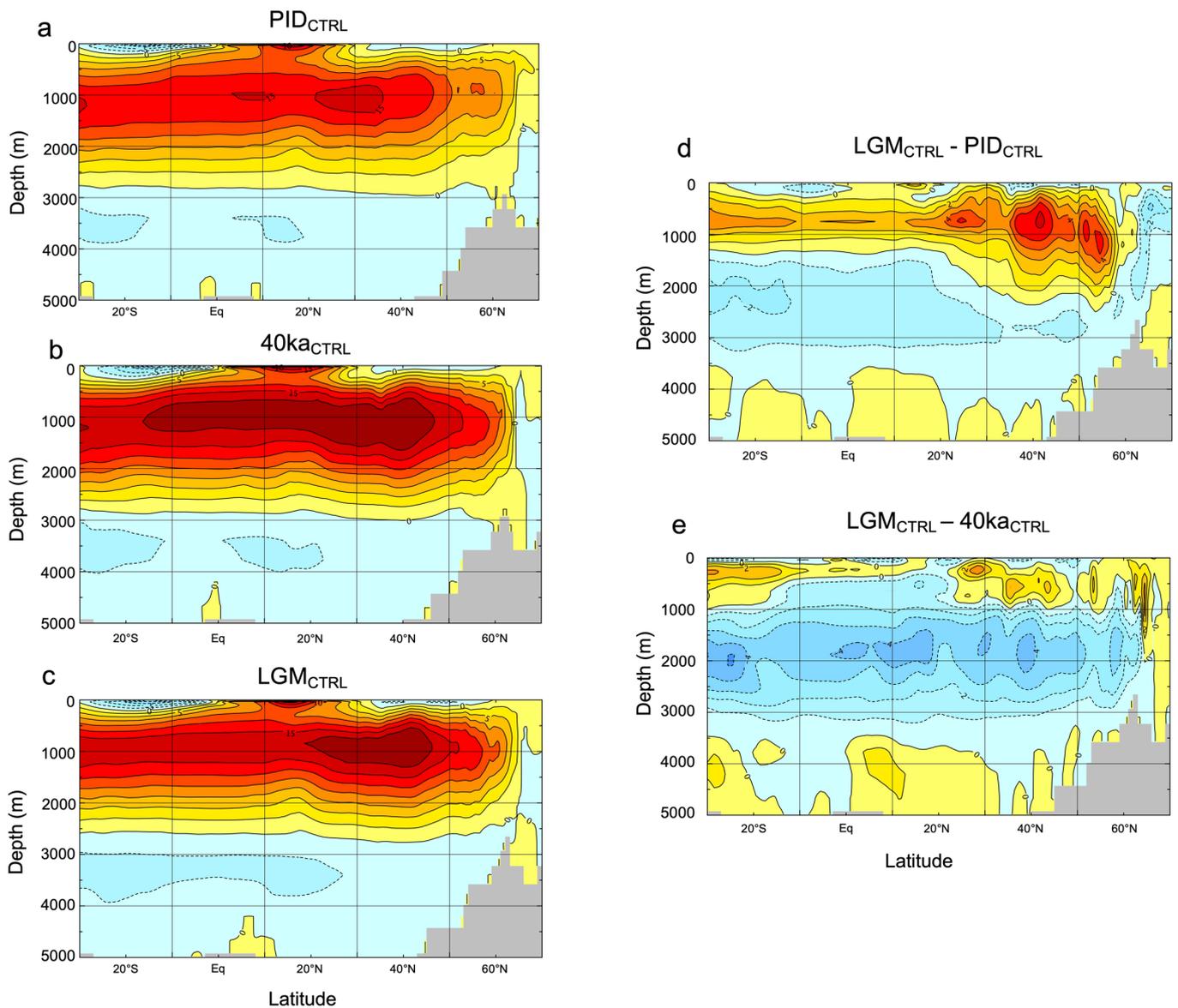
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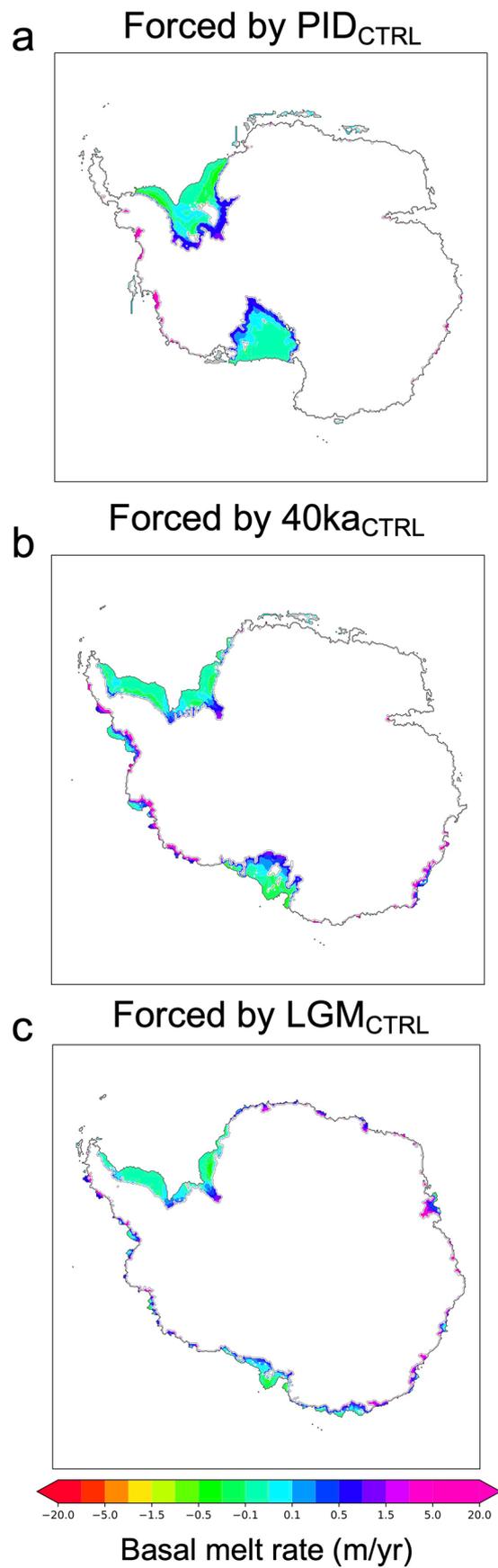
**Extended Data Fig. 1 | Simulated density changes with depth for pre-industrial conditions and a hybrid state of Last Glacial Maximum temperature and pre-industrial salinity conditions averaged in the Atlantic sector.** Shown are potential density ( $\sigma_t$ ) changes with depth (in  $\text{kg m}^{-4}$ ) for **a**) preindustrial conditions and **b**) a pseudo-modern state that combines simulated temperature and salinity fields for glacial and pre-industrial conditions, respectively. For the preindustrial state positive values indicate a stable stratification in the whole domain of panel a). In contrast, the zero-contour line in panel b) indicates that an unstable stratification (negative values) would arise for such a pseudo-hybrid state in large parts of the deep Atlantic below depths of ~2800 m. Simulated glacial and pre-industrial temperature and salinity fields are shown in Fig. 2.



**Extended Data Fig. 2 | Simulated 40ka temperature and salinity distribution and the Last Glacial Maximum anomalies to this state as zonal averages in the Atlantic sector.** Temperature ( $^{\circ}C$ ) and salinity (PSU, white contours) distribution for **a**) intermediate climate conditions ( $40ka_{CTRL}$ ). Panels **b**) and **c**) show the glacial - 40ka difference ( $LGM_{CTRL} - 40ka_{CTRL}$ ) for temperature and salinity, respectively. All fields represent 100-year averages.

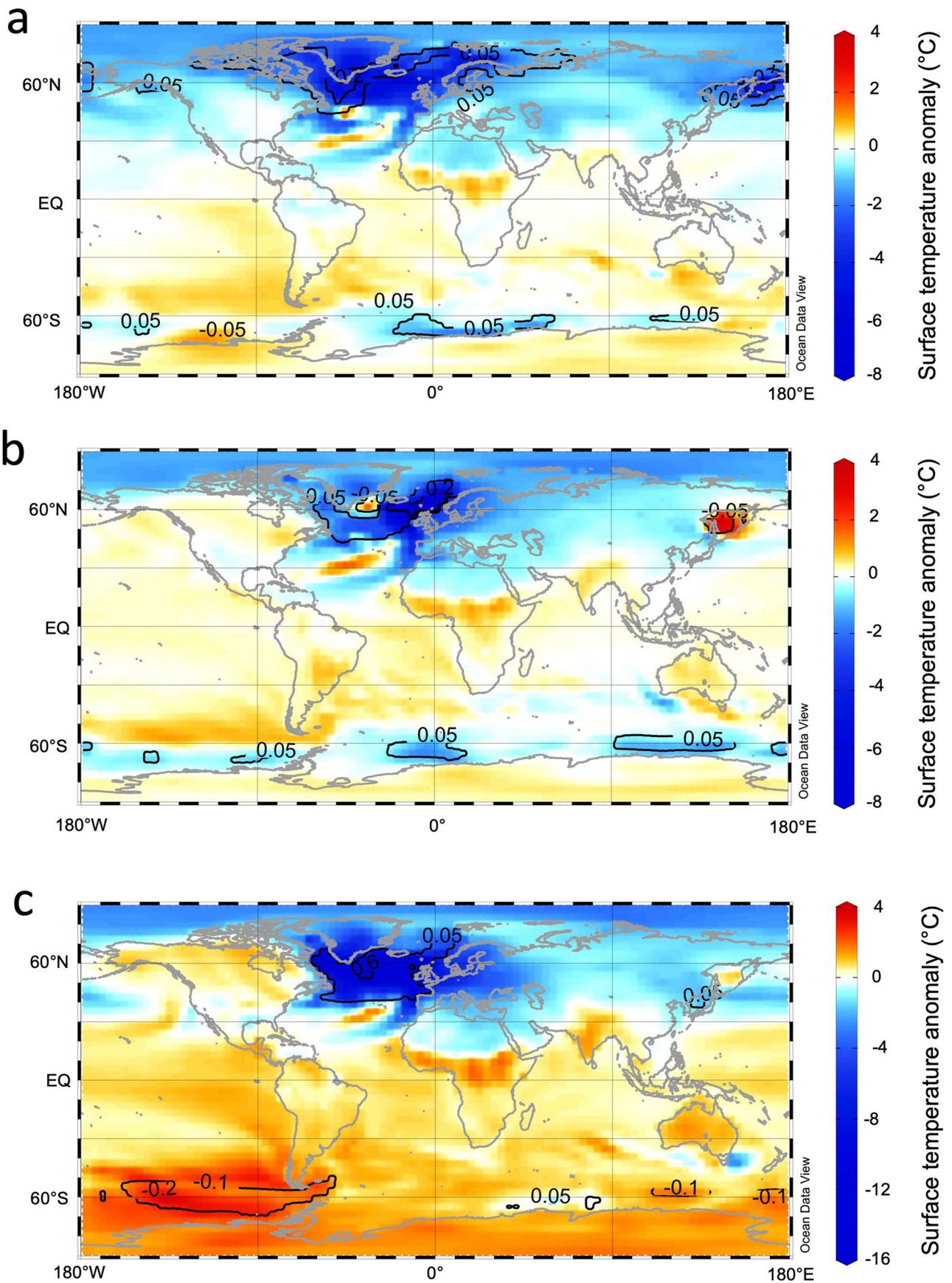


**Extended Data Fig. 3 | Simulated Atlantic overturning stream-function for different climate conditions and the differences between the glacial state and the interglacial and intermediate climate states.** Atlantic overturning stream-functions ( $S_v$ ,  $10^6 \text{ m}^3 \text{ s}^{-1}$ ) for **a**) interglacial ( $PID_{CTRL}$ ), **b**) intermediate ( $40ka_{CTRL}$ ) and **c**) glacial ( $LGM_{CTRL}$ ) climate conditions, as well as the differences **d**) between glacial and interglacial conditions ( $LGM_{CTRL} - PID_{CTRL}$ ) and **e**) between glacial and intermediate climate conditions ( $LGM_{CTRL} - 40ka_{CTRL}$ ). Shown are 100 year averages. Solid contours indicate clockwise NADW circulation and dotted contours anti-clockwise AABW circulation.



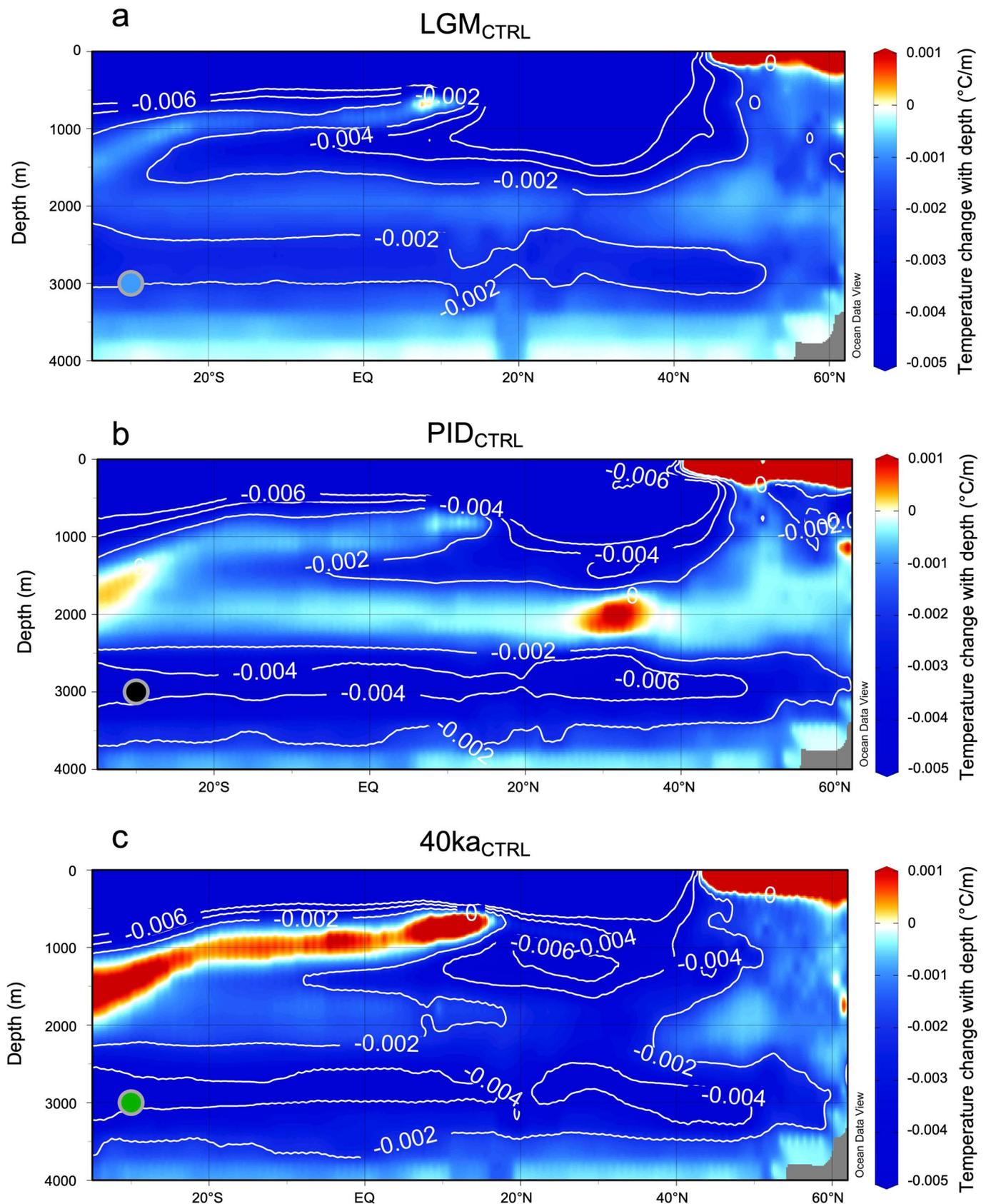
Extended Data Fig. 4 | See next page for caption.

**Extended Data Fig. 4 | Simulated basal melt rates for different thermo-haline conditions.** Basal melt rates as simulated using PISM-PICO for **a**) interglacial ( $PID_{CTRL}$ ), **b**) intermediate ( $40ka_{CTRL}$ ) and **c**) glacial ( $LGM_{CTRL}$ ) climate conditions. All simulations are identical except the thermo-haline fields that are applied as the ocean boundary condition (cf. Methods). These conditions stem from the 100 yr climatological mean of the respective climate states. Differences in ice sheet extent (white area) are caused by grounding line migration.

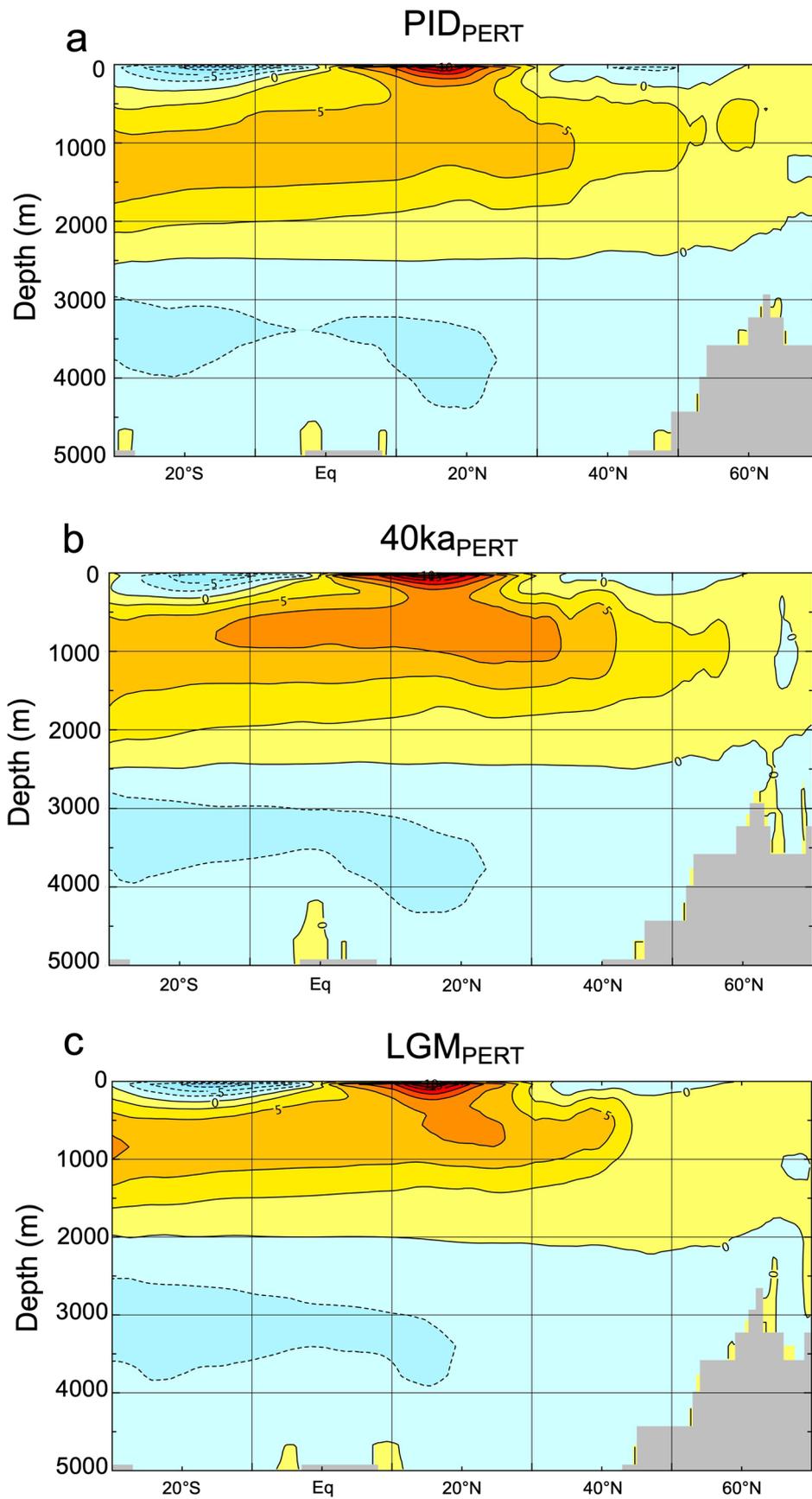


Extended Data Fig. 5 | See next page for caption.

**Extended Data Fig. 5 | Temperature and sea ice changes in response to North Atlantic freshwater perturbation.** Shown are surface temperature and sea-ice area fraction changes (iso-lines) for **a**) interglacial ( $PID_{PERT} - PID_{CTRL}$ ), **b**) 40ka ( $40ka_{PERT} - 40ka_{CTRL}$ ) and **c**) Last Glacial Maximum ( $LGM_{PERT} - LGM_{CTRL}$ ) conditions, averaged between model years 700-800. The north Atlantic cooling is most pronounced in the LGM perturbation experiment, which is linked to relatively pronounced sea ice cover increase and associated feedbacks. Together with the southern high latitude responses this is indicating a more pronounced interhemispheric seesaw response in surface temperature in the LGM perturbation experiment than under 40 ka or pre-industrial conditions.

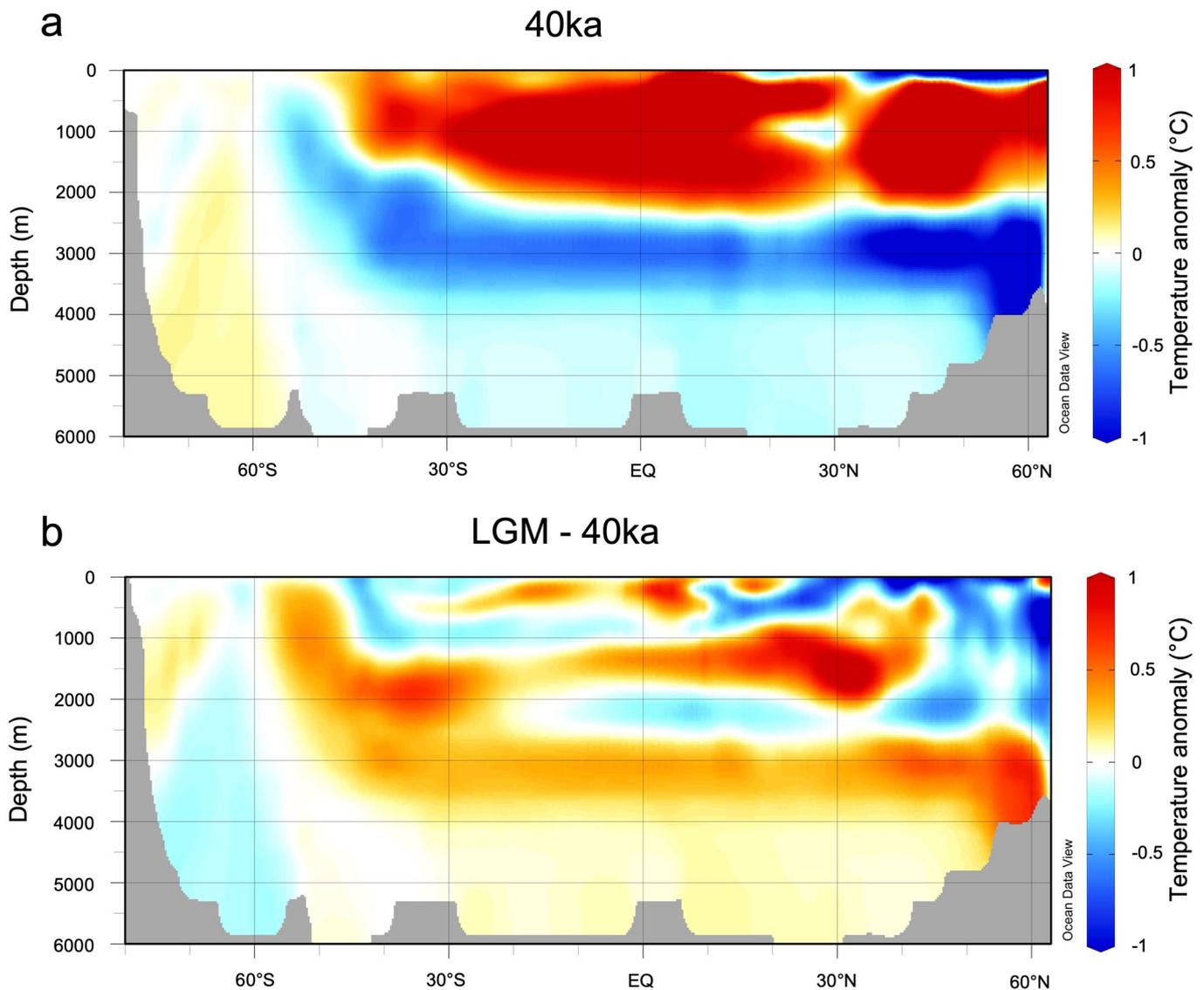


**Extended Data Fig. 6 | Vertical temperature gradient for different climate states in the Atlantic basin.** Zonally averaged temperature gradient ( $dT/dz$ ) for **a**) glacial (LGM<sub>CTRL</sub>), **b**) interglacial (PID<sub>CTRL</sub>) and **c**) intermediate (40ka<sub>CTRL</sub>) climate conditions. The coloured circles indicate the geographical location of the climate states shown in Fig. 3.

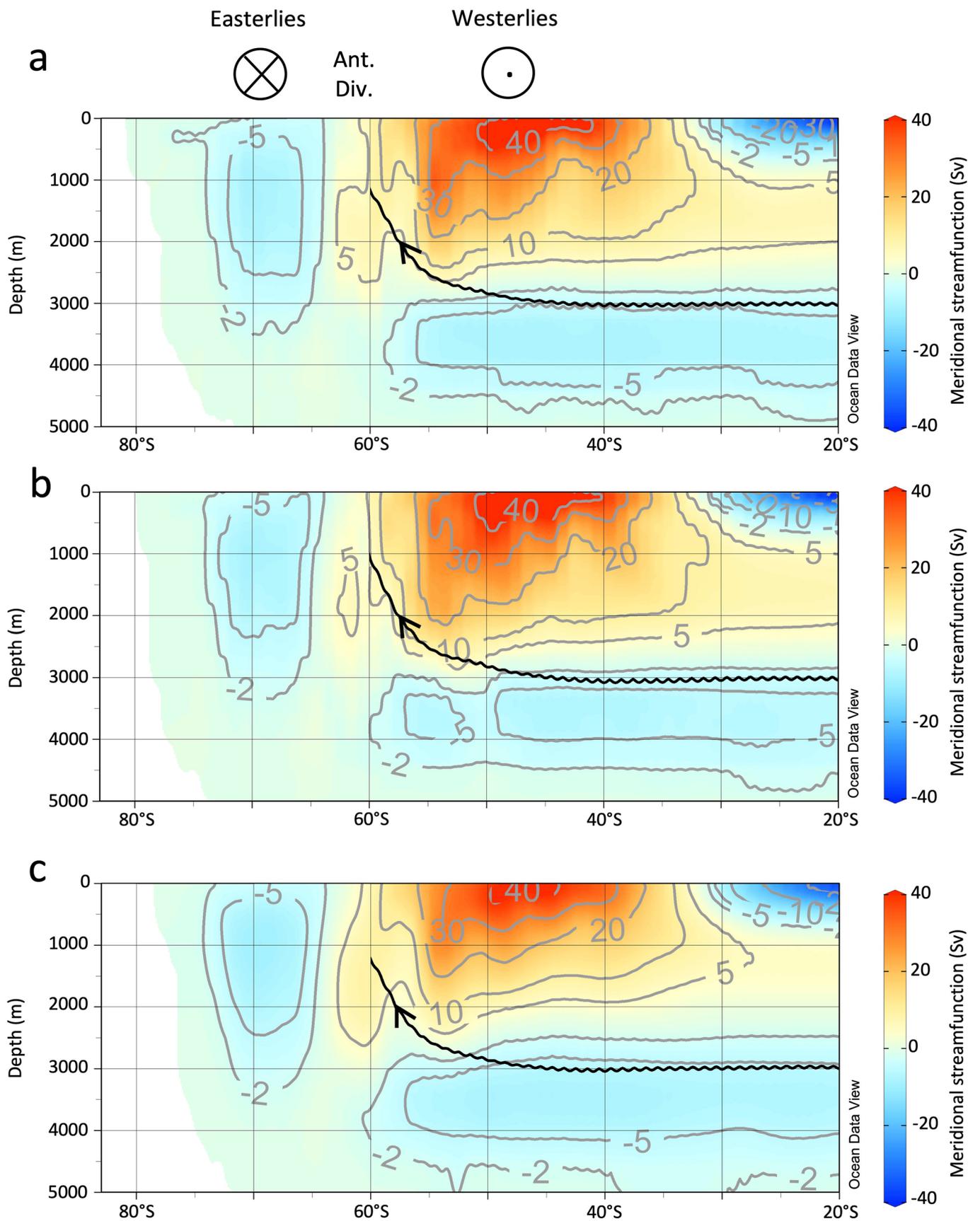


Extended Data Fig. 7 | See next page for caption.

**Extended Data Fig. 7 | Perturbed Atlantic overturning stream-function (Sv) at interglacial, intermediate and glacial climate conditions.** Atlantic overturning stream-function (Sv,  $10^6 \text{ m}^3 \text{ s}^{-1}$ ) in the perturbation experiments representing the average between model years 100-200 for **a**) interglacial ( $\text{PID}_{\text{PERT}}$ ), **b**) intermediate ( $40\text{ka}_{\text{PERT}}$ ) and **c**) glacial maximum ( $\text{LGM}_{\text{PERT}}$ ) conditions (cf. Fig. 4a). Solid contours indicate clockwise NADW circulation and dotted contours anti-clockwise AABW circulation.

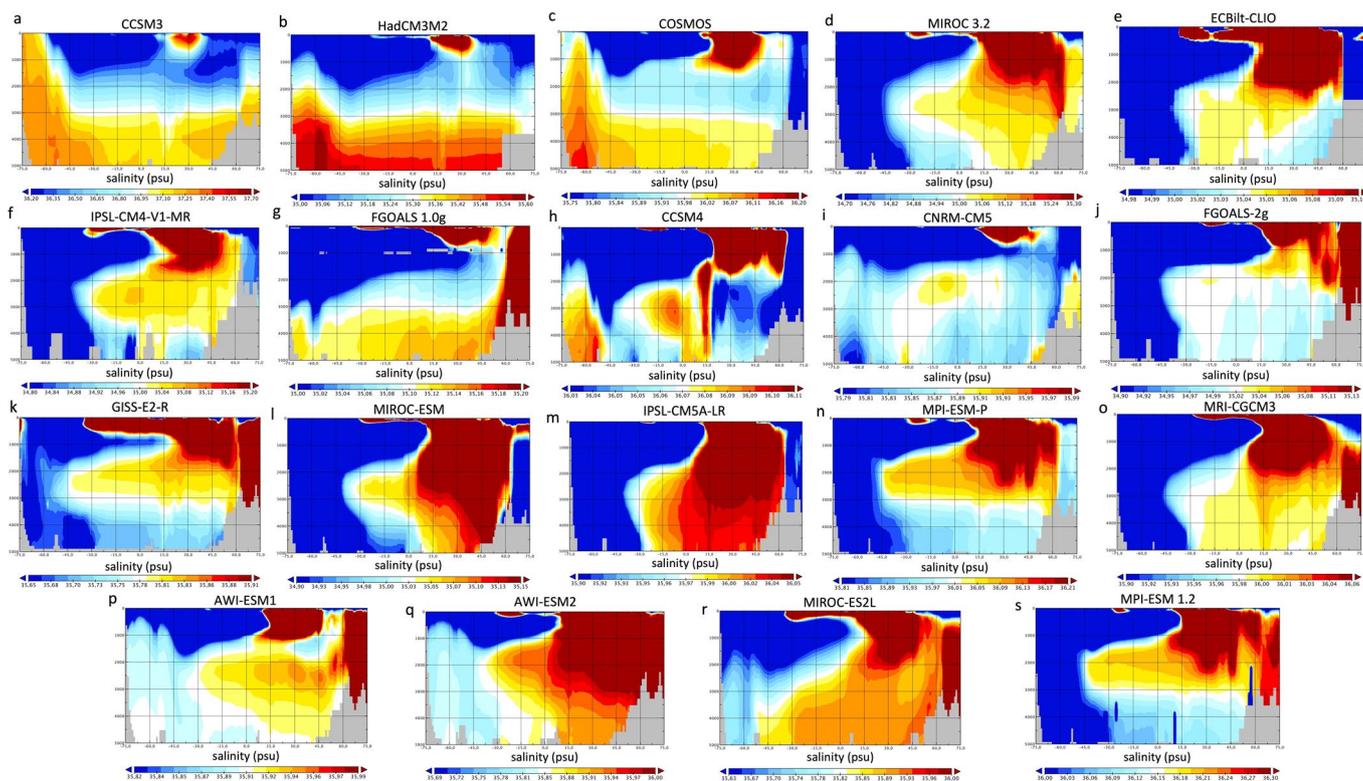


**Extended Data Fig. 8 | Temperature anomaly for the intermediate climate state and the anomaly difference to the glacial state.** Temperature anomaly by AMOC weakening **a**) for intermediate climate conditions ( $40\text{ka}_{\text{PERT}} - 40\text{ka}_{\text{CTRL}}$ ) between model years 100-200 as a 100 year mean, zonally averaged in the Atlantic sector. Panel **b**) shows the difference between glacial (Fig. 5b) and intermediate (panel (a) this figure) climate condition anomalies.



Extended Data Fig. 9 | See next page for caption.

**Extended Data Fig. 9 | Meridional overturning circulation in the Southern Ocean.** Shown are the Eulerian-mean meridional overturning stream function ( $S_v$ ,  $10^6 \text{ m}^3 \text{ s}^{-1}$ ) in depth coordinates in the perturbation experiments representing the average between model years 100-200 for **a**) interglacial ( $\text{PID}_{\text{PERT}}$ ), **b**) intermediate ( $40\text{ka}_{\text{PERT}}$ ) and **c**) glacial maximum ( $\text{LGM}_{\text{PERT}}$ ) conditions (cf. Extended Data Fig. 6). Positive (negative) values indicate clockwise (anticlockwise) circulation. The latitudinal variation in zonal wind-stress induces upwelling between the easterly and westerly wind maxima, and downwelling outside this corridor. Upwelled water can be returned as part of the wind-driven cells south or north of the Antarctic divergence (Ant. Div.) that is linked to a southbound Ekman transport under the influence of Polar easterlies and a northbound Ekman transport under the influence of the westerlies. Along these transport pathways heat and freshwater fluxes at the surface are part of the generalized circulation pattern that feed into the lower or upper limb of the recirculation, respectively. However, the actual upwelling transport largely occurs along surfaces of equal density, with wind and eddy processes playing a central role<sup>16</sup>. To illustrate such an isopycnal upwelling pathway, we additionally show the  $\sigma_3$  density surface (black line) that characterizes conditions at 3000 m depth and 30°S and upwells towards the latitude belt of the Antarctic Divergence that is subject to strong mixing during austral fall and winter seasons<sup>78</sup>.



**Extended Data Fig. 10 | Glacial Atlantic salinity structure in models of the last three PMIP/CMIP generations.** PMIP2/CMIP3 generation models are presented in panels **a**) to **g**), PMIP3/CMIP5 in panels **h**) to **o**) and PMIP4/CMIP6 in panels **p**) to **s**). The individual panels show the zonal mean salinities (PSU) in the Atlantic sector averaged over 100 years for the following models **a**) CCSM3<sup>85</sup>; **b**) HadCM3M2<sup>86</sup>; **c**) COSMOS<sup>29,31</sup>; **d**) MIROC 3.2<sup>89</sup>; **e**) ECBilt-CLIO<sup>90</sup>; **f**) IPSL-CM4\_V1-MR<sup>91</sup>; **g**) FGOALS-1g<sup>92</sup>; **h**) CCSM4<sup>87</sup>; **i**) CNRM-CM5<sup>93</sup>; **j**) FGOALS-2g<sup>94</sup>; **k**) GISS-E2-R<sup>95</sup>; **l**) MIROC-ESM<sup>96</sup>; **m**) IPSL-CM5A-LR<sup>97</sup>; **n**) MPI-ESM-P<sup>98</sup>; **o**) MRI-CGCM3<sup>99</sup>; **p**) AWI-ESM1<sup>100,101</sup>; **q**) AWI-ESM2<sup>102</sup>; **r**) MIROC-ES2L<sup>103,104</sup> and **s**) MPI-ESM 1.2<sup>105,106</sup>. Depths and latitudes are displayed in meters (m) and degrees north, respectively. The individual colorbars have been adjusted to reveal mid-depth and abyssal ocean salinity structures.